

Theoretische Teilchenphysik 1

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Exercise sheet 1

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Exercise 1: Gamma matrix representations (P) *(2+2+1+2+2+1 = 10 points)*

We consider the Dirac matrices γ^μ ($\mu = 0, \dots, 3$) in the Weyl representation

$$\gamma_\chi^\mu = \left(\begin{pmatrix} 0 & 1_{2 \times 2} \\ 1_{2 \times 2} & 0 \end{pmatrix}, \begin{pmatrix} 0 & \vec{\sigma} \\ -\vec{\sigma} & 0 \end{pmatrix} \right)$$

as well as in the Dirac representation

$$\gamma_D^\mu = \left(\begin{pmatrix} 1_{2 \times 2} & 0 \\ 0 & -1_{2 \times 2} \end{pmatrix}, \begin{pmatrix} 0 & \vec{\sigma} \\ -\vec{\sigma} & 0 \end{pmatrix} \right),$$

where $1_{2 \times 2}$ is the 2×2 unit matrix and $\vec{\sigma} = (\sigma_1, \sigma_2, \sigma_3)$ with σ_i ($i = 1, 2, 3$) represent the Pauli matrices.

- Show that both representations obey the Clifford algebra $\{\gamma^\mu, \gamma^\nu\} = 2g^{\mu\nu}$ by inserting the explicit representations separately into the anti-commutation relations.
- The Weyl and Dirac representations are connected by a unitary transformation U such that

$$\gamma_\chi^\mu = U^\dagger \gamma_D^\mu U$$

holds. Up to an arbitrary phase, determine the unitary 4×4 matrix U explicitly.

- A fifth Dirac matrix, γ^5 , can be defined by $\gamma_5 := i\gamma^0\gamma^1\gamma^2\gamma^3$. Show that the explicit forms of γ_D^5 and γ_χ^5 for both the Dirac and the Weyl representation, respectively, are given by

$$\gamma_D^5 = \begin{pmatrix} 0 & 1_{2 \times 2} \\ 1_{2 \times 2} & 0 \end{pmatrix}, \quad \gamma_\chi^5 = \begin{pmatrix} -1_{2 \times 2} & 0 \\ 0 & 1_{2 \times 2} \end{pmatrix}.$$

- Show that $\{\gamma^5, \gamma^\mu\} = 0$ holds for both the Dirac and Weyl representation separately by inserting the Dirac matrices explicitly in these representations.

- (e) We now define chirality projectors ω_{\mp} by

$$\omega_{\mp} := \frac{1_{4 \times 4} \mp \gamma^5}{2} ,$$

where ω_- is the left-chiral projector and ω_+ is the right-chiral projector. By using the explicit representation of γ^5 in the Weyl basis only, show that the ω_{\mp} obey the following projector properties:

$$\omega_{\mp}^2 = \omega_{\mp} , \quad \omega_- \omega_+ = \omega_+ \omega_- = 0 .$$

- (f) We now define a bispinor $\Psi := (\psi_L, \psi_R)$ in the Weyl basis, where ψ_L and ψ_R are left- and right-chiral spinors, respectively. Calculate the following bispinors by inserting the explicit representation of γ^5 in the Weyl basis:

$$\Psi_L := \omega_- \Psi , \quad \Psi_R := \omega_+ \Psi .$$

Interpret the result and explain why the Weyl basis is often called chiral basis.

Exercise 2: Dirac algebra in 4 dimensions (P)

(4+6 = 10 points)

Independent of any representation, the Dirac matrices γ^μ and γ^5 in 4 space-time dimensions obey the Clifford algebra $\{\gamma^\mu, \gamma^\nu\} = 2g^{\mu\nu}$ as well as $\{\gamma^5, \gamma^\mu\} = 0$. The γ^μ can always be chosen to be unitary so that $(\gamma^\mu)^\dagger = (\gamma^\mu)^{-1}$ holds.

- (a) Prove the following Dirac algebra relations by using the unitarity and anti-commutation relations of the Dirac matrices, only (i.e. do not use the explicit representations of the Dirac matrices from Exercise 1):

$$\begin{aligned} (\gamma^\mu)^\dagger &= \gamma^0 \gamma^\mu \gamma^0 , \\ \gamma^5 \gamma^5 &= 1_{4 \times 4} , \\ \gamma_\mu \gamma^\mu &= 4 \cdot 1_{4 \times 4} , \\ \gamma_\mu \gamma^\alpha \gamma^\mu &= -2\gamma^\alpha , \\ \gamma_\mu \gamma^\alpha \gamma^\beta \gamma^\mu &= 4g^{\alpha\beta} . \end{aligned}$$

- (b) Prove the following trace identities by using the anti-commutation relations of the Dirac matrices and the general properties of a matrix trace, only:

$$\begin{aligned} \text{tr}(\gamma^\mu) &= 0 , \\ \text{tr}(\gamma^5) &= 0 , \\ \text{tr}(\gamma^\mu \gamma^\nu \gamma^5) &= 0 , \\ \text{tr}(\gamma^\mu \gamma^\nu \gamma^\rho) &= 0 , \\ \text{tr}(\gamma^\mu \gamma^\nu) &= 4g^{\mu\nu} , \\ \text{tr}(\gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma) &= 4(g^{\mu\nu} g^{\rho\sigma} - g^{\mu\rho} g^{\nu\sigma} + g^{\mu\sigma} g^{\nu\rho}) . \end{aligned}$$