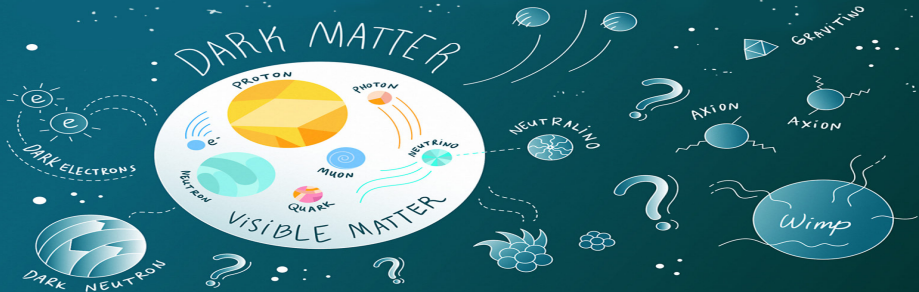


DM Collider Searches

Introduction

Jonas Müller | 8.7.2019

INSTITUT FÜR THEORETISCHE PHYSIK (ITP)



- 1 Motivation/Expectation
- 2 Approaches
- 3 Top-Down
- 4 Bottom-Up

1 Motivation/Expectation

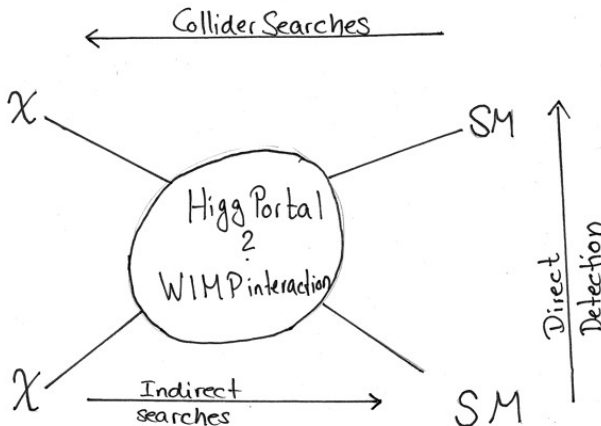
2 Approaches

3 Top-Down

4 Bottom-Up

Motivation/Expectation

- If DM is a fundamental new particle *and* interacts with SM particles
 - Production in collider is possible
 - Decay chains into DM particles are possible



What can we expect?

- Suppose we have
 - $\dot{\mathcal{L}} = 5 \cdot 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$
 - $m_\chi = 100 \text{ GeV}$ and $\sigma_{LHC} = G_F^2 m_\chi^2$
 - Isotropic production in an distance $R = 10 \text{ cm}$
 - $\rho_{DM} = 0.3 \text{ GeV/cm}^3$ and $v_{DM} = 220 \text{ km/s}$

$$\Phi_{LHC} = 4\pi R^2 \dot{\mathcal{L}} \sigma_{LHC} \sim 10^4 \text{ cm}^2/\text{s}$$

$$\Phi_{halo} = \frac{\rho_{DM}}{m_\chi} \pi R^2 v_{DM} \sim 10^9 \text{ cm}^2/\text{s}$$

⇒ The expected flux in the LHC collider is much smaller than the halo flux!

1 Motivation/Expectation

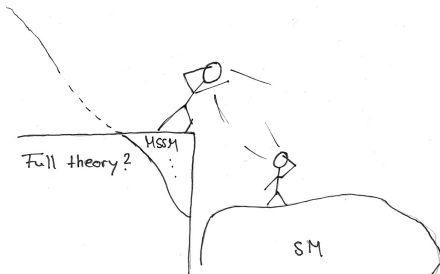
2 **Approaches**

3 Top-Down

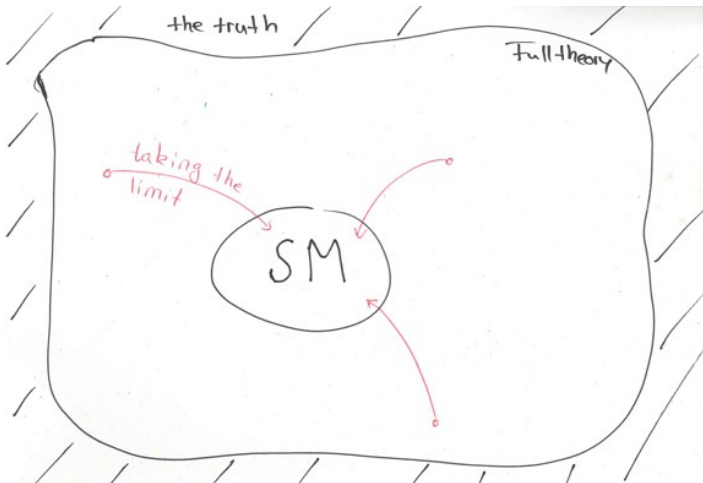
4 Bottom-Up

Extend the Standard Model

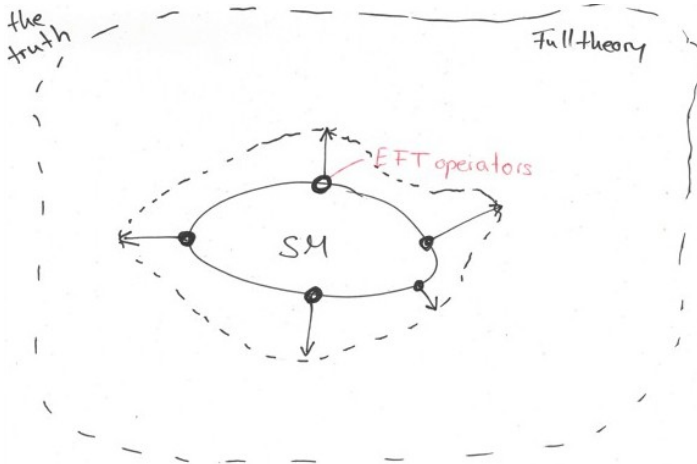
- top-down approach
- bottom-up approach



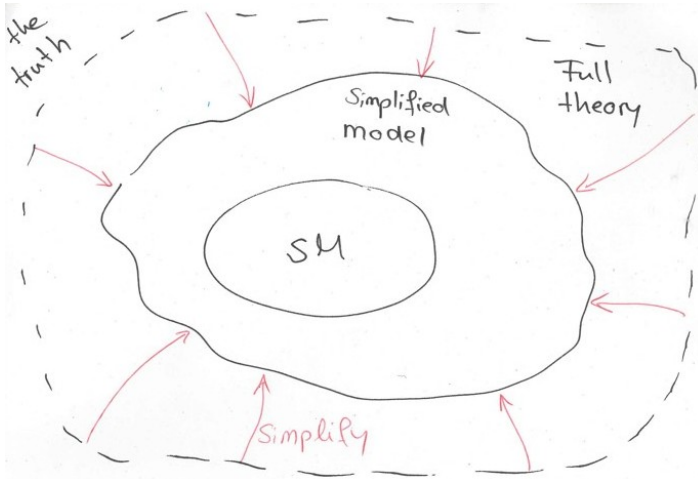
top-down Approach



bottom-up Approach - EFT



bottom-up Approach - Simplified Models



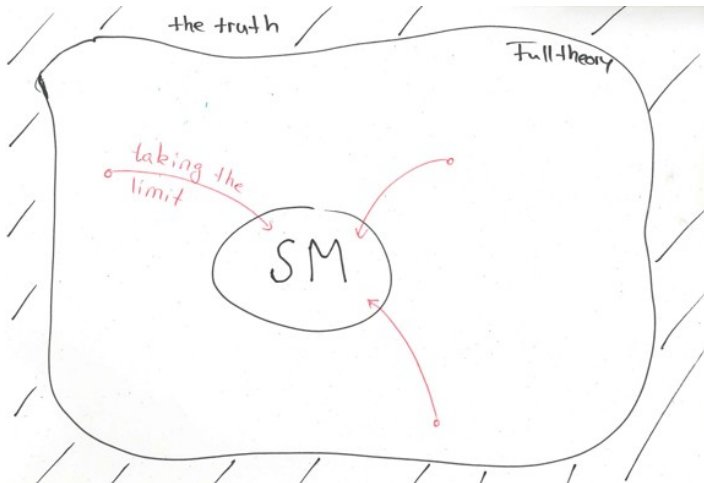
1 Motivation/Expectation

2 Approaches

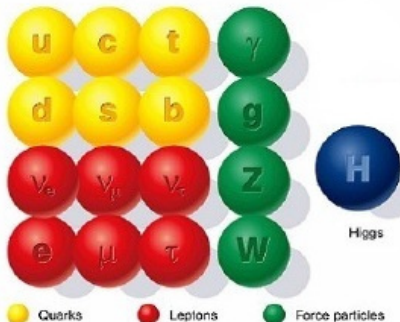
3 Top-Down

4 Bottom-Up

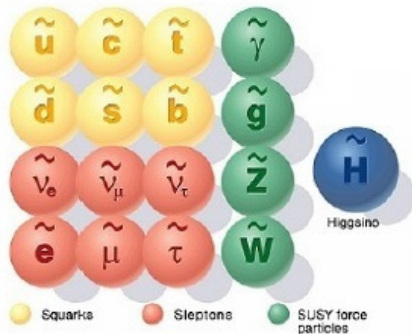
Top-Down Approach - SUSY



SUPERSYMMETRY



Standard particles



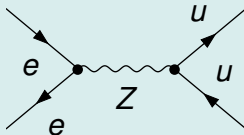
SUSY particles

Neutralino is a mixing state of the gauge and Higgs boson superpartners

$$\begin{pmatrix} \tilde{\chi}_1^0 \\ \tilde{\chi}_2^0 \\ \tilde{\chi}_3^0 \\ \tilde{\chi}_4^0 \end{pmatrix} = N \cdot \begin{pmatrix} \tilde{W}^0 \\ \tilde{B}^0 \\ \tilde{H}_a^0 \\ \tilde{H}_b^0 \end{pmatrix}$$

the lightest neutralino $\tilde{\chi}_1^0$ is the stable DM candidate!

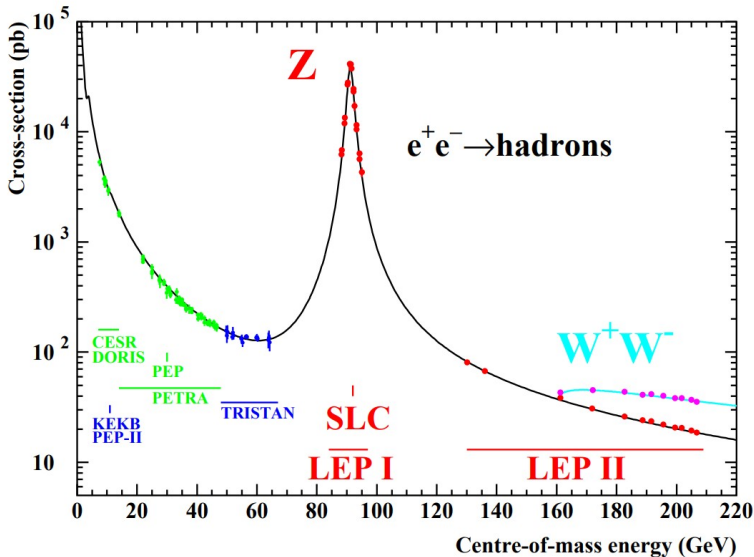
Z-Pole Measurement



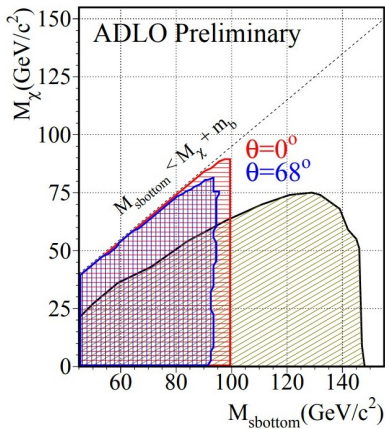
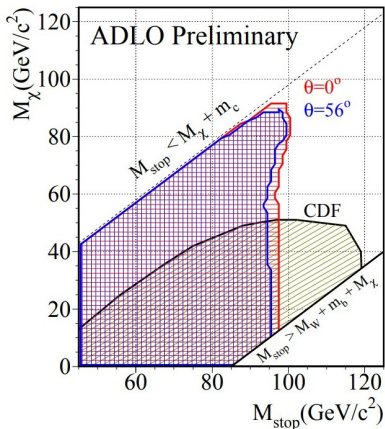
- $\sqrt{s} \sim m_Z$
- Measurement of the total width $\Delta\Gamma_Z$
- Looking for events with high missing transverse momentum/energy
- Allows for an upper bound of $\Delta\Gamma_{inv}$

⇒ Bound on $\Delta\Gamma_{inv}$ translates on an lower bound on $m_{\tilde{\chi}}$!

$$\tilde{f} \rightarrow f + \tilde{\chi}_1^0$$

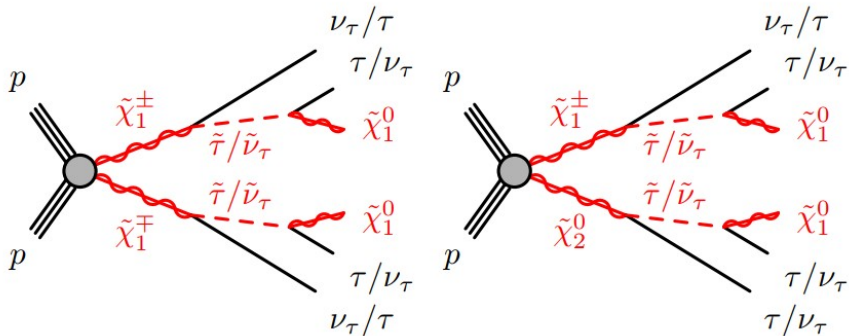


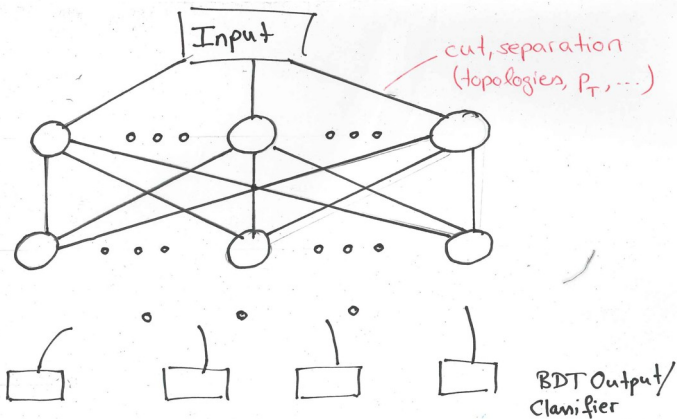
$$\tilde{f} \rightarrow f + \tilde{\chi}_1^0$$



- More Energy, but more problems!
- General Idea of the searches:
 - Take your favorite model (e.g. MSSM, NMSSM...)
 - Calculate the production crosssection, branching ratios...
 - Generate an event simulation based on the model (and for the detector!)
 - Generate an event simulation with SM only (background)
 - Apply *smart* parameter cuts to maximise S/B
 - Hope for a detectable signal!
- Typical DM signatures:
 - High \cancel{E}_T
 - Mono-X signatures
- The *art* of the parameter cuts
 - Boosted Decision Trees (BDTs)
 - Deep Learning (DL)

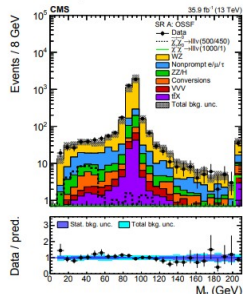
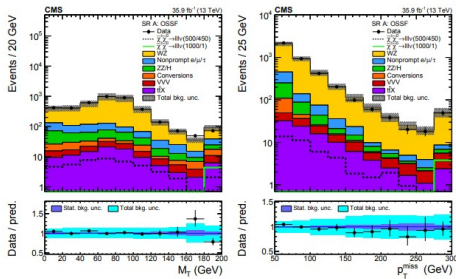
Example Search Process

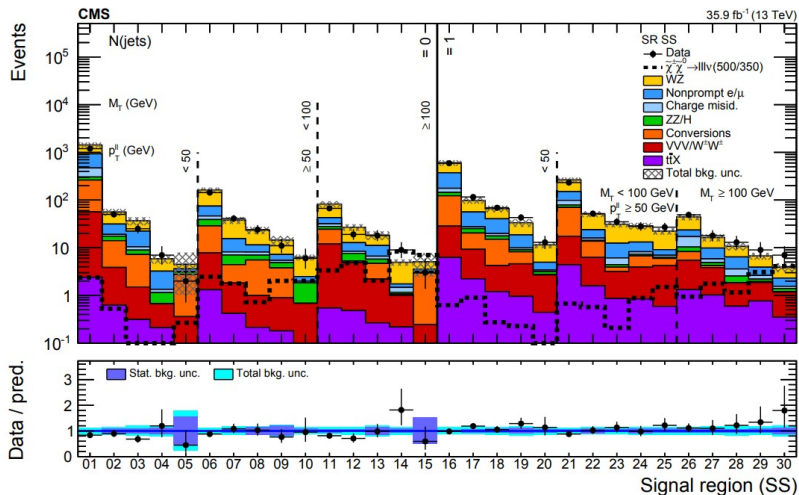




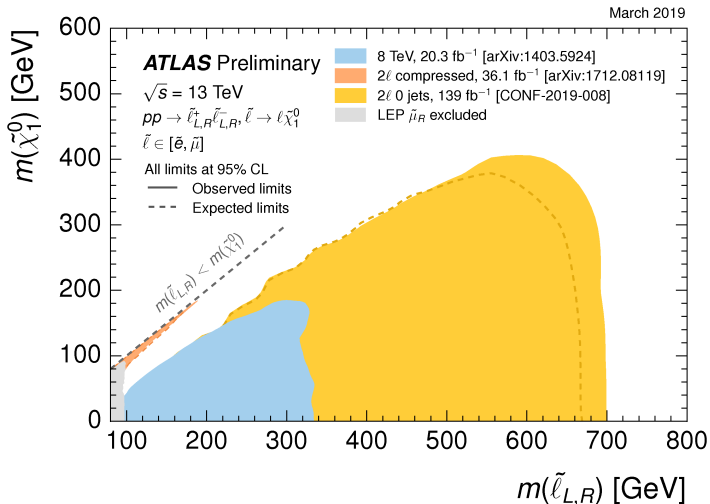
What is used to select events?

- Event Topologies (0-jet, 1-jet, ...)
- p_T cuts
- Invariant mass distributions
- Missing transverse energy \cancel{E}_T





Exclusion Limits for lightest neutralino and slepton masses



Exclusion Summary-ATLAS

ATLAS SUSY Searches* - 95% CL Lower Limits

March 2019

ATLAS Preliminary

$\sqrt{s} = 13 \text{ TeV}$

Model	Signature	$\int \mathcal{L} dt [\text{fb}^{-1}]$	Mass limit	Reference		
Inclusive Searches	$\tilde{g}\tilde{g}, \tilde{q}\tilde{q} \rightarrow \text{gluino}$	0 e, μ mono-jet	2-6 jets 1-3 jets	E_{T}^{miss} 36.1 E_{T}^{miss} 36.1	$m(\tilde{g}) > 100 \text{ GeV}$ $m(\tilde{q}) > m(\tilde{g}) - 5 \text{ GeV}$	1712.02332 1711.03931
	$\tilde{g}\tilde{g}, \tilde{q}\tilde{q} \rightarrow \text{gluino}$	0 e, μ	2-6 jets	E_{T}^{miss} 36.1	\tilde{g}	1712.02332
	$\tilde{g}\tilde{g}, \tilde{q}\tilde{q} \rightarrow \text{gluino}$	3 e, μ $\text{no } \tau$	4 jets	E_{T}^{miss} 36.1	\tilde{g}	1712.02332
	$\tilde{g}\tilde{g}, \tilde{q}\tilde{q} \rightarrow \text{gluino}$	3 e, μ 2 jets	4 jets	E_{T}^{miss} 36.1	\tilde{g}	1712.02332
	$\tilde{g}\tilde{g}, \tilde{q}\tilde{q} \rightarrow \text{gluino}$	0 e, μ 3 e, μ	7-11 jets 4 jets	E_{T}^{miss} 36.1 E_{T}^{miss} 36.1	\tilde{g}	1712.02332
	$\tilde{g}\tilde{g}, \tilde{q}\tilde{q} \rightarrow \text{gluino}$	0-1 e, μ 3 e, μ	3 jets 4 jets	E_{T}^{miss} 79.8 E_{T}^{miss} 36.1	\tilde{g}	1706.60731 1706.60731
3 rd gen. squarks direct production	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{t}_1^* \tilde{b}_1^*$	Multiple	36.1	\tilde{b}_1	1708.09286, 1711.03931	
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{t}_1^* \tilde{b}_1^*$	Multiple	36.1	\tilde{b}_1	1708.09286	
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{t}_1^* \tilde{b}_1^*$	Multiple	139	\tilde{b}_1	1706.60731	
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{t}_1^* \tilde{t}_1^*$	0-2 e, μ	0-2 jets+1-2 b	E_{T}^{miss} 36.1	\tilde{t}_1	1508.08616, 1709.04183, 1711.11520
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{t}_1^* \tilde{t}_1^*$	Multiple	36.1	\tilde{t}_1	1709.04183, 1711.11520	
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{t}_1^* \tilde{t}_1^*$	1 $\tau + 1 e, \mu, \tau$	2 jets+1 b	E_{T}^{miss} 36.1	\tilde{t}_1	1803.10178
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{t}_1^* \tilde{t}_1^*$	0 e, μ	2 c	E_{T}^{miss} 36.1	\tilde{t}_1	1805.01649
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{t}_1^* \tilde{t}_1^*$	0 e, μ	mono-jet	E_{T}^{miss} 36.1	\tilde{t}_1	1805.01649
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{t}_1^* \tilde{t}_1^*$	1-2 e, μ	4 b	E_{T}^{miss} 36.1	\tilde{t}_1	1711.03931
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{t}_1^* \tilde{t}_1^*$	1-2 e, μ	4 b	E_{T}^{miss} 36.1	\tilde{t}_1	1706.03986
EW direct	$\tilde{t}_1^*\tilde{t}_1^*$ via WZ	2-3 e, μ	36.1	\tilde{t}_1^*	1403.5294, 1806.02250	
	$\tilde{t}_1^*\tilde{t}_1^*$ via WZ	$\text{no } \tau$	E_{T}^{miss} 36.1	\tilde{t}_1^*	1712.08119	
	$\tilde{t}_1^*\tilde{t}_1^*$ via WW	2 e, μ	2 b	E_{T}^{miss} 139	\tilde{t}_1^*	ATLAS CONF-2019-008
	$\tilde{t}_1^*\tilde{t}_1^*$ via WW	0-1 e, μ	2 b	E_{T}^{miss} 36.1	\tilde{t}_1^*	1812.09402
	$\tilde{t}_1^*\tilde{t}_1^*$ via WW	2 e, μ	2 τ	E_{T}^{miss} 139	\tilde{t}_1^*	ATLAS CONF-2019-008
	$\tilde{t}_1^*\tilde{t}_1^*$ via WW	2 e, μ	2 τ	E_{T}^{miss} 36.1	\tilde{t}_1^*	1708.07875
	$\tilde{t}_1^*\tilde{t}_1^*$ via WW	2 e, μ	0 jets	E_{T}^{miss} 139	\tilde{t}_1^*	1708.07875
	$\tilde{t}_1^*\tilde{t}_1^*$ via WW	2 e, μ	≥ 1	E_{T}^{miss} 36.1	\tilde{t}_1^*	ATLAS CONF-2019-008
	$\tilde{t}_1^*\tilde{t}_1^*$ via WW	0 e, μ	≥ 3 b	E_{T}^{miss} 36.1	\tilde{t}_1^*	1712.08119
	$\tilde{t}_1^*\tilde{t}_1^*$ via WW	4 e, μ	0 jets	E_{T}^{miss} 36.1	\tilde{t}_1^*	1806.04030
Long-lived particles	Direct $\tilde{t}_1^*\tilde{t}_1^*$ prod., long-lived \tilde{t}_1^*	Disapp. trk	1 jet	E_{T}^{miss} 36.1	\tilde{t}_1^*	1712.02118
	Stable \tilde{g} R-hadron	Multiple	36.1	\tilde{g}	Pure Wino	
RPV	LFV $\tilde{g}\tilde{g} \rightarrow \tilde{g} + X, \tilde{g} \rightarrow \text{gluino} + \text{jet}$	$\text{no } \tau, \mu, \tau$	3-2	E_{T}^{miss} 36.1	\tilde{g}	1607.08079
	$\tilde{t}_1^*\tilde{t}_1^* \rightarrow W\tilde{t}_1^* \tilde{t}_1^*$	4 e, μ	0 jets	E_{T}^{miss} 36.1	\tilde{t}_1^*	1804.03902
	$\tilde{t}_1^*\tilde{t}_1^* \rightarrow W\tilde{t}_1^* \tilde{t}_1^*$	4-6 (large-N) jets	36.1	E_{T}^{miss} 36.1	\tilde{t}_1^*	1804.03902
	$\tilde{t}_1^*\tilde{t}_1^* \rightarrow W\tilde{t}_1^* \tilde{t}_1^*$	Multiple	36.1	E_{T}^{miss} 36.1	\tilde{t}_1^*	1710.04901, 1806.04095
	$\tilde{t}_1^*\tilde{t}_1^* \rightarrow W\tilde{t}_1^* \tilde{t}_1^*$	Multiple	36.1	E_{T}^{miss} 36.1	\tilde{t}_1^*	1710.04901, 1806.04095
	$\tilde{t}_1^*\tilde{t}_1^* \rightarrow W\tilde{t}_1^* \tilde{t}_1^*$	Multiple	36.1	E_{T}^{miss} 36.1	\tilde{t}_1^*	1710.04901, 1806.04095
	$\tilde{t}_1^*\tilde{t}_1^* \rightarrow W\tilde{t}_1^* \tilde{t}_1^*$	Multiple	36.1	E_{T}^{miss} 36.1	\tilde{t}_1^*	1710.04901, 1806.04095
	$\tilde{t}_1^*\tilde{t}_1^* \rightarrow W\tilde{t}_1^* \tilde{t}_1^*$	Multiple	36.1	E_{T}^{miss} 36.1	\tilde{t}_1^*	1710.04901, 1806.04095
	$\tilde{t}_1^*\tilde{t}_1^* \rightarrow W\tilde{t}_1^* \tilde{t}_1^*$	Multiple	36.1	E_{T}^{miss} 36.1	\tilde{t}_1^*	1710.04901, 1806.04095
	$\tilde{t}_1^*\tilde{t}_1^* \rightarrow W\tilde{t}_1^* \tilde{t}_1^*$	Multiple	36.1	E_{T}^{miss} 36.1	\tilde{t}_1^*	1710.04901, 1806.04095
$\tilde{t}_1^*\tilde{t}_1^* \rightarrow W\tilde{t}_1^* \tilde{t}_1^*$	Multiple	36.1	E_{T}^{miss} 36.1	\tilde{t}_1^*	1710.04901, 1806.04095	

*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, e.g. refs. for the assumptions made.

10⁻¹

1

Mass scale [TeV]

Exclusion Summary-ATLAS

ATLAS SUSY Searches* - 95% CL Lower Limits

March 2019

ATLAS Preliminary

$\sqrt{s} = 13 \text{ TeV}$

Model	Signature	$\mathcal{L} \cdot \text{Br} \cdot (\text{fb}^{-1})$	Mass limit	Reference					
Inclusive Searches	$\tilde{q}\tilde{q} \rightarrow q\tilde{q}\tilde{q}^0$	0 e, μ mono-jet	2-6 jets 1-3 jets	E_{T}^{miss} E_{T}^{miss}	36.1 36.1	\tilde{g} [DW, Six Degim] 0.43 \tilde{g} [Tks, Six Degim] 0.71 \tilde{g} 0.98 \tilde{g} 1.55	$m(\tilde{L}) = 100 \text{ GeV}$ $m(\tilde{g}) = m(\tilde{t}) = 5 \text{ GeV}$	1712.02302 1712.02301	
	$\tilde{g}\tilde{g} \rightarrow q\tilde{q}\tilde{q}\tilde{q}^0$	0 e, μ 2-6 jets	E_{T}^{miss}	36.1	\tilde{g} \tilde{g} Forbiddn 0.95-1.6 \tilde{g} 2.0	$m(\tilde{L}) = 200 \text{ GeV}$ $m(\tilde{L}) = 900 \text{ GeV}$	1712.02302 1712.02302		
	$\tilde{g}\tilde{g} \rightarrow q\tilde{q}\tilde{q}\tilde{q}^0$	3 e, μ $e\mu, \mu\mu$	4 jets 2 jets	E_{T}^{miss}	36.1 36.1	\tilde{g} \tilde{g} 1.2 \tilde{g} 1.85	$m(\tilde{L}) = 800 \text{ GeV}$ $m(\tilde{g}) = m(\tilde{t}) = 50 \text{ GeV}$	1706.03731 1805.11381	
	$\tilde{g}\tilde{g} \rightarrow q\tilde{q}\tilde{W}\tilde{Z}\tilde{L}^0$	0 e, μ 4 jets	7-11 jets 4 jets	E_{T}^{miss}	36.1 36.1	\tilde{g} \tilde{g} 0.96 \tilde{g} 1.8	$m(\tilde{L}) = 400 \text{ GeV}$ $m(\tilde{g}) = m(\tilde{t}) = 200 \text{ GeV}$	1708.02794 1706.03731	
	$\tilde{g}\tilde{g} \rightarrow t\tilde{t}\tilde{L}^0$	0-1 e, μ 3 jets	4 jets	E_{T}^{miss}	79.9 36.1	\tilde{g} \tilde{g} 1.25 \tilde{g} 2.25	$m(\tilde{L}) = 200 \text{ GeV}$ $m(\tilde{g}) = m(\tilde{t}) = 350 \text{ GeV}$	ATLAS-COBF-2018-041 1706.03731	
	1 st gen. squarks direct production	$\tilde{h}_1 \tilde{h}_1, \tilde{h}_1 \rightarrow \tilde{h}_1^0 \tilde{h}_1^0$	Multiple Multiple	36.1 36.1	E_{T}^{miss} E_{T}^{miss}	36.1 36.1	Forbiddn 0.9 Forbiddn 0.58-0.82 Forbiddn 0.7	$m(\tilde{L}) = 300 \text{ GeV}, \text{BR}(\tilde{h}_1^0 \rightarrow \tilde{t}\tilde{t}) = 1$ $m(\tilde{L}) = 200 \text{ GeV}, \text{BR}(\tilde{h}_1^0 \rightarrow \text{BR}(\tilde{L})^0) = 0.5$ $m(\tilde{L}) = 200 \text{ GeV}, m(\tilde{L}) = 300 \text{ GeV}, \text{BR}(\tilde{L})^0 = 1$	1708.09206, 1711.03361 1708.09206 1706.03731
		$\tilde{h}_1 \tilde{h}_1, \tilde{h}_1 \rightarrow \tilde{h}_1^0 \tilde{h}_1^0 \rightarrow \tilde{h}_1^0 \tilde{h}_1^0$	0 e, μ 6 b	E_{T}^{miss}	139	\tilde{h}_1 \tilde{h}_1 Forbiddn 0.23-0.48 \tilde{h}_1 0.23-1.35	$m(\tilde{L}) = 130 \text{ GeV}, m(\tilde{L}) = 100 \text{ GeV}$ $\text{BR}(\tilde{L}^0 \rightarrow \tilde{L}^0 \tilde{L}^0) = 100 \text{ GeV}, m(\tilde{L}^0) = 0 \text{ GeV}$	SUSY-2018-31 SUSY-2018-31	
		$\tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow W\tilde{b}\tilde{L}^0$ or $\tilde{t}_1^0 \tilde{t}_1^0$	0-2 e, μ Multiple	0-2 jets+1-2 b Multiple	E_{T}^{miss}	36.1 36.1	\tilde{t}_1 \tilde{t}_1 0.48-0.84 \tilde{t}_1 1.0	$m(\tilde{L}) = 1 \text{ GeV}$ $m(\tilde{L}) = 150 \text{ GeV}, m(\tilde{L}) = m(\tilde{L}) = 8 \text{ GeV}, \tilde{L} = \tilde{L}$	1506.06616, 1709.04183, 1711.11520
		$\tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow \tilde{t}_1 \tilde{b}\tilde{L}^0, \tilde{t}_1 \rightarrow \tilde{t}_1 \tilde{L}^0$	1 + 1 e, μ, τ 2 jets+1 b	E_{T}^{miss}	36.1	\tilde{t}_1 \tilde{t}_1 0.46 \tilde{t}_1 0.85	$m(\tilde{L}) = 800 \text{ GeV}$ $m(\tilde{L}) = 0 \text{ GeV}$	1803.10178 1805.01649	
		$\tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow \tilde{t}_1 \tilde{L}^0, \tilde{t}_1 \rightarrow \tilde{t}_1 \tilde{L}^0$	0 e, μ 2 c	E_{T}^{miss}	36.1	\tilde{t}_1 \tilde{t}_1 0.42 \tilde{t}_1 0.85	$m(\tilde{L}) = m(\tilde{L}) = 50 \text{ GeV}$ $m(\tilde{L}) = m(\tilde{L}) = 5 \text{ GeV}$	1805.01649 1711.03361	
$\tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow \tilde{t}_1 \tilde{L}^0 + \tilde{b}$		0 e, μ mono-jet	E_{T}^{miss}	36.1	\tilde{t}_1 \tilde{t}_1 0.32-0.88	$m(\tilde{L}) = 0 \text{ GeV}, m(\tilde{L}) = m(\tilde{L}) = 180 \text{ GeV}$	1706.03966		
$\tilde{t}_1^0 \tilde{t}_1^0$ via WZ		2-3 e, μ $e\mu, \mu\mu$	2-3 jets ≥ 1	E_{T}^{miss} E_{T}^{miss}	36.1 36.1	$\tilde{t}_1^0 \tilde{t}_1^0$ $\tilde{t}_1^0 \tilde{t}_1^0$ 0.17 $\tilde{t}_1^0 \tilde{t}_1^0$ 0.42 $\tilde{t}_1^0 \tilde{t}_1^0$ 0.68 $\tilde{t}_1^0 \tilde{t}_1^0$ 1.0 $\tilde{t}_1^0 \tilde{t}_1^0$ 0.76 $\tilde{t}_1^0 \tilde{t}_1^0$ 0.22 $\tilde{t}_1^0 \tilde{t}_1^0$ 0.7	$m(\tilde{L}) = 0$ $m(\tilde{L}) = m(\tilde{L}) = 10 \text{ GeV}$	1403.5294, 1806.02293 1712.02119	
$\tilde{t}_1^0 \tilde{t}_1^0$ via WW		2 e, μ 2 jets	E_{T}^{miss}	139	$\tilde{t}_1^0 \tilde{t}_1^0$	$m(\tilde{L}) = 0$	ATLAS-COBF-2019-008		
$\tilde{t}_1^0 \tilde{t}_1^0$ via $\tilde{L}\tilde{L}^0$		0-1 e, μ 2 jets	E_{T}^{miss}	36.1	$\tilde{t}_1^0 \tilde{t}_1^0$	$m(\tilde{L}) = 0$	ATLAS-COBF-2019-008		
$\tilde{t}_1^0 \tilde{t}_1^0, \tilde{t}_1^0 \tilde{t}_1^0 \rightarrow \tilde{t}_1^0 \tilde{t}_1^0 + \tilde{L}\tilde{L}^0, \tilde{t}_1^0 \tilde{t}_1^0 \rightarrow \tilde{t}_1^0 \tilde{t}_1^0 + \tilde{L}\tilde{L}^0$		2 τ	E_{T}^{miss}	36.1	$\tilde{t}_1^0 \tilde{t}_1^0$ $\tilde{t}_1^0 \tilde{t}_1^0$ 0.22 $\tilde{t}_1^0 \tilde{t}_1^0$ 0.76	$m(\tilde{L}) = 0.5 m(\tilde{L}) = m(\tilde{L}) = 100 \text{ GeV}$ $m(\tilde{L}) = 0.5 m(\tilde{L}) = m(\tilde{L}) = 100 \text{ GeV}$ $m(\tilde{L}) = m(\tilde{L}) = 100 \text{ GeV}, m(\tilde{L}) = 0.5 m(\tilde{L}) = m(\tilde{L}) = 100 \text{ GeV}$	1708.07875 1708.07875		
$\tilde{L}\tilde{L}, \tilde{L}\tilde{L}, \tilde{L} \rightarrow \tilde{L}\tilde{L}^0$	2 e, μ 2 e, μ	0 jets ≥ 1	E_{T}^{miss} E_{T}^{miss}	139 36.1	\tilde{L} \tilde{L} 0.18 \tilde{L} 0.13-0.23	$m(\tilde{L}) = 0$ $m(\tilde{L}) = m(\tilde{L}) = 5 \text{ GeV}$	ATLAS-COBF-2019-008 1712.08119		
$\tilde{A}\tilde{A}, \tilde{A} \rightarrow \tilde{A}\tilde{G}(\tilde{Z}\tilde{G})$	0 e, μ 4 jets	2 b 0 jets	E_{T}^{miss} E_{T}^{miss}	36.1 36.1	\tilde{A} \tilde{A} 0.3 \tilde{A} 0.29-0.88	$\text{BR}(\tilde{A}) \rightarrow \tilde{A}\tilde{L} = 1$ $\text{BR}(\tilde{A}) \rightarrow \tilde{Z}\tilde{L} = 1$	1806.04030 1804.03602		
Direct $\tilde{L}\tilde{L}^0$ prod., long-lived \tilde{L}^0	Disapp. trk	1 jet	E_{T}^{miss}	36.1	\tilde{L}^0 \tilde{L}^0 0.15 \tilde{L}^0 0.46	Pure Winos Pure Higgsino	1712.02118 ATLAS-CONF-2017-019		
	Stable \tilde{g} R-hadron Metastable \tilde{g} R-hadron, $\tilde{g} \rightarrow \tilde{q}\tilde{q}^0$	Multiple Multiple	36.1 36.1	E_{T}^{miss} E_{T}^{miss}	36.1 36.1	\tilde{g} \tilde{g} 2.0 \tilde{g} 2.05, 2.4	$m(\tilde{L}) = 100 \text{ GeV}$	1902.01626, 1808.04095 1710.04001, 1808.04095	
RPV	LFV $\tilde{p}\tilde{p} \rightarrow \tilde{b}, X, \tilde{L}_i \rightarrow \nu\tilde{q}\tilde{q}^0/\tilde{q}\tilde{q}^0$	$e\mu, \mu\tau$	3-2	E_{T}^{miss}	36.1	\tilde{L}_i 0.62 \tilde{L}_i 1.33 \tilde{L}_i 1.9	$X_{112} = 0.11, A_{11233333} = 0.07$ $m(\tilde{L}) = 100 \text{ GeV}$	1607.98079 1804.03602	
	$\tilde{L}_i^0 \tilde{L}_i^0 \rightarrow W\tilde{Z}\tilde{L}^0/\tilde{L}^0\tilde{L}^0$	4 e, μ 4-5 large- β jets	0 jets Multiple	E_{T}^{miss}	36.1 36.1	$\tilde{L}_i^0 \tilde{L}_i^0$ 1.0 $\tilde{L}_i^0 \tilde{L}_i^0$ 1.8 $\tilde{L}_i^0 \tilde{L}_i^0$ 2.0	Large \tilde{L}_i^0 $m(\tilde{L}) = 200 \text{ GeV}, \text{bino-like}$	1804.03602 ATLAS-COBF-2018-003	
	$\tilde{L}_i \tilde{L}_i \rightarrow \tilde{L}_i^0 \tilde{L}_i^0$	Multiple	Multiple	E_{T}^{miss}	36.1	\tilde{L}_i 0.35 \tilde{L}_i 1.05 \tilde{L}_i 1.05	$m(\tilde{L}) = 200 \text{ GeV}, \text{bino-like}$	ATLAS-COBF-2018-003	
	$\tilde{L}_i \tilde{L}_i \rightarrow \tilde{L}_i^0 \tilde{L}_i^0 + \tilde{b}\tilde{L}^0$	Multiple	Multiple	E_{T}^{miss}	36.1	\tilde{L}_i 0.42 \tilde{L}_i 0.61 \tilde{L}_i 1.05	$m(\tilde{L}) = 200 \text{ GeV}, \text{bino-like}$	ATLAS-COBF-2018-003	
	$\tilde{L}_i \tilde{L}_i \rightarrow \tilde{L}_i^0 \tilde{L}_i^0$	2 e, μ 2 jets + 2 b	Multiple	E_{T}^{miss}	36.7 36.1	\tilde{L}_i 0.42 \tilde{L}_i 0.61 \tilde{L}_i 1.0	$m(\tilde{L}) = 200 \text{ GeV}, \text{bino-like}$	1710.07711	
	$\tilde{L}_i \tilde{L}_i \rightarrow \tilde{L}_i^0 \tilde{L}_i^0$	2 e, μ μ	2 b DV	E_{T}^{miss}	36.1 136	\tilde{L}_i 0.4-1.5 \tilde{L}_i 1.0 \tilde{L}_i 1.6	$\text{BR}(\tilde{L}_i \rightarrow \tilde{L}_i^0/\tilde{L}_i^0) = 20\%$ $\text{BR}(\tilde{L}_i \rightarrow \tilde{q}\tilde{q}^0) = 100\%, 0.09\% = 1$	1710.05544 ATLAS-COBF-2019-006	

*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

10⁻¹ 1 Mass scale [TeV]

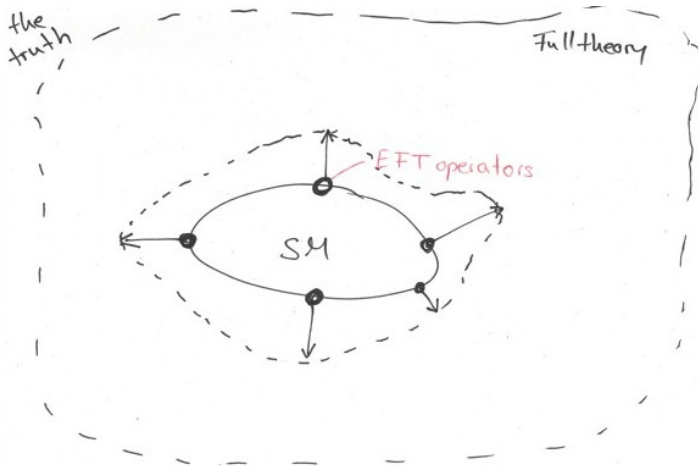
1 Motivation/Expectation

2 Approaches

3 Top-Down

4 Bottom-Up

Bottom-Up Approach - EFT



- We want to extend the SM with minimal additional assumptions:
 - Additional fundamental DM particles \rightarrow fermionic/scalar
 - \mathbb{Z}_2 symmetry to ensure stability of the DM candidate
 - DM interaction with the SM particles (e.g. with quarks)

$$\mathcal{L} = \mathcal{L}_{SM} + \underbrace{i\bar{\chi}\gamma_\mu\partial^\mu\chi - m_\chi\bar{\chi}\chi}_{DM\text{-kinetic terms}} + \underbrace{\sum_q \sum_{i,j} \frac{G_{qij}}{\sqrt{2}} [\bar{\chi}\Gamma_i^X\chi] [\bar{q}\Gamma_q^j q]}_{DM\text{-SM interaction}}$$

- To be general

$$\mathcal{L}_{int.} = \sum_i c_i \mathcal{O}_i$$

Label	Operator	Usual coefficient	Dimension
\mathcal{O}_{M1}	$\bar{\chi}\chi\bar{q}q$	$m_q/2M_*^3$	6
\mathcal{O}_{M2}	$\bar{\chi}i\gamma_5\chi\bar{q}q$	$m_q/2M_*^3$	6
\mathcal{O}_{M3}	$\bar{\chi}\chi\bar{q}i\gamma_5q$	$m_q/2M_*^3$	6
\mathcal{O}_{M4}	$\bar{\chi}i\gamma_5\chi\bar{q}i\gamma_5q$	$m_q/2M_*^3$	6
\mathcal{O}_{M5}	$\bar{\chi}\gamma^\mu\gamma_5\chi\bar{q}\gamma_\mu q$	$1/2M_*^2$	6
\mathcal{O}_{M6}	$\bar{\chi}\gamma^\mu\gamma_5\chi\bar{q}\gamma_\mu\gamma_5q$	$1/2M_*^2$	6
\mathcal{O}_{M7}	$\bar{\chi}\chi G_{\mu\nu}G^{\mu\nu}$	$\alpha_S/8M_*^3$	7
\mathcal{O}_{M8}	$\bar{\chi}\gamma_5\chi G_{\mu\nu}G^{\mu\nu}$	$i\alpha_S/8M_*^3$	7
\mathcal{O}_{M9}	$\bar{\chi}\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$\alpha_S/8M_*^3$	7
\mathcal{O}_{M10}	$\bar{\chi}\gamma_5\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$i\alpha_S/8M_*^3$	7

Table 2: Operators for Majorana DM.

Label	Operator	Usual coefficient	Dimension
θ_{C1}	$\phi^* \phi \bar{q} q$	m_q / M_*^2	5
θ_{C2}	$\phi^* \phi \bar{q} i \gamma_5 q$	m_q / M_*^2	5
θ_{C3}	$\phi^* i \overleftrightarrow{\partial}_\mu \phi \bar{q} \gamma^\mu q$	$1 / M_*^2$	6
θ_{C4}	$\phi^* i \overleftrightarrow{\partial}_\mu \phi \bar{q} \gamma^\mu \gamma_5 q$	$1 / M_*^2$	6
θ_{C5}	$\phi^* \phi G_{\mu\nu} G^{\mu\nu}$	$\alpha_S / 4M_*^2$	6
θ_{C6}	$\phi^* \phi G_{\mu\nu} \tilde{G}^{\mu\nu}$	$\alpha_S / 4M_*^2$	6

Table 3: Operators for Complex Scalar DM.

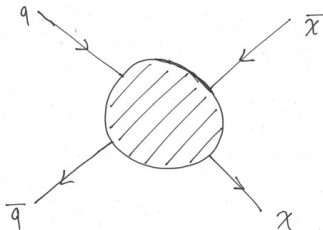
Label	Operator	Usual coefficient	Dimension
\mathcal{O}_{R1}	$\phi^2 \bar{q}q$	$m_q/2M_*^2$	5
\mathcal{O}_{R2}	$\phi^2 \bar{q}i\gamma_5 q$	$m_q/2M_*^2$	5
\mathcal{O}_{R3}	$\phi^2 G_{\mu\nu} G^{\mu\nu}$	$\alpha_S/8M_*^2$	6
\mathcal{O}_{R4}	$\phi^2 G_{\mu\nu} \tilde{G}^{\mu\nu}$	$\alpha_S/8M_*^2$	6

Table 4: Operators for Real Scalar DM.

Let's keep things easy

- Assuming fermionic DM and only vector-like DM-quark interactions

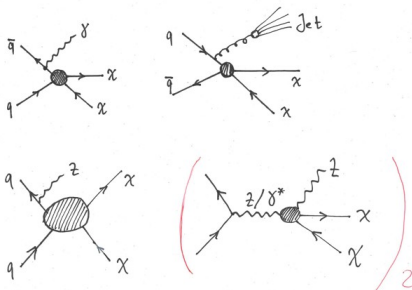
$$\Rightarrow \mathcal{O}_5 = \frac{1}{M_*} (\bar{\chi} \gamma_\mu \chi) (\bar{q} \gamma^\mu q)$$



- Pro: Only two open parameters: m_χ, M_*

Mono-X Searches

- Looking for events with high missing transverse energy (DM-particles are lost in the detection)
- Mono-X events have one particle/jet with high transverse momentum "without" anything to recoil against

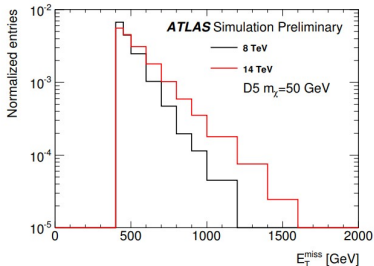
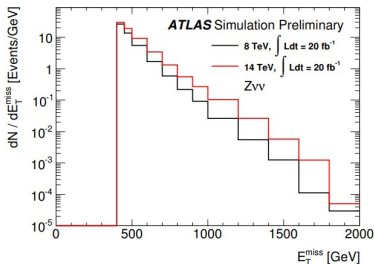


- irreducible background

$$\bar{q}q \rightarrow Z \rightarrow \bar{\nu}\nu$$

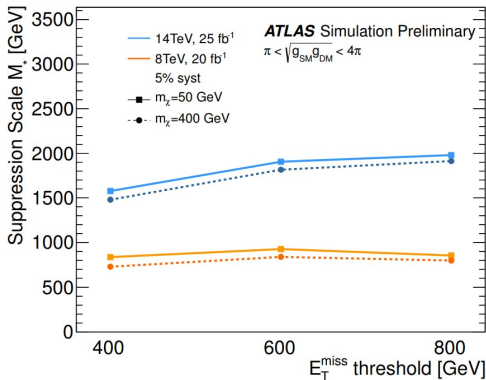
Sensitivity to WIMP Dark Matter in the Final States Containing Jets and Missing Transverse Momentum

- Atlas mono-jet search to pair production of WIMP DM
- $\sqrt{s} = 14\text{TeV}$
- Sensitivity projection
- Monte Carlo of the background (left) and signal (right):



Sensitivity of the Atlas Analysis

- 95% CL lower limit on the suppression scale M_*
- Assumed 5% systematic uncertainty on the SM background
- EFT approach is assumed to be valid



Validity of the EFT

What is the range of validity of the effective theory?

- The suppression scale M_* can be *mapped* to the "UV" theory with a mediator of a mass M_{med} and two couplings g_{SM}, g_χ describing the coupling of the mediator to DM and SM

$$\frac{g_{SM}g_\chi}{Q^2 - M_{med}^2} = -\frac{g_{SM}g_\chi}{M_{med}^2} \left(1 + \frac{Q^2}{M_{med}^2} + \mathcal{O}\left(\frac{Q^4}{M_{med}^4}\right) \right) \approx -\frac{1}{M_*^2}$$

$$\Rightarrow M_* \sim M_{med} / \sqrt{g_{SM}g_\chi}$$

- Pair-annihilation requires

$$M_{med} > 2m_\chi$$

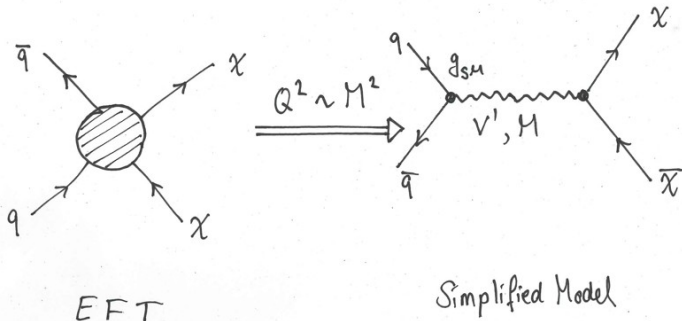
- Perturbativity requires

$$g_{SM}g_\chi \lesssim (4\pi)^2$$

- Validity for

$$m_\chi \lesssim 2\pi M_* \quad \text{and} \quad Q^2 \ll M_{med}^2$$

Bottom-Up Approach - Simplified Model



- Extend the SM with additional minimal field content
 - DM: majorana/dirac fermions, real/complex scalars
 - Mediator: Scalar, Pseudoscalar, Vector ...
- Write down your favorite effective Lagrangian (renormalizable!)
 - 0s0 model: Mediator spin $\rightarrow 0$; DM spin $\rightarrow 0$

$$\mathcal{L}_{0s0} = \underbrace{\frac{1}{2} (\partial_\mu \phi)^2}_{\text{Kinetic terms}} - \underbrace{\frac{1}{2} m_\phi^2 \phi^2 - \frac{\lambda_\phi}{4} \phi^2 H^\dagger H}_{\text{Higgs portal}}$$

- 0s1/2 model: Mediator spin $\rightarrow 0$; DM spin $\rightarrow 1/2$

$$\mathcal{L}_{0s\frac{1}{2}} = \frac{1}{2} (\partial_\mu S)^2 - \frac{m_S^2}{2} S^2 + \bar{\chi}(i\not{\partial} - m_\chi)\chi - g_\chi S \bar{\chi}\chi - g_{SM} S \sum_f \frac{y_f}{\sqrt{2}} \bar{f}f$$

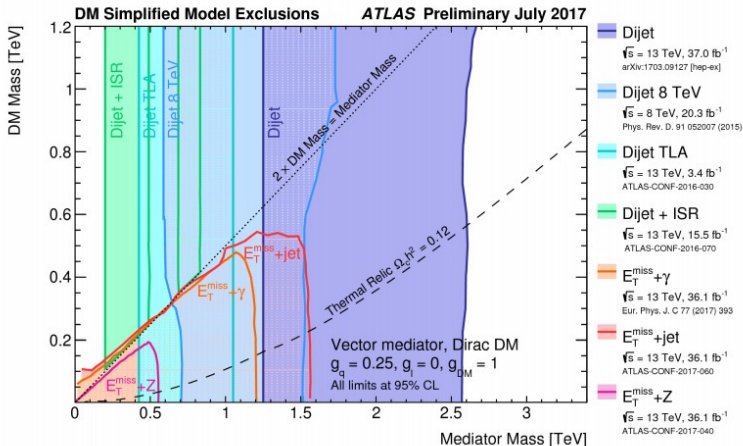
- Vector Mediator: Mediator spin $\rightarrow 1$; DM spin $\rightarrow 1/2$

$$\mathcal{L}_{vec} \supset \frac{1}{2} M_{med}^2 V_\mu V^\mu - g_{DM} V_\mu \bar{\chi} \gamma^\mu \chi - \sum_q g_{SM}^q V_\mu \bar{q} \gamma^\mu q$$

- Start calculating...

Simplified Model Results

V1



- DM collider searches are *challenging* (small fluxes, small production rates,...)
- Two general approaches for model building/searches
 - Top-down
 - + Particle Content is *known*
 - +/- Direct limits on model parameters
 - Each model needs a detailed and involved analysis
 - Bottom up
 - + *Model-independent* searches **and** limits!
 - Huge amount of free parameters in the EFT operators (most of the time only a subset of operators)
 - + Simplified models can handle single problems without (*unnecessary*) huge amount of workload
 - At the most shows only the direction; not the end of the story!

Thank you for your attention!

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Thank you for your attention!

- Title Picture

<https://www.symmetrismagazine.org/article/december-2013/four-things-you-might-not-know-about-dark-matter>

- SUSY searches:

- ArXiv: hep-ex/9809031 (OPAL search)
- ArXiv:1709.05406v2 (LEP search)

- Simplified/EFT:

- ArXiv: 1603.08002v2
- ATL-PHYS-PUB-2014-007