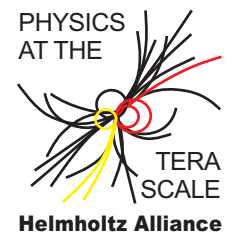




Weak boson scattering at the LHC

RADCOR 2009

Barbara Jäger
University of Würzburg

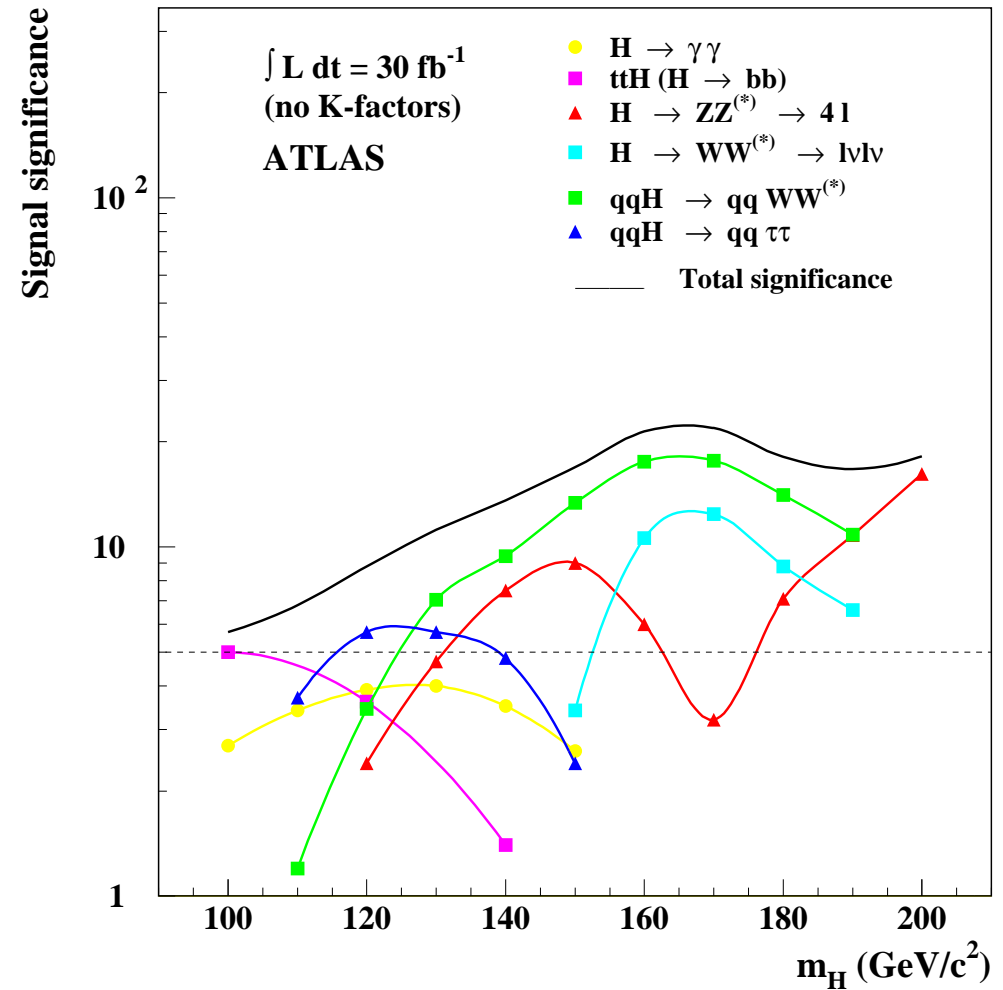


- ❖ weak boson fusion at the LHC:
 - Higgs production at high precision
 - $VVjj$ production @ NLO QCD
- ❖ strongly interacting gauge boson systems
 - heavy Higgs scenario
 - Warped Higgsless model
- ❖ summary & conclusions

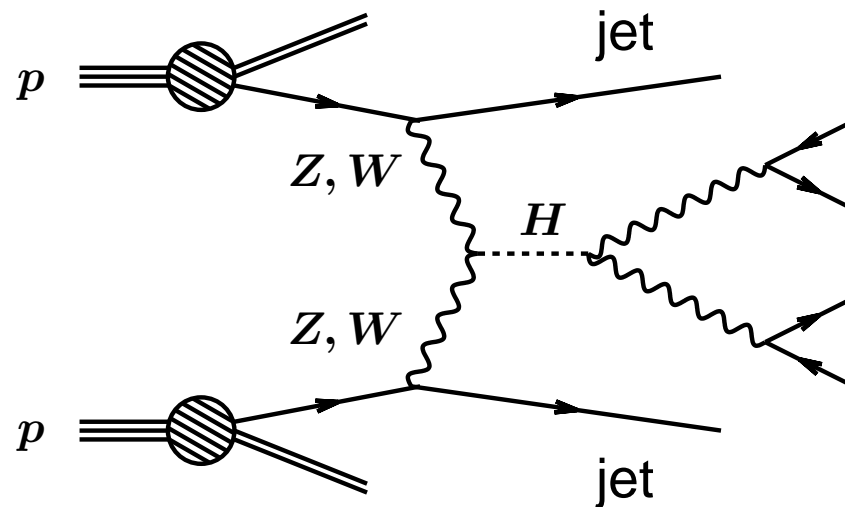
LHC discovery potential

most promising
production / decay modes
for Higgs search at the LHC
depend on M_H :

- VBF $qq \rightarrow qqH$ with
 $H \rightarrow W^\pm W^\mp \rightarrow e^\pm \mu^\mp p_T$
important for
 $M_H \gtrsim 110 \text{ GeV}$



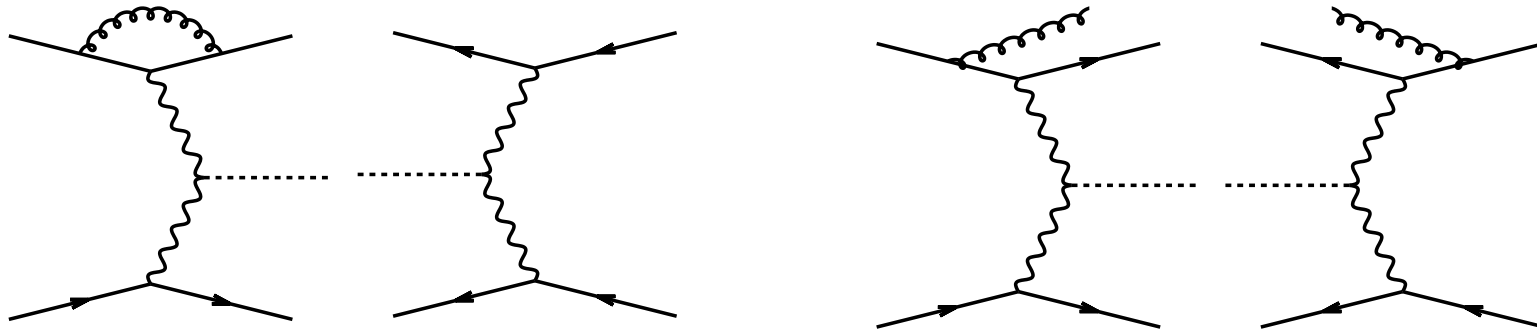
a promising search channel: VBF



suppressed color exchange between quark lines gives rise to

- ❖ little jet activity in central rapidity region
- ❖ scattered quarks \rightarrow two forward tagging jets (energetic; large rapidity)
- ❖ Higgs decay products typically between tagging jets

Higgs production in VBF @ NLO QCD



NLO QCD:

inclusive cross section:

Han, Valencia, Willenbrock (1992)

distributions:

Figy, Oleari, Zeppenfeld (2003)

Berger, Campbell (2004)



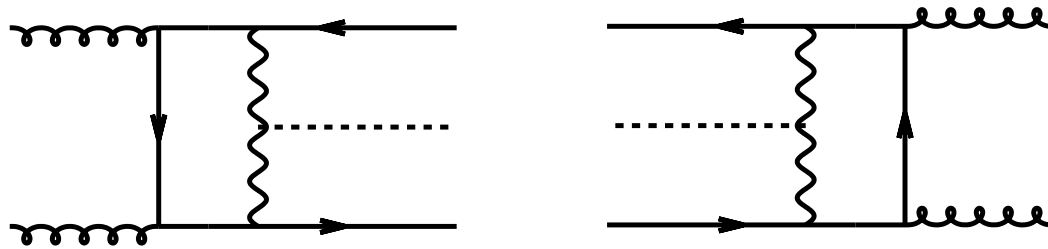
**NLO QCD corrections
moderate**

and well under control
(order 10% or less)

higher orders of QCD in VBF

Harlander, Vollinga, Weber (2007):

gauge invariant, finite sub-class of virtual
two-loop QCD corrections to $pp \rightarrow Hjj$ via VBF



important due to large
gluon luminosity at LHC?

$$gg \rightarrow q\bar{q}H, q\bar{q} \rightarrow ggH,$$
$$qg \rightarrow qgH, \bar{q}g \rightarrow \bar{q}gH$$

minimal set of cuts: $\sigma_{\text{gluon}}^{2\text{-loop}} \sim 2\%$ of $\sigma_{\text{VBF}}^{\text{LO}}$

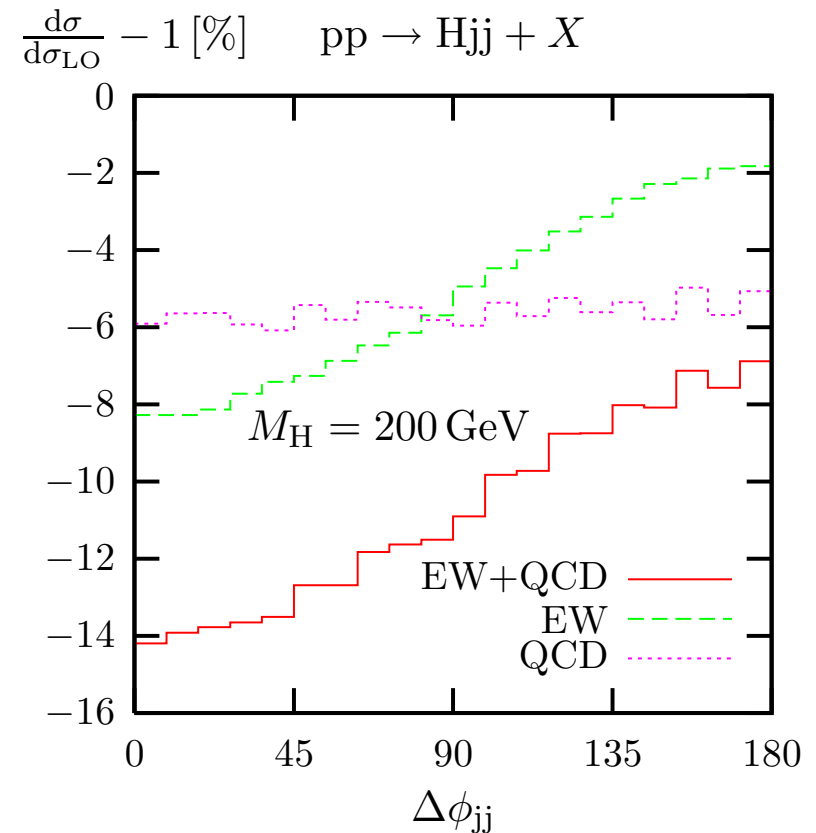
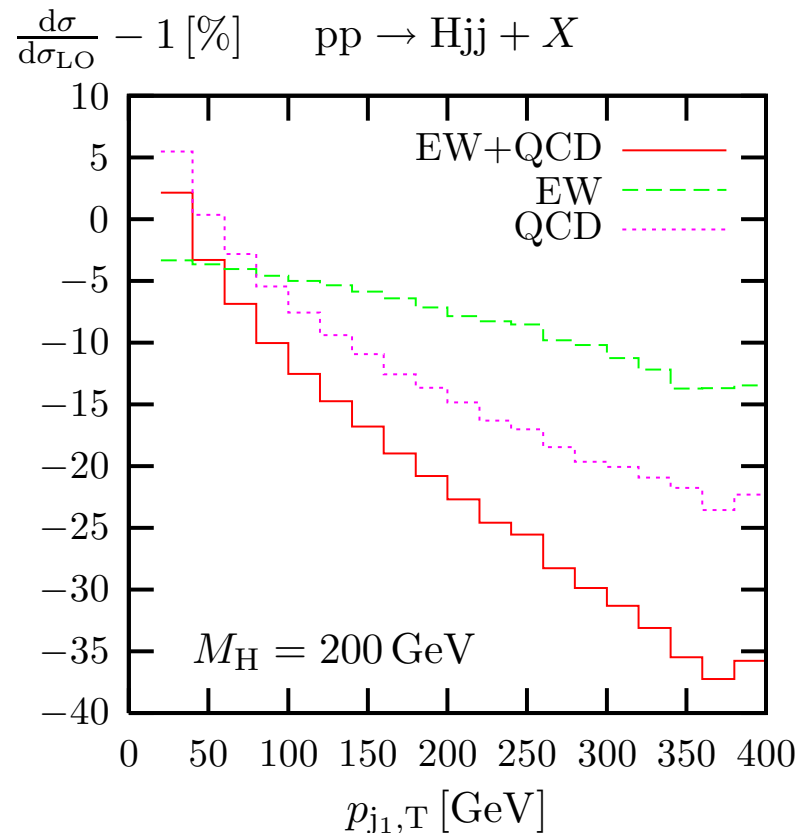
VBF cuts: relative suppression by additional order of magnitude

Higgs production in VBF @ NLO EW

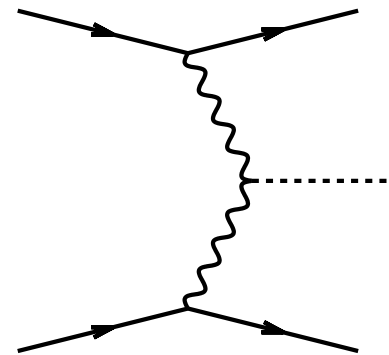
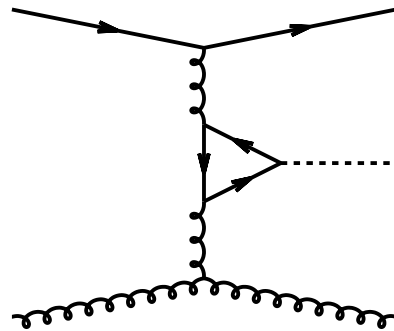
Ciccolini, Denner, Dittmaier (2007):

NLO EW corrections to inclusive cross sections and distributions

- ➔ **NLO EW corrections non-negligible**, modify K factors and distort distributions by up to 10%



VBF Higgs signal can be faked by
double real corrections to $gg \rightarrow H$ (“gluon fusion”)



complete LO calculation (including pentagons):

Del Duca, Kilgore, Oleari, Schmidt, Zeppenfeld (2001)

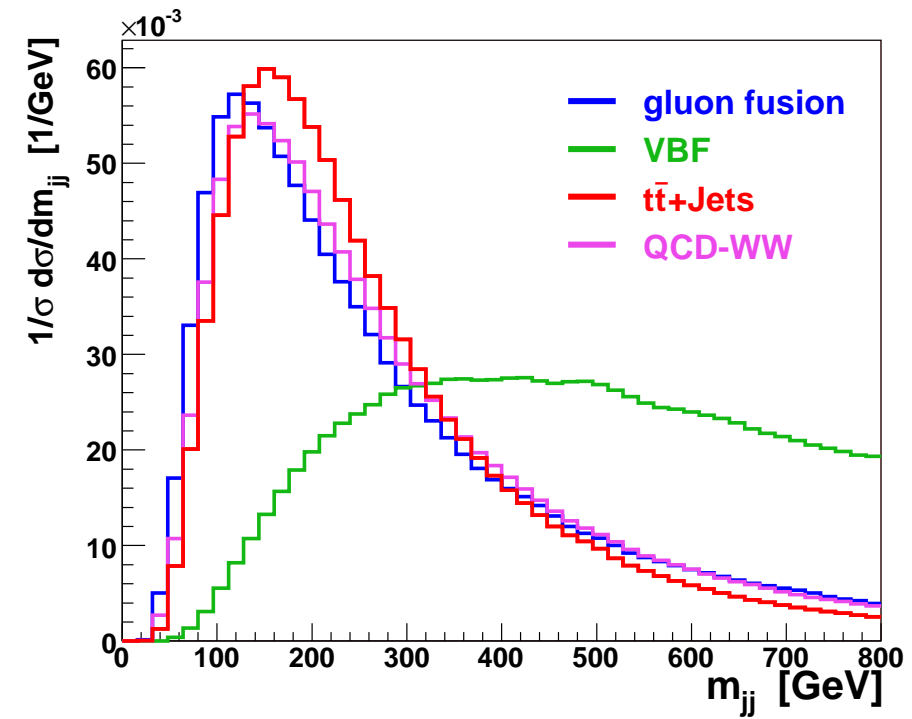
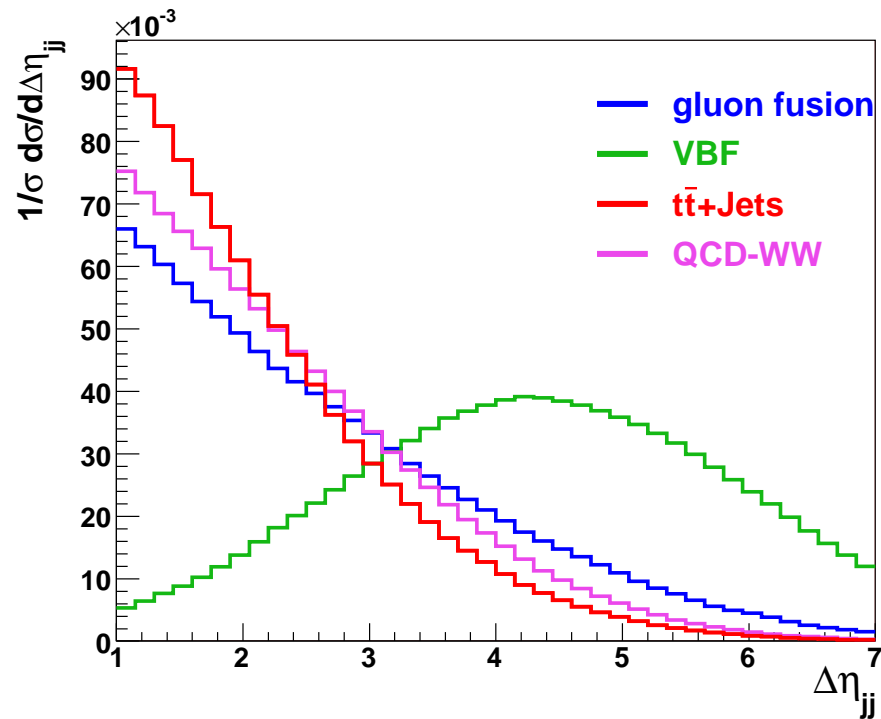
NLO QCD calculation in $m_t \rightarrow \infty$ limit:

Campbell, Ellis, Zanderighi (2006)

need to understand phenomenology of both processes to
distinguish between them

$pp \rightarrow Hjj$: VBF versus GF

apply **cuts** to separate VBF signal from gluon fusion (GF) background



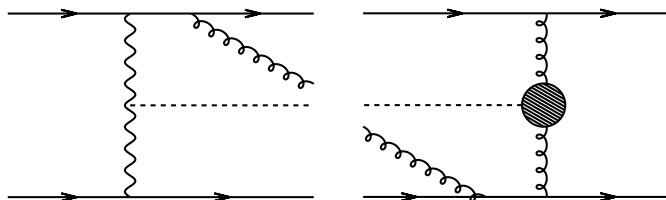
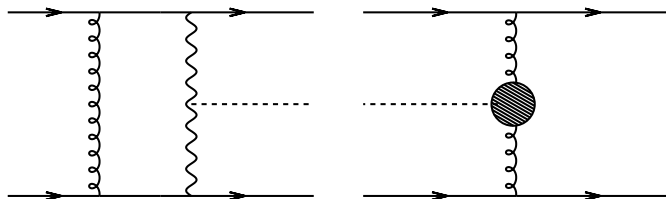
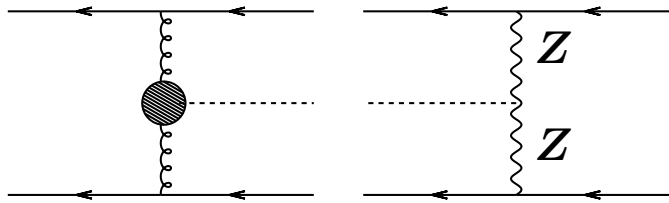
Klämke, Zeppenfeld (2007)

$pp \rightarrow Hjj$ via $VBF \times GF$

can $VBF \times GF$ interference pollute the clean VBF signature?

Georg (2005) & Andersen, Smillie (2006):

- ◆ neutral current graphs
(no charged current interference)
- ◆ identical quark contributions
with $t \leftrightarrow u$ crossing

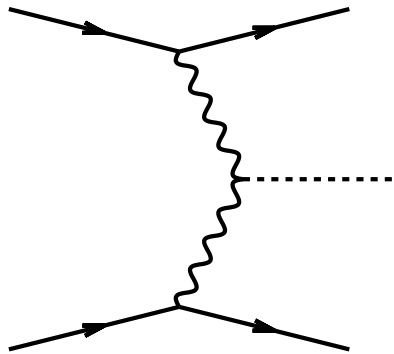


Andersen et al. (2007)

Bredenstein, Hagiwara, B. J. (2008):

- ◆ strong cancelation effects
between contributions of
different flavor

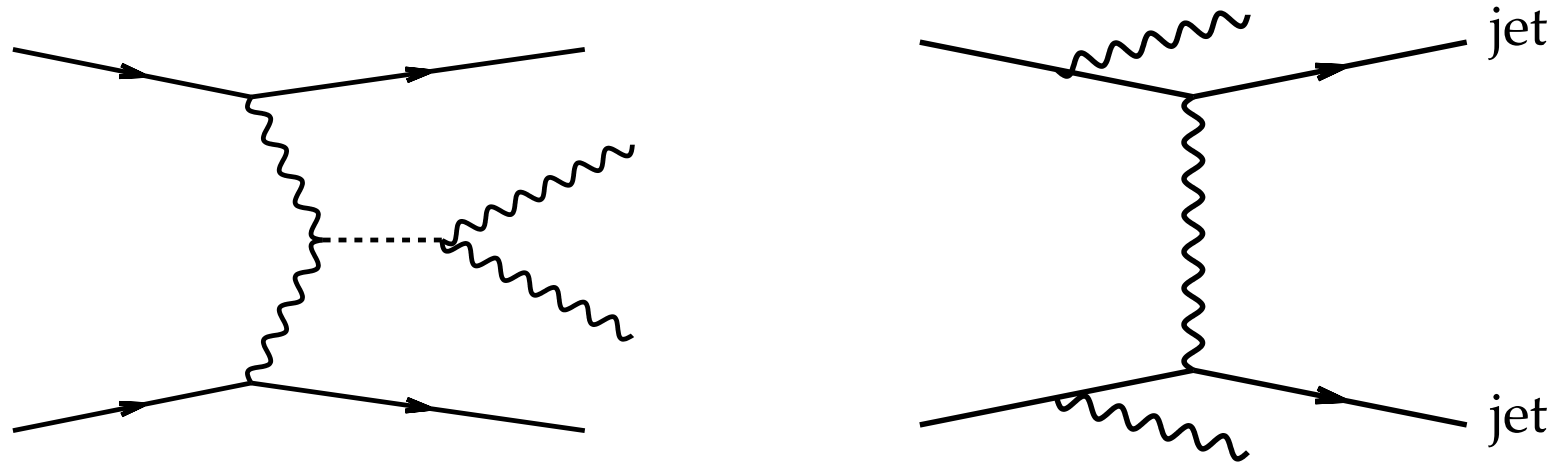
☞ interference effects are **completely negligible**



- ❖ $pp \rightarrow Hjj$ via VBF under excellent control
 - ❖ QCD & EW NLO corrections at 10% level
 - ❖ dominant NNLO QCD corrections small
 - ❖ SUSY QCD corrections small
- ☞ *Michael Rauch's talk*
- ❖ interference with GF Hjj production negligible
 - ❖ small PDF uncertainties $\lesssim 4\%$

but: establishing a signal for the Higgs boson in VBF
requires also
calculation of **background contributions**

EW $VVjj$ production



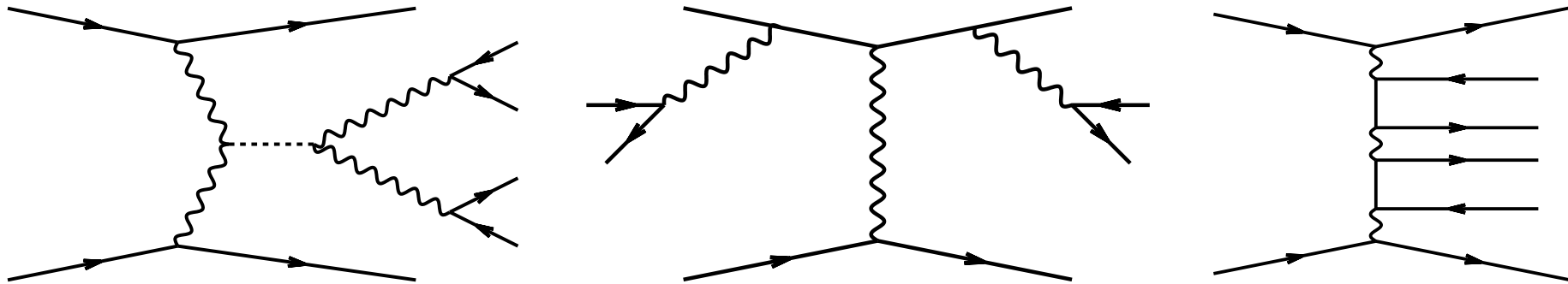
$pp \rightarrow VV + jj$ via VBF

irreducible background to Higgs signal process
in $H \rightarrow VV$ decay mode



accurate predictions essential

EW $VVjj$ production



experiment: don't observe $VVjj$ final state, but
hadronic or leptonic decay products

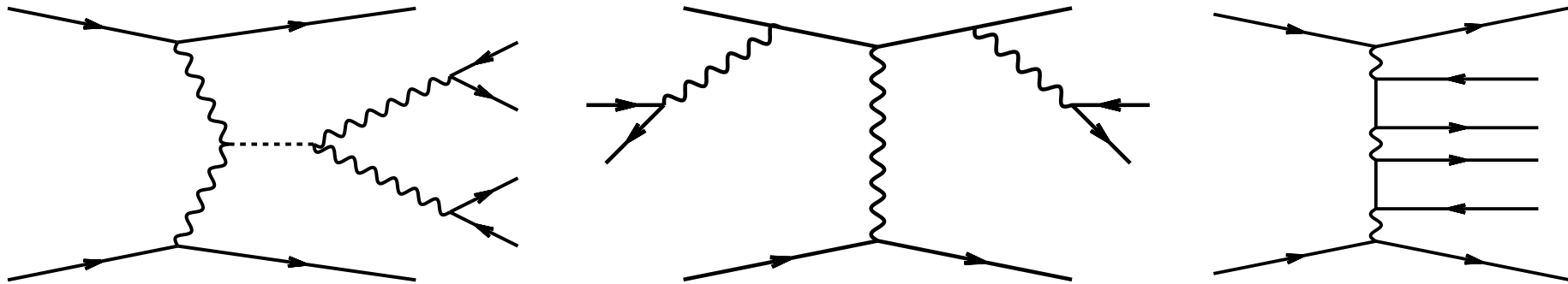
$4\text{jets} + jj$

high statistics
large backgrounds

$4\text{leptons} + jj$

low statistics
clean signature

EW $VVjj$ production



experiment: don't observe $VVjj$ final state, but
hadronic or leptonic decay products

focus on
 $pp \rightarrow 4\text{leptons} + jj$
via VBF
(short: " $pp \rightarrow VVjj$ ")

4leptons + jj

low statistics
clean signature

need **stable, fast & flexible Monte Carlo program** allowing for

❖ computation of various jet observables for
 W^+W^-jj , $ZZjj$, $W^\pm Zjj$, and $W^\pm W^\pm jj$

production via VBF at NLO-QCD accuracy

(leptonic decay correlations are fully taken into account)

❖ straightforward implementation of cuts

C. Oleari, D. Zeppenfeld, B. J. (2006)

G. Bozzi, C. Oleari, D. Zeppenfeld, B. J. (2007)

C. Oleari, D. Zeppenfeld, B. J. (2009)

[c.f. also Dieter Zeppenfeld's talk]

$pp \rightarrow \ell\bar{\ell} \ell'\bar{\ell}' jj$: the leading order

need to compute numerical value for

$$|\mathcal{M}_B|^2 = \left| \text{[diagram 1]} + \text{[diagram 2]} + \text{[diagram 3]} + \dots \right|^2$$

at each generated phase space point in 4 dim (finite)

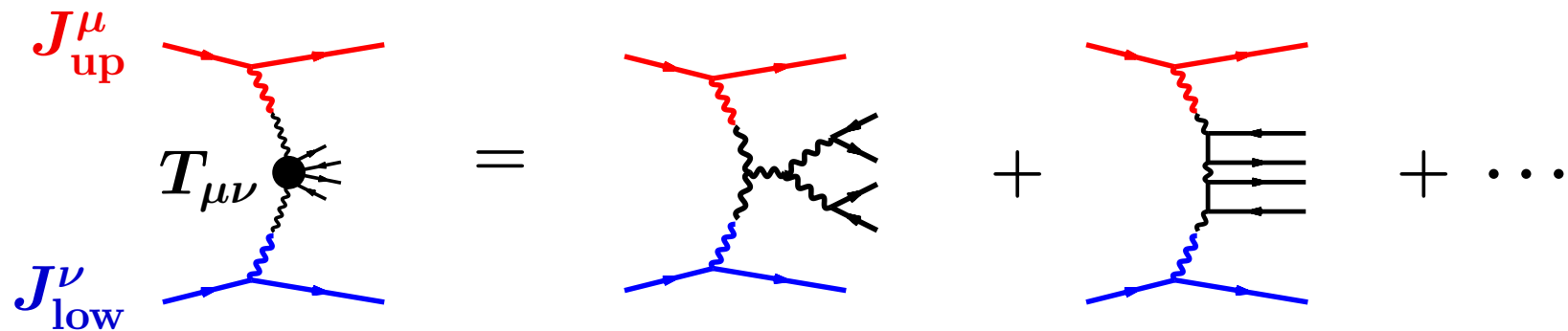
... depending on leptonic final state: up to 580 diagrams

essential: organize calculation **economically**



- employ amplitude techniques to **evaluate \mathcal{M} first** (numerically) for specific helicities of external particles, **then square**
- **avoid multiple evaluation** of recurring building blocks

develop modular structure with



... and evaluate each building block only once per phase space point
(related sub-diagrams, various flavor combinations,
crossed processes ...)

such recycling is used to a very small extent by
automatized programs like MadGraph/MadEvent

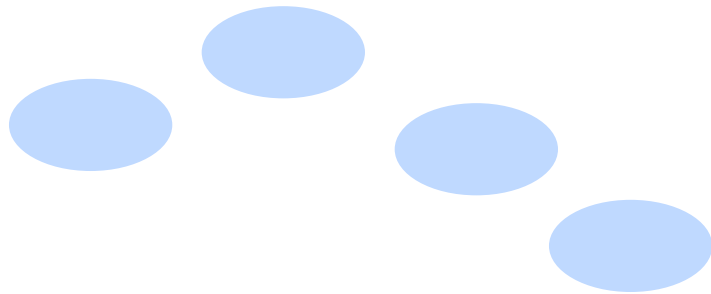
not included:

- ❖ interference effects from **identical fermions**
(t - and u -channel interference)
 - ❖ identical flavor **annihilation processes**
(s -channel contributions)
- ... strongly suppressed in phase-space region
where VBF can be observed experimentally

Oleari, Zeppenfeld (2003), Georg (2005)

Andersen, Smillie (2006)

Ciccolini, Denner, Dittmaier (2007)

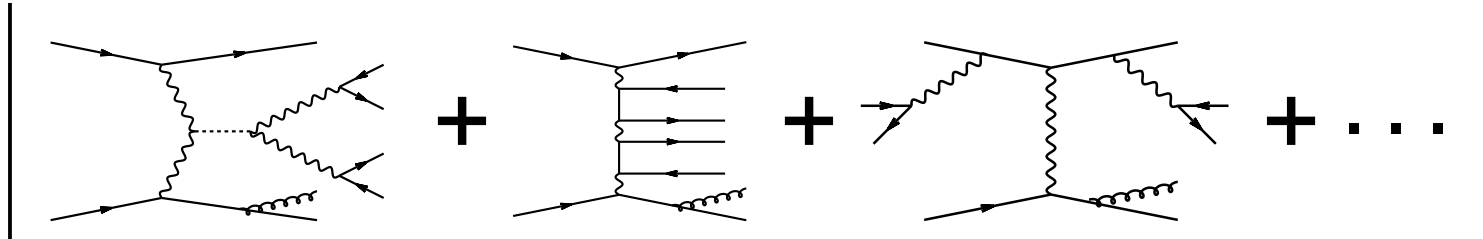


the next-to-leading
order:

- real emission
- subtraction terms
- virtual corrections

real emission contributions

needed: numerical value for almost 3000 diagrams

$$|\mathcal{M}_R|^2 = \left| \begin{array}{c} \text{diagram 1} \\ + \\ \text{diagram 2} \\ + \\ \text{diagram 3} \\ + \dots \end{array} \right|^2$$
The equation shows the squared magnitude of the real emission amplitude, |M_R|^2, as a sum of several Feynman diagrams enclosed in large vertical bars with a superscript 2. The diagrams include: 1) a tree-level diagram with a wavy line and a dashed line; 2) a tree-level diagram with a wavy line and a solid line; 3) a tree-level diagram with a wavy line and a solid line in a different configuration; and 4) an ellipsis indicating further diagrams.

complication: real emission contribution **diverges** as unobserved parton becomes **soft or collinear**

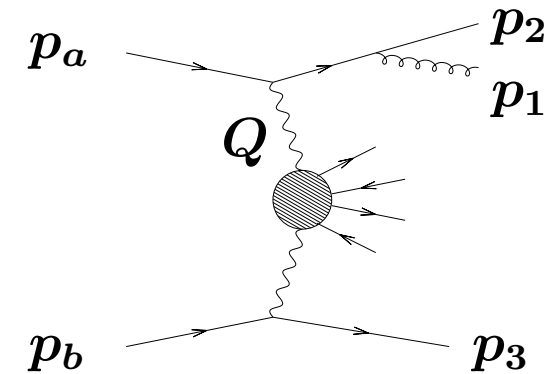


- ❖ analytic calculation: divergencies canceled directly by respective singularities in virtual contributions
- ❖ numerical approach: apply **subtraction formalism** (phase space slicing, dipole subtraction, ...)
- ☞ divergencies are absorbed by auxiliary **counterterms**

dipole subtraction: $qq' \rightarrow qq'(g)VV$ via VBF

continuous interpolation
between
soft and collinear
gluon radiation:

$$\sim \frac{x^2 + z^2}{(1-x)(1-z)} |\mathcal{M}_B(\tilde{p})|^2$$



analytical integration over
one-particle phase space:

$$\sim |\mathcal{M}_B(p)|^2 \left[\frac{2}{\epsilon^2} + \frac{3}{\epsilon} + \text{const.} \right]$$

Catani, Seymour (1996)

$$\sigma^{NLO} = \int_{m+1} [d\sigma_{\epsilon=0}^R - d\sigma_{\epsilon=0}^A] + \int_m \left[d\sigma^V + \int_1 d\sigma^A \right]_{\epsilon=0}$$

virtual contributions

$$\mathcal{M}_V = \text{[diagram 1]} + \text{[diagram 2]} + \text{[diagram 3]} + \dots$$

$$= \mathcal{M}_B F(Q) \left[-\frac{2}{\epsilon^2} - \frac{3}{\epsilon} \right] + \tilde{\mathcal{M}}_V^{finite}$$

determined numerically

[c. f. Denner, Dittmaier (2002,2005)]

combination of real emission, virtuals,
and subtraction terms:
poles canceled analytically → **finite** results

phase-space integration can be performed numerically (Vegas)

Monte Carlo program for **cross sections and distributions**
which allows for the implementation of
realistic experimental selection cuts

➔ embedded in more general framework
for various VBF-type processes

vbfnlo

publicly available from

<http://www-itp.physik.uni-karlsruhe.de/~vbfnloweb/>

use design energy of 14 TeV, apply k_T jet algorithm,
CTEQ6 parton distributions, and typical VBF cuts:

tagging jets

$$p_{Tj} \geq 20 \text{ GeV}, \quad |y_j| \leq 4.5,$$

$$\Delta y_{jj} = |y_{j_1} - y_{j_2}| > 4,$$

$$M_{jj} > 600 \text{ GeV}$$

jets located in opposite hemispheres

$$p_{T\ell} \geq 20 \text{ GeV}, \quad |\eta_\ell| \leq 2.5,$$

$$\Delta R_{j\ell} \geq 0.4,$$

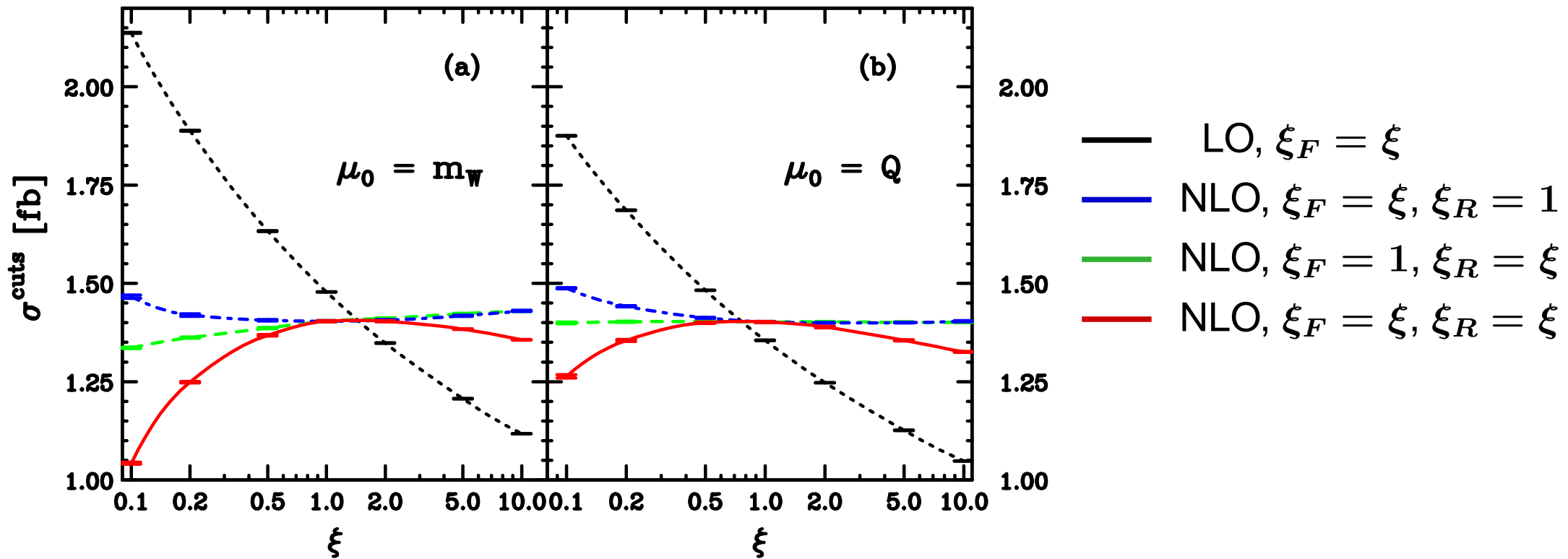
$$y_{j,\min} < \eta_\ell < y_{j,\max}$$

charged leptons

scale uncertainty: $pp \rightarrow W^+W^+jj$

choose default scale $\mu_0 = m_W$ or $\mu_0 = Q$
set $\mu_R = \xi_R \mu_0$ and $\mu_F = \xi_F \mu_0$, with variable ξ

Oleari, Zeppenfeld, B. J. (2009)

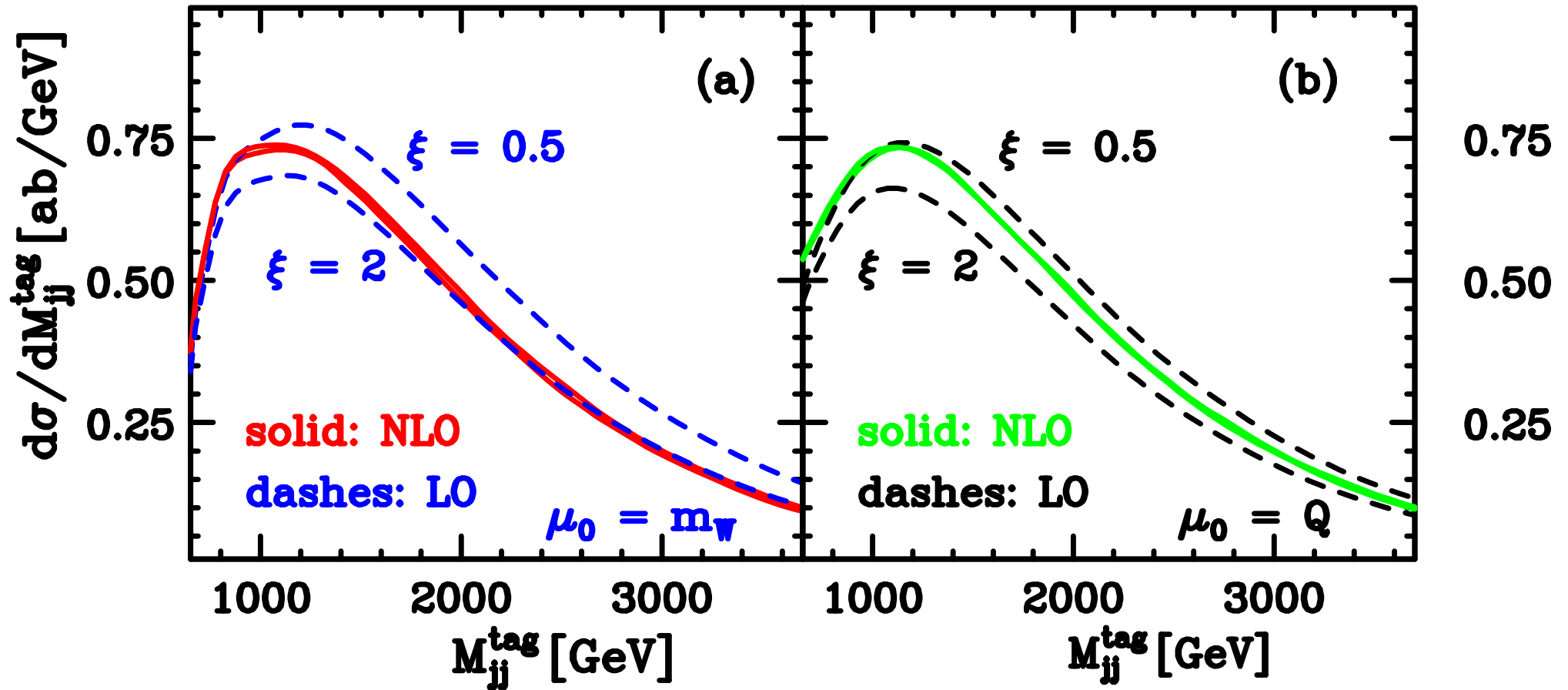


LO: no control on scale

NLO QCD: scale dependence strongly reduced

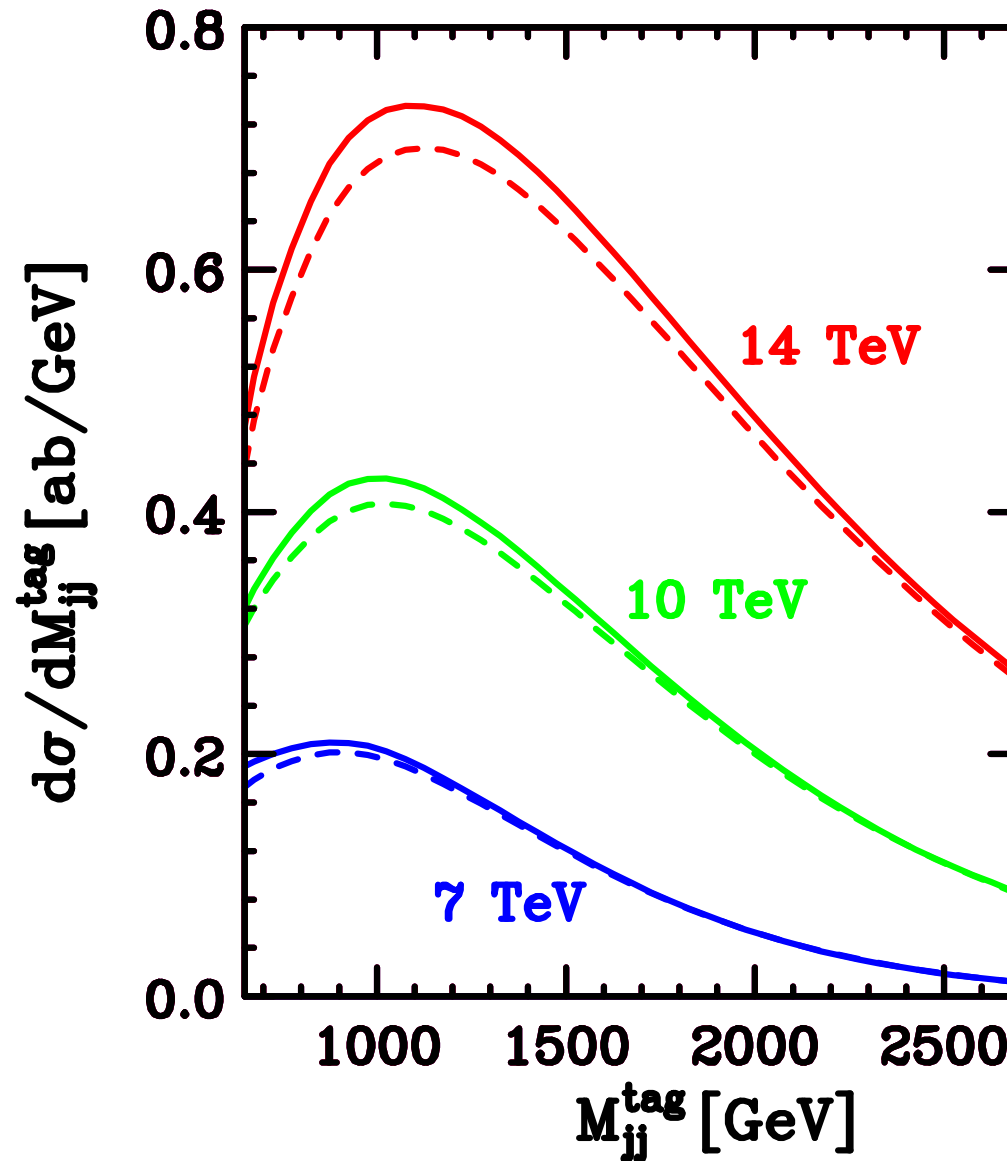
W^+W^+jj distributions: invariant mass of tagging jets

Oleari, Zeppenfeld, B. J. (2009)

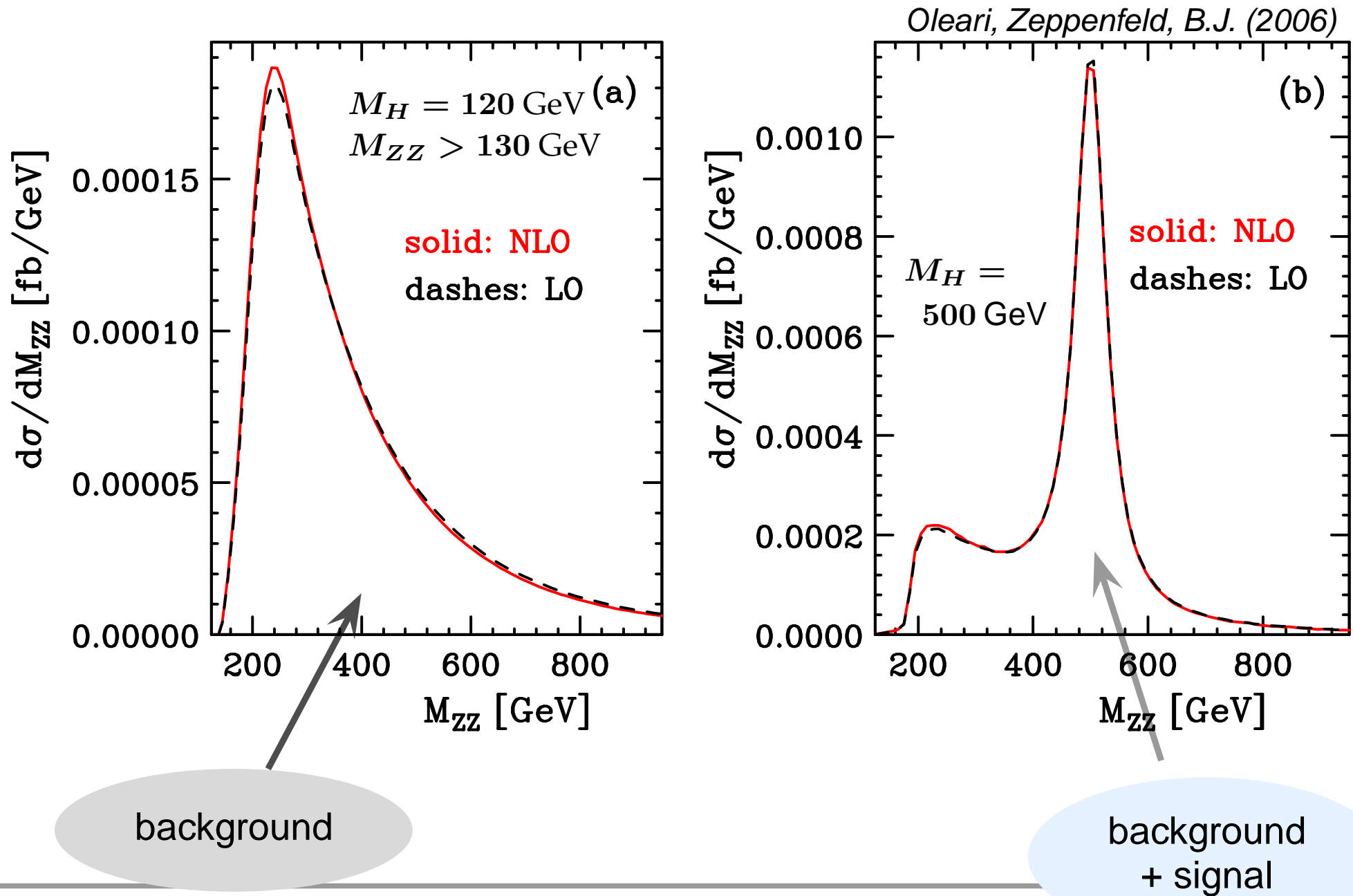


$$\mu = \xi \mu_0$$

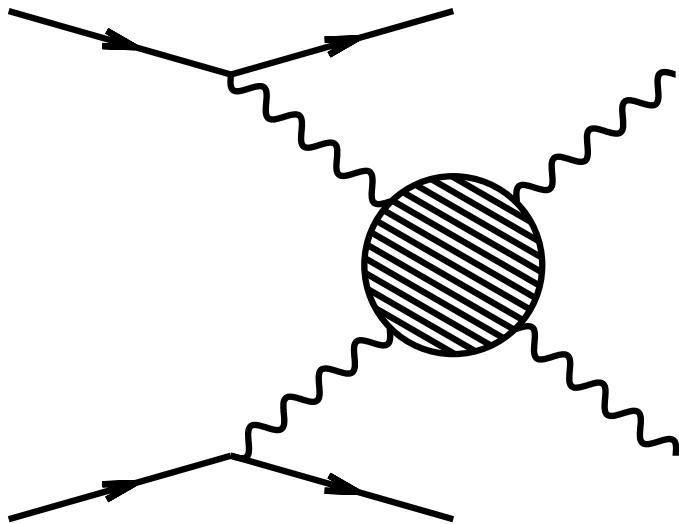
$pp \rightarrow W^+W^+jj$: energy dependence



M_{VV} distribution: $pp \rightarrow \ell^+ \ell^- \ell'^+ \ell'^- jj$



Oleari, Zeppenfeld, B.J. (2006)



VBF processes are extremely
sensitive to new interactions in the
gauge boson sector



can we spot signatures of
non-standard scenarios for
electroweak symmetry breaking?

VV scattering & unitarity

$$W_L^+ W_L^- \rightarrow W_L^+ W_L^-$$

$$\text{with } \epsilon_L^\mu \sim \frac{\sqrt{s}}{M_W}$$

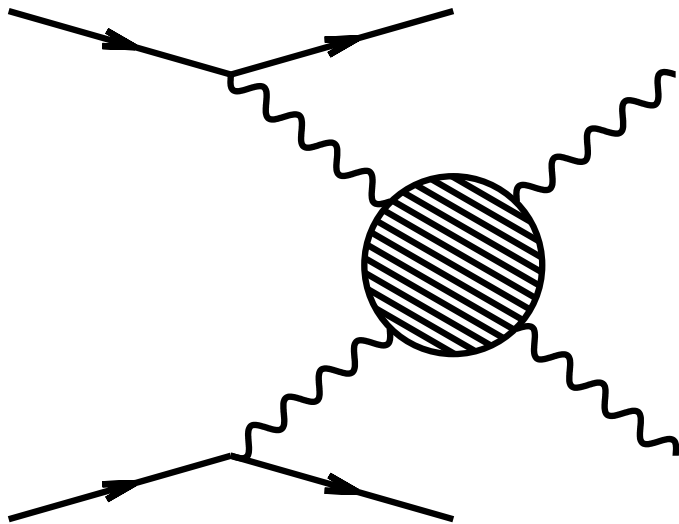
$$\mathcal{M} = \text{[s-channel diagram]} + \text{[t-channel diagram]} + \text{[u-channel diagram]} \sim \frac{s}{M_W^2}$$

growth violates unitarity \rightarrow need:

$$\text{[t-channel diagram]} + \text{[s-channel diagram]}$$

Higgs with $M_H \lesssim 1$ TeV
or new physics at TeV scale

can we distinguish signatures of SM-type Higgs mechanism from other scenarios of EW symmetry breaking?



comprehensive analysis of signal and backgrounds needed

cf. Bagger et al. (1993, 1995)

Englert, Worek, Zeppenfeld, B. J. (2008)

- ➡ minimize backgrounds with respect to signal
- ➡ maximize number of surviving signal events

consider two “prototype” scenarios for the VBF signal:

- ❖ SM with **heavy Higgs boson** ($M_H = 1 \text{ TeV}$, $\Gamma_H = 0.5 \text{ TeV}$)

naive estimate of strongly coupled sector with
scalar, iso-scalar resonance at the TeV scale

- ❖ **Warped Higgsless model** with extra **vector resonances**

$$\left(\begin{array}{l} m_{W_2} = 700 \text{ GeV}, \Gamma = 13.7 \text{ GeV}, \\ m_{Z_2} = 695 \text{ GeV}, \Gamma = 18.7 \text{ GeV}, \\ m_{Z_3} = 718 \text{ GeV}, \Gamma = 6.4 \text{ GeV} \end{array} \right)$$

the Warped Higgsless model

consider gauge boson sector of **Randall-Sundrum scenario with one compactified extra dimension** and AdS₅ metric

$$ds^2 = \frac{R^2}{y^2} \left\{ g_{\mu\nu} dx^\mu dx^\nu - dy^2 \right\}$$

$$R \leq y \leq R'$$

Planck brane

TeV brane

boundary conditions along extra dimension

[Csáki, Grojean, Murayama, Pilo, Terning]

⇒ Kaluza-Klein decomposition of the gauge fields

the Warped Higgsless model

$$W_\mu(x, y) = \sum_k \psi_k^{(W)} W_\mu^{(k)}(x)$$

$$Z_\mu(x, y) = \sum_k \psi_k^{(Z)} Z_\mu^{(k)}(x)$$

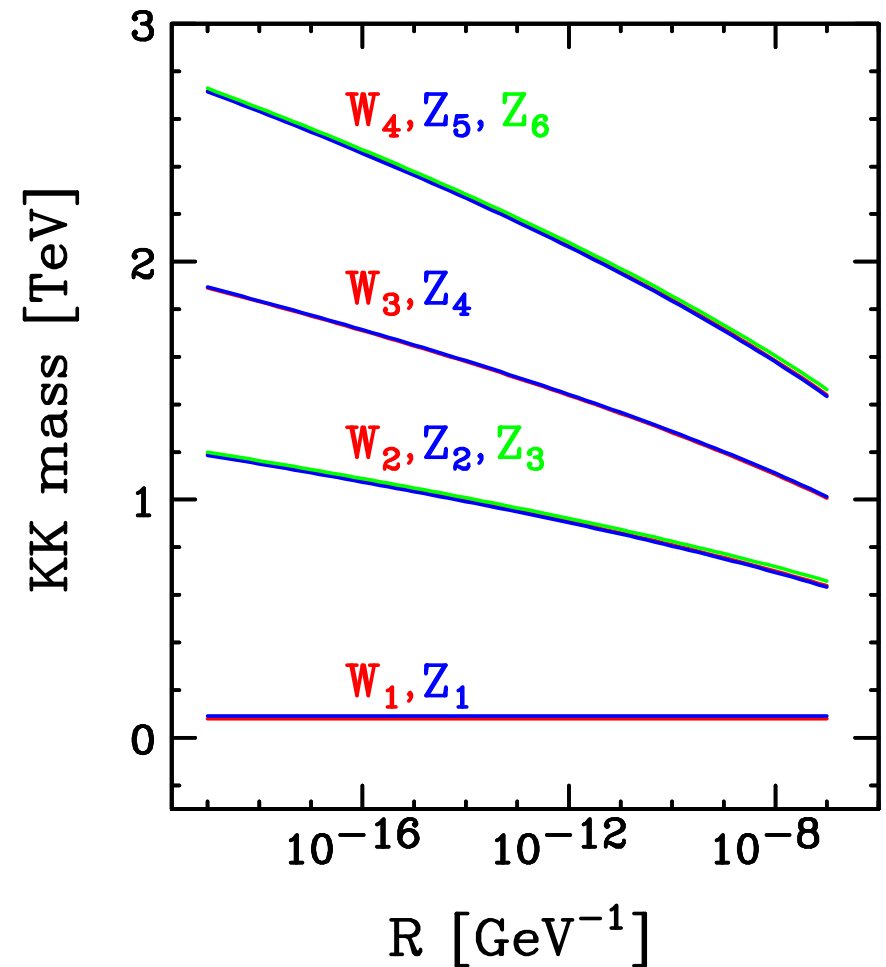
$k = 0$: photon

$k = 1$: SM Z, W^\pm

$k > 1$: KK Z_k, W_k^\pm

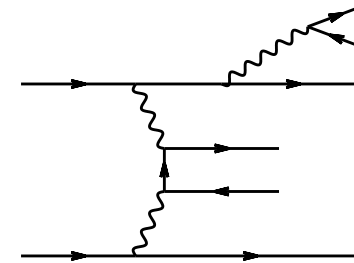
model fully determined by R

Englert, Worek, Zeppenfeld, B. J. (2008)

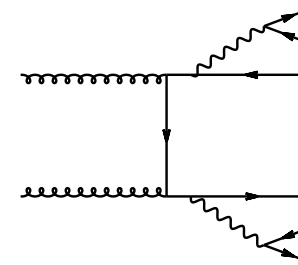


backgrounds to the strongly interacting gauge boson signal
in the **heavy Higgs (HH)** and **Kaluza-Klein (KK)** scenarios:

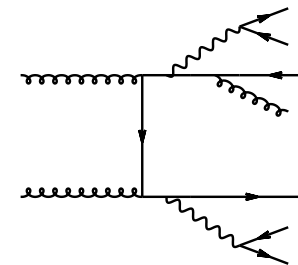
❖ EW $VVjj$ production



❖ QCD $VVjj$ production



❖ $t\bar{t}$ + jets production (with $t \rightarrow Wb$)



inclusive cuts

$$\begin{aligned} p_{Tj}^{\text{tag}} &> 30 \text{ GeV}, |\eta_j| < 4.5, \\ \Delta R_{jj} &> 0.7, \Delta R_{\ell j} > 0.4, \\ p_{T\ell} &> 20 \text{ GeV}, |\eta_\ell| < 2.5, \\ m_{\ell\ell} &> 15 \text{ GeV} \end{aligned}$$

VBF cuts

$$\begin{aligned} \eta_{j,\min}^{\text{tag}} &< \eta_\ell < \eta_{j,\max}^{\text{tag}}, \\ \eta_{j_1}^{\text{tag}} \times \eta_{j_2}^{\text{tag}} &< 0, \\ \Delta\eta_{jj} &> 4, m_{jj} > m_{jj}^{\min} \end{aligned}$$

leptonic cuts
(process-specific)

central jet veto
b-tagging veto

in contrast to backgrounds, **signal** processes feature **energetic leptons** of high p_T and large invariant mass

details of leptonic cuts depend on decay channel

$$\underline{pp \rightarrow W^+ W^- jj:}$$

$$p_{T\ell} > 100 \text{ GeV}$$

$$\Delta p_T(\ell\ell) > 250 \text{ GeV}$$

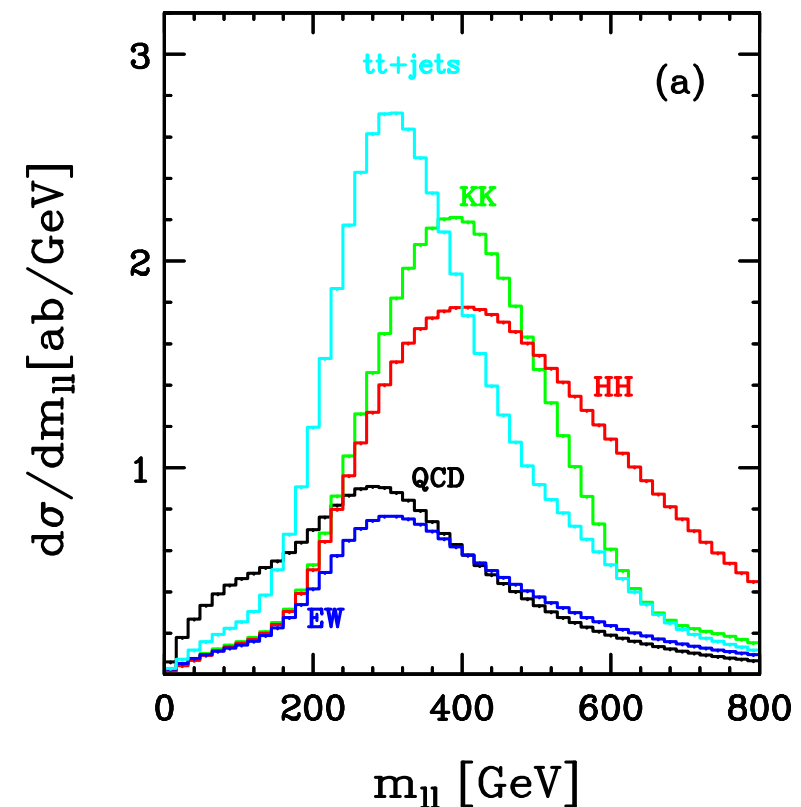
$$\min(m_{\ell j}) > 180 \text{ GeV}$$



inclusive cuts, VBF cuts, CJV & b -veto

$$p_{T\ell} > 100 \text{ GeV}, \min(m_{\ell j}) > 180 \text{ GeV}$$

Englert, Worek, Zeppenfeld, B. J. (2008)



in contrast to backgrounds, **signal** processes feature **energetic leptons** of high p_T and large invariant mass

details of leptonic cuts depend on decay channel

$$\underline{pp \rightarrow W^+ W^- jj:}$$

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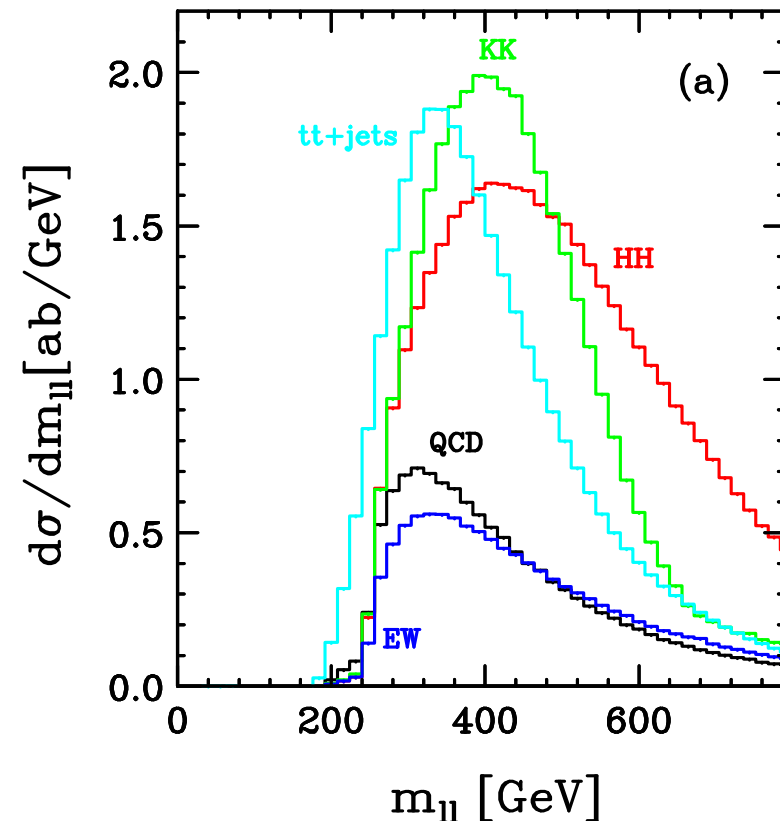
$$m_{\ell\ell} > 200 \text{ GeV}$$

$$\min(m_{\ell j}) > 180 \text{ GeV}$$



... final level of cuts

Englert, Worek, Zeppenfeld, B. J. (2008)



results: scalar resonance

Process	σ_S	σ_B	S/B	S/\sqrt{B}	N_{signal}	$N_{\text{bkgd.}}$
<i>$ZZjj \rightarrow 4\ell jj$</i>	0.048	0.021	2.2	5.7	14	6
<i>$ZZjj \rightarrow 2l2\nu jj$</i>	0.27	0.10	2.7	14.8	81	30
<i>W^+W^-jj</i>	0.51	0.78	0.6	10.0	153	234
$W^\pm Zjj$	0.031	0.386	0.1	0.9	9	116

final level of cuts & integrated luminosity ... 300 fb⁻¹

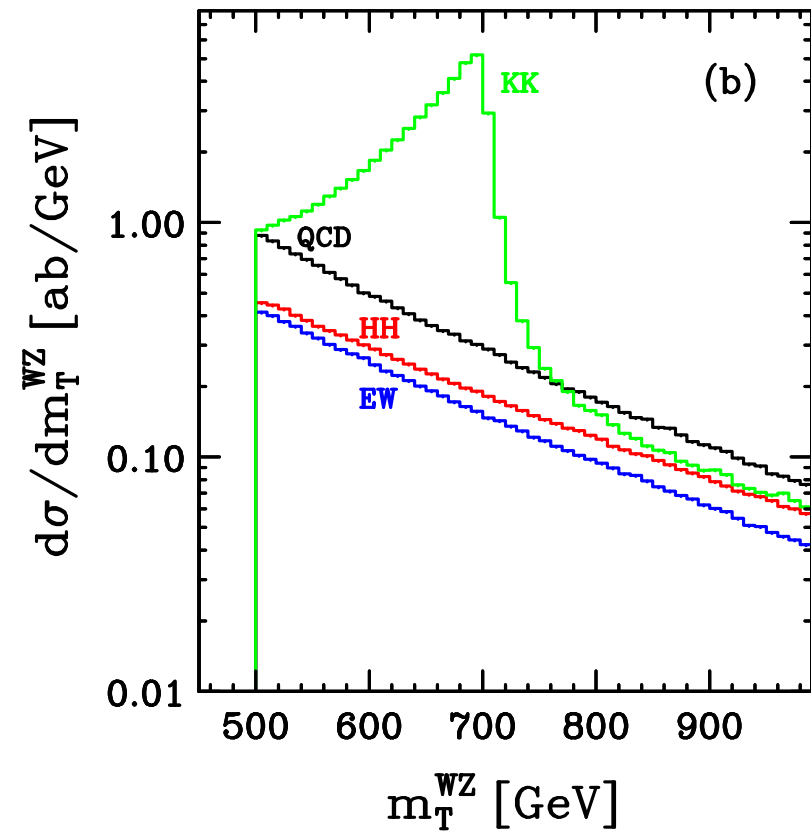
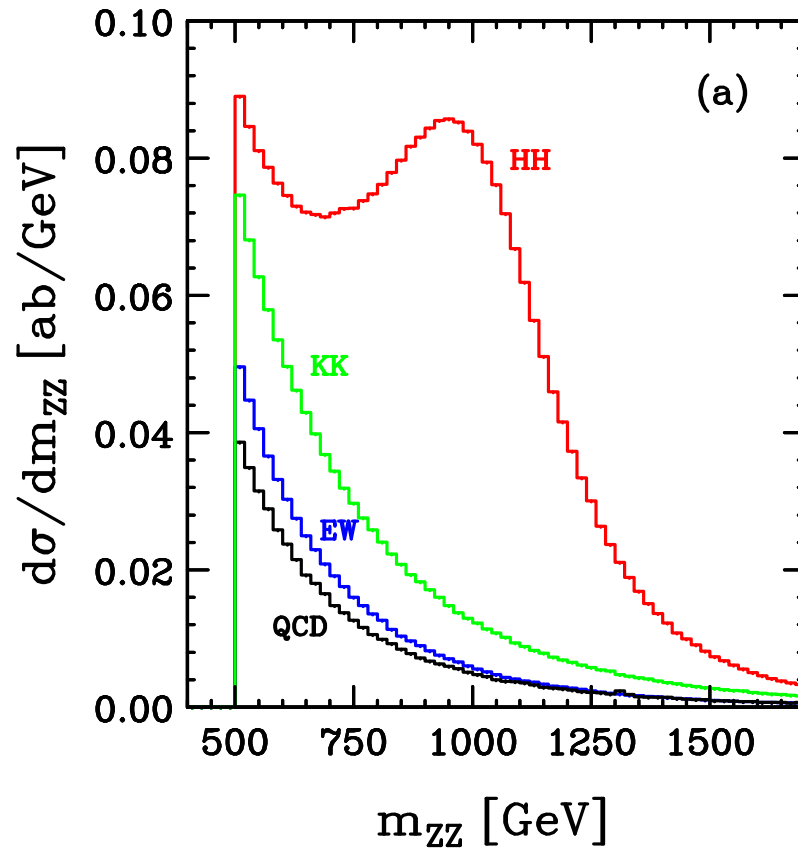
results: vector resonance

Process	σ_S	σ_B	S/B	S/\sqrt{B}	N_{signal}	$N_{\text{bkgd.}}$
$W^\pm Z jj$	0.68	0.39	1.7	18.9	204	117
$W^+ W^- jj$	0.40	0.78	0.5	7.9	120	234
$ZZ jj \rightarrow 4\ell jj$	0.009	0.021	0.4	1.1	3	6
$ZZ jj \rightarrow 2\ell 2\nu jj$	0.05	0.10	0.5	2.7	15	30

final level of cuts & integrated luminosity ... 300 fb⁻¹

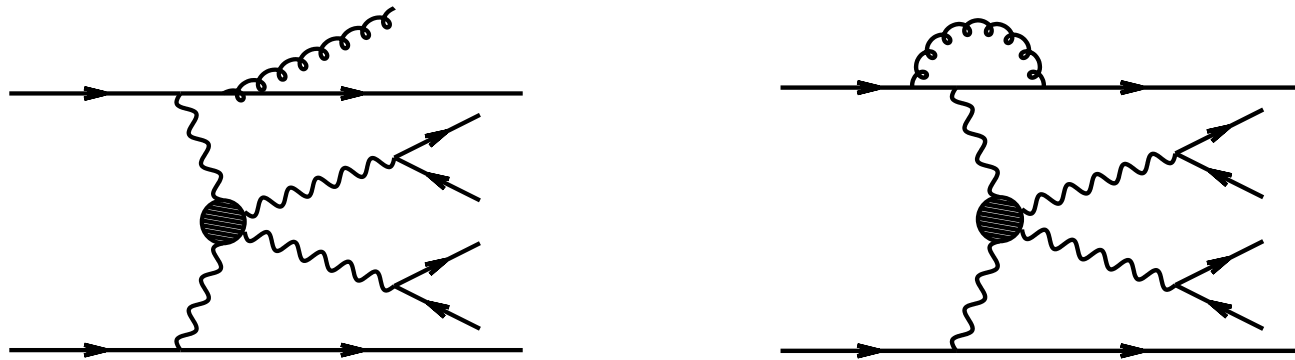
LO results: scalar / vector resonances

Englert, Worek, Zeppenfeld, B.J. (2008)



$ZZjj \rightarrow 4\ell jj$

W^+Zjj



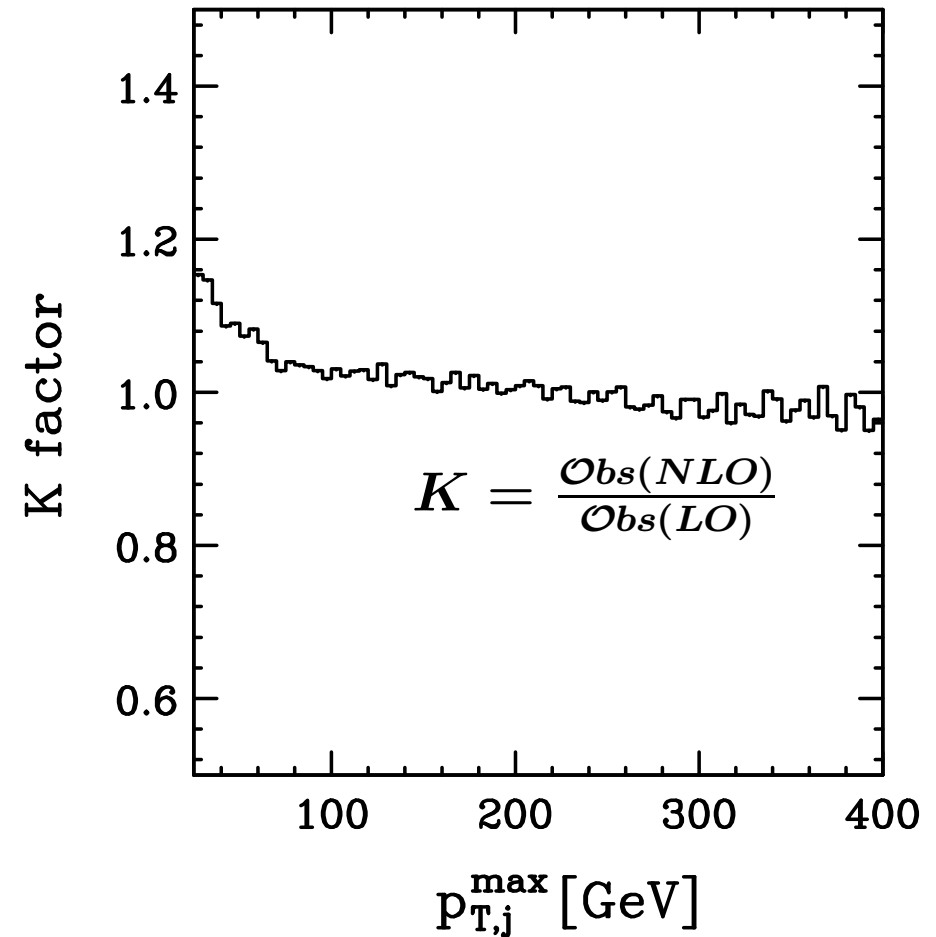
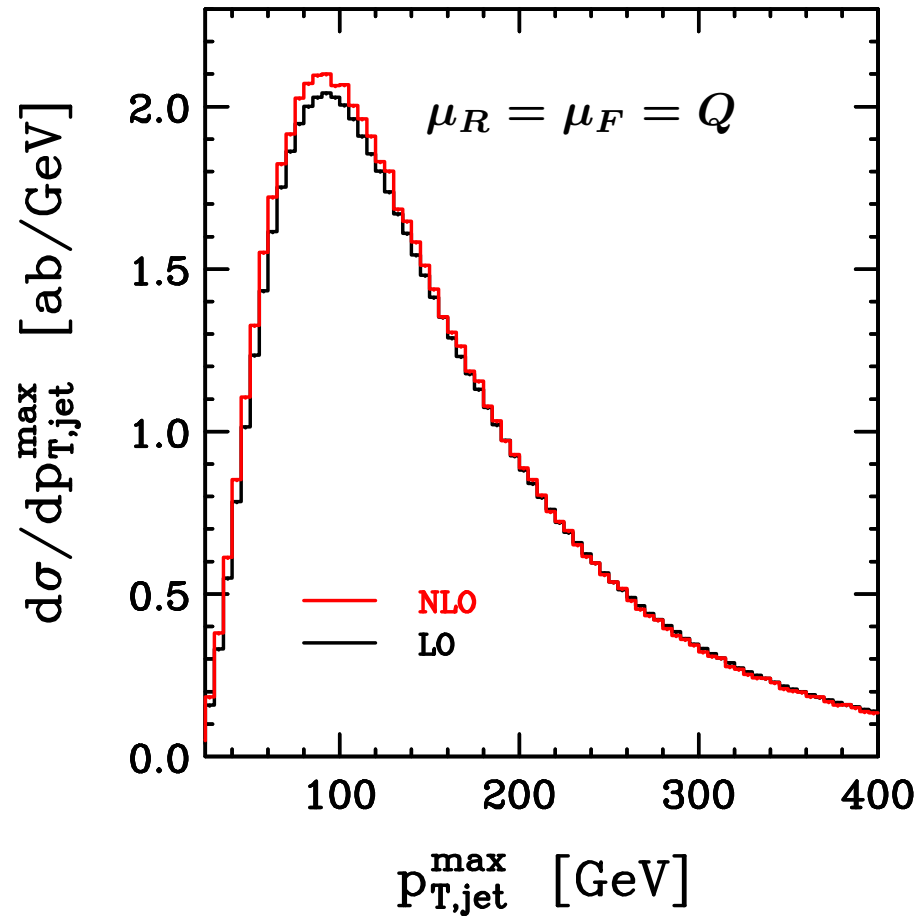
- ❖ compute **real emission** diagrams
- ❖ compute **virtual corrections**

$$2 \operatorname{Re}[\mathcal{M}_V \mathcal{M}_B^*] \propto \frac{\alpha_s(\mu_r)}{2\pi} |\mathcal{M}_B|^2 \times \left[-\frac{2}{\epsilon^2} - \frac{3}{\epsilon} + \text{const.} \right] + 2 \operatorname{Re}[\tilde{\mathcal{M}}_V \mathcal{M}_B^*]$$

- ❖ handle IR divergencies by **dipole subtraction** approach (*Catani, Seymour*)

impact of NLO-QCD corrections

Englert, Zeppenfeld, B. J. (2008)



NLO-QCD corrections always in the few-percent range

❖ explicit calculations revealed that VBF reactions are **perturbatively well-behaved** (moderate NLO QCD and EW corrections, negligible higher order and interference effects)

❖ **backgrounds** are well under control

signatures of new physics in the gauge boson sector should be observable at the LHC



VBF crucial for understanding mechanism of electroweak symmetry breaking

