How to model WIMP interactions  ${\scriptstyle 000000000000}$ 

How to detect a WIMP anyway?

Conclusions

### An introduction to direct dark matter detection

Markus Eichhorn

24th of June 2019

How to model WIMP interactions

How to detect a WIMP anyway?

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### DM interactions

DM must interact via gravity, but maybe there are further interactions.

Common models are

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In this talk we focus on

• WIMP DM with mass ranges  $\mathcal{O}(\text{GeV})$ 

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fermionic DM

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- Higgs exchange (Higgs portal models)
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In this talk we focus on

- WIMP DM with mass ranges  $\mathcal{O}(\text{GeV})$
- fermionic DM
- WIMP: Dirac fermion DM (no Majorana Fermion)

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#### Astrophysical assumptions

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- a galactic DM halo with a Maxwell-Boltzmann velocity distribution f(v)
- ▶ escape velocity of  $v_{\rm esc} \sim 500 \, {\rm km \over \rm s}$  is maximal velocity for particle DM

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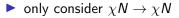
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#### Assumptions for the scattering



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- ▶ only consider  $\chi N \rightarrow \chi N$
- ▶ no excitation  $\chi'$  of WIMP particle

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- ▶ only consider  $\chi N \rightarrow \chi N$
- no excitation  $\chi'$  of WIMP particle
- ▶ also no excitation N' of the nucleus  $\Rightarrow$  small or vanishing momentum transfer in the elastic scattering process  $\mathcal{O}(MeV)$

Conclusions

# Kinematics of elastic scattering I before scattering

In laboratory frame we have  $oldsymbol{p}_{\chi}=m_{\chi}oldsymbol{v}_{\chi}$  and  $oldsymbol{p}_{\mathrm{N}}=0$ 

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$$\Rightarrow \tilde{\mathbf{p}}_{\chi} = \mu \mathbf{v}_{\chi}$$

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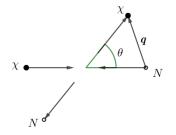
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Conclusions

## Kinematics of elastic scattering II after scattering

We consider elastic scattering, so  $|\tilde{\boldsymbol{p}}_{\chi}| = |\tilde{\boldsymbol{p}}_{\chi}'|$ But we also have the momentum transfer  $\boldsymbol{q} = \tilde{\boldsymbol{p}}_{\chi}' - \tilde{\boldsymbol{p}}_{\chi}$ , and thus

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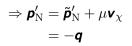


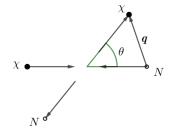
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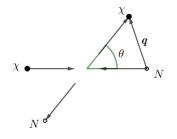
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$$\Rightarrow \mathbf{p}'_{\mathrm{N}} = \widetilde{\mathbf{p}}'_{\mathrm{N}} + \mu \mathbf{v}_{\chi}$$
  
=  $-\mathbf{q}$ 

where  $q^2$  evaluates to

$$oldsymbol{q}^2 = ilde{oldsymbol{p}}_{\chi}^{\prime 2} - 2 ilde{oldsymbol{p}}_{\chi}^{\prime} ilde{oldsymbol{p}}_{\chi} + ilde{oldsymbol{p}}_{\chi}^2 \ = 2(\mu oldsymbol{v}_{\chi})^2 (1 - \cos heta)$$



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## The recoil energy

The only momentum the nucleus carries is the momentum transfer  ${\pmb q}$  thus

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▶ if  $m_{\chi} \gg m_{\rm N}$  then  $\mu \approx m_{\rm N}$ ▶ if  $m_{\chi} \approx m_{\rm N}$  then  $\mu \approx \frac{m_{\rm N}}{2}$ 

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If we consider  $E_{
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If we consider  $E_{\rm R}^{\rm th}$  such that only  $E_{\rm R} > E_{\rm R}^{\rm th}$  is detectable. The minimal velocity for a scattering angle  $\theta = 180^{\circ}$  is given by

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- $\blacktriangleright$  lower threshold  $E_{\rm R}^{\rm th}$
- ► achieve  $m_{\rm N} \approx m_{\chi}$

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## Recoil spectrum

The recoil spectrum reads

$$\frac{\mathrm{d}R}{\mathrm{d}E_{\mathrm{R}}} = N_{\mathrm{T}}n_{\chi}\left\langle v_{\chi}\frac{\mathrm{d}\sigma}{\mathrm{d}E_{\mathrm{R}}}\right\rangle$$

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- $\langle \cdot \rangle$  denotes the average over velocities

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#### Final form of recoil spectrum

We find for a fixed velocity v

$$\mathrm{d}E_{\mathrm{R}} = \mathrm{d}\cos\theta \; \frac{\mu^2}{m_{\mathrm{N}}} v^2$$

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where f(v) is a halo model dependent velocity propability function

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#### Now onto finding the cross section

To find the cross section we need the squared matrix elements

 $\sigma \sim \overline{|\mathcal{M}|^2}$ 

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 $\Rightarrow$  Need SM extentions

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Conclusions

### What to bear in mind?

momentum transfer in MeV region

elastic scattering without exitation of nucleus or WIMP

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## What to bear in mind?

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- consider effective field theories like Fermi's interaction

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Many operators are possible

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contact interaction with quark

$$\mathcal{O}_{\mathrm{D1}} = \bar{\chi}\chi\bar{q}q \quad c_{\mathrm{D1}}^{\mathrm{q}} = \frac{m_{\mathrm{q}}}{M_{*}^{3}}$$

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#### axial interaction with quarks

$$\mathcal{O}_{\mathrm{D8}} = \bar{\chi} \gamma^{\mu} \gamma_5 \chi \bar{q} \gamma_{\mu} \gamma_5 q \quad c_{\mathrm{D8}}^{\mathrm{q}} = \frac{1}{M_*^2}$$

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## Some remarks

 $\blacktriangleright$   $M_*$  is the energy scale of the interaction

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$$\begin{split} \bar{\chi} \Gamma \chi \bar{q} \Gamma' q \\ \bar{\chi} \Gamma \chi G_{\mu\nu} G^{\mu\nu} \quad \bar{\chi} \Gamma \chi G_{\mu\nu} \tilde{G}^{\mu\nu} \end{split}$$

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►  $\Gamma, \Gamma' \in \{\gamma_{\mu}, i\gamma_5, \gamma_{\mu}\gamma_5, \sigma_{\mu\nu}\}$  is a fermion bilinear structure ► full list is given in [1603.08002]

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- ►  $\Gamma, \Gamma' \in \{\gamma_{\mu}, i\gamma_5, \gamma_{\mu}\gamma_5, \sigma_{\mu\nu}\}$  is a fermion bilinear structure
- full list is given in [1603.08002]
- two more important operators are

$$\mathcal{O}_5 = ar{\chi} \gamma^\mu \chi ar{q} \gamma_\mu q$$
  
 $\mathcal{O}_7 = ar{\chi} \gamma^\mu \chi ar{q} \gamma_\mu \gamma_5 q$ 

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#### How to go to the nucleon level

Sum over operators with same tensor structure,

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## How to go to the nucleon level

Sum over operators with same tensor structure, example for  $\mathcal{O}_{D1}$  and  $\mathcal{O}_{D11}$ 

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#### How to go to the nucleon level

# Sum over operators with same tensor structure, example for $\mathcal{O}_{D1}$ and $\mathcal{O}_{D11}$ Reminder

$$\mathcal{O}_{\mathrm{D1}} = \bar{\chi} \chi \bar{q} q \quad \mathcal{O}_{\mathrm{D11}} = \bar{\chi} \chi \mathcal{G}_{\mu\nu} \mathcal{G}^{\mu\nu}$$

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Sum over operators with same tensor structure, example for  $\mathcal{O}_{D1}$  and  $\mathcal{O}_{D11}$  Reminder

$$\mathcal{O}_{\mathrm{D}1} = \bar{\chi} \chi \bar{\boldsymbol{q}} \boldsymbol{q} \quad \mathcal{O}_{\mathrm{D}11} = \bar{\chi} \chi \boldsymbol{G}_{\mu\nu} \boldsymbol{G}^{\mu\nu}$$

We now consider

$$ilde{\mathcal{O}} := \sum_{\mathrm{q}} oldsymbol{c}_{\mathrm{D1}}^{\mathrm{q}} \mathcal{O}_{\mathrm{D1}}^{\mathrm{q}} + \sum_{\mathrm{g}} oldsymbol{c}_{\mathrm{D11}}^{\mathrm{g}} \mathcal{O}_{\mathrm{D11}}^{\mathrm{g}}$$

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Quarks and gluons are all constituents of a nucleon

Conclusions

# Going to the nucleon level

We now consider matrix elements like  $\langle N | \tilde{\mathcal{O}}' | N \rangle$  where  $N \in \{n, p\}$ and  $\tilde{\mathcal{O}}' := \bar{\chi} \chi \mathcal{A}$ 

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part of the nucleus, thus

$$\langle N|\bar{\chi}\chi\mathcal{A}|N\rangle = \bar{\chi}\chi\langle N|\mathcal{A}|N\rangle$$

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$$\langle N | \bar{\chi} \chi \mathcal{A} | N \rangle = \bar{\chi} \chi \langle N | \mathcal{A} | N \rangle$$

We map the combinations of operators on parton level to an operator at nucleon level

$$\mathcal{O}_{\mathrm{D1}}^{\mathrm{N}} = C \bar{\chi} \chi \bar{N} N$$

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# What is C?

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The coefficient C is determined by

 $\blacktriangleright$  previous coefficients  $\textit{c}_{D1}^{\rm q}$  and  $\textit{c}_{D11}^{\rm g}$ 

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# What is C?

- $\blacktriangleright$  previous coefficients  $\textit{c}_{D1}^{\rm q}$  and  $\textit{c}_{D11}^{\rm g}$
- form factors

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# What is C?

- previous coefficients  $c_{D1}^{q}$  and  $c_{D11}^{g}$
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  - main constituents are up, down quark and also strange as sea quark

Conclusions

# What is C?

- previous coefficients  $c_{D1}^{q}$  and  $c_{D11}^{g}$
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  - gluons play a secondary role while remaining sea quarks charm, top, bottom play a minor role

Conclusions

## What is C?

The coefficient C is determined by

- previous coefficients  $c_{\mathrm{D1}}^{\mathrm{q}}$  and  $c_{\mathrm{D11}}^{\mathrm{g}}$
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- the full expression reads

$${\cal C} = \sum_{
m q=u,d,s} c_{
m D1}^{
m q} rac{m_{
m N}}{m_{
m q}} f_{
m q}^{
m (N)} + rac{2}{27} f_{
m G}^{
m (N)} \left( \sum_{
m q=c,b,t} c_{
m D1}^{
m q} rac{m_{
m N}}{m_{
m q}} - rac{1}{3\pi} c_{
m D11}^{
m g} m_{
m N} 
ight)$$

How to detect a WIMP anyway?

Conclusions

#### Understanding the expression for C

$$C = \sum_{\text{q=u,d,s}} c_{\text{D1}}^{\text{q}} \frac{m_{\text{N}}}{m_{\text{q}}} f_{\text{q}}^{(\text{N})} + rac{2}{27} f_{\text{G}}^{(\text{N})} \left( \sum_{\text{q=c,b,t}} c_{\text{D1}}^{\text{q}} \frac{m_{\text{N}}}{m_{\text{q}}} - rac{1}{3\pi} c_{\text{D11}}^{\text{g}} m_{\text{N}} 
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How to model WIMP interactions

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How to model WIMP interactions

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  - identify contribution of heavier quarks by gluon contribution.
     A detailed calculation yields the prefactors

How to detect a WIMP anyway?

Conclusions

### General case

In general we consider

$$\mathcal{L}_{ ext{eff}} = \sum_{ ext{q},i} oldsymbol{c}_i^{ ext{q}} \mathcal{O}_i^{ ext{q}} + \sum_{ ext{g},i} oldsymbol{c}_i^{ ext{g}} \mathcal{O}_i^{ ext{g}}$$

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Conclusions

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How to detect a WIMP anyway? 000000000

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again we have  $\mathcal{O}_k^{\mathrm{N}} = \bar{\chi} \Gamma \chi \bar{N} \Gamma' N$  with  $\Gamma, \Gamma' \in \{\gamma_\mu, i\gamma_5, \gamma_\mu \gamma_5, \sigma_{\mu\nu}\}$  a full list is given in [1603.08002].

How to detect a WIMP anyway?

Conclusions

#### Important operators on nucleon level

$$\mathcal{O}_{\mathrm{D1}}^{\mathrm{N}} = \bar{\chi} \chi \bar{N} N$$
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How to detect a WIMP anyway?

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$$\mathcal{O}_{\mathrm{D8}}^{\mathrm{N}} = \bar{\chi}\gamma^{\mu}\gamma_{5}\chi\bar{N}\gamma_{\mu}\gamma_{5}N$$
$$\mathcal{O}_{\mathrm{D9}}^{\mathrm{N}} = \bar{\chi}\sigma^{\mu\nu}\chi\bar{N}\sigma_{\mu\nu}N$$

$$\sigma^{\mu\nu} = \frac{\mathrm{i}}{2} \left[ \gamma^{\mu}, \gamma^{\nu} \right]$$

Conclusions

## Exploiting the non relativistic (NR) limit

Since the WIMPs are non relativistic, the operators can be simplified.

Conclusions 00

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The expression of the nucleon operator yields the relations

$$egin{aligned} \langle \mathcal{O}_{\mathrm{D8}}^{\mathrm{N}} 
angle &= -rac{1}{2} \langle \mathcal{O}_{\mathrm{D9}}^{\mathrm{N}} 
angle &= -16 m_{\chi} m_{\mathrm{N}} \mathcal{O}_{4}^{\mathrm{NR}} \ \langle \mathcal{O}_{\mathrm{D1}}^{\mathrm{N}} 
angle &= \langle \mathcal{O}_{\mathrm{D5}}^{\mathrm{N}} 
angle &= 4 m_{\chi} m_{\mathrm{N}} \mathcal{O}_{1}^{\mathrm{NR}} \end{aligned}$$

How to detect a WIMP anyway?

Conclusions

### Further NR operators

We had

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How to detect a WIMP anyway?

Conclusions

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$$v^2$$
,  $i(\boldsymbol{S}_{\chi} \times \boldsymbol{q}) \cdot \boldsymbol{v}$ ,  $\boldsymbol{v}^{\perp} \cdot \boldsymbol{S}_N$ 

Conclusions 00

### Nucleus level

The cross section with the nucleus is given by the coherent sum of the interaction with a nucleon

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spin independent needs no extra input besides number of protons and neutrons

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 For non zero momentum transfer additional form factors are needed

How to model WIMP interactions

How to detect a WIMP anyway? •000000000

Conclusions

## On to detecting(?) the WIMP



now we want to measure the nuclear recoil

How to model WIMP interactions

How to detect a WIMP anyway? ••••••• Conclusions

# On to detecting(?) the WIMP

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How to detect a WIMP anyway? •••••••• Conclusions

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How to detect a WIMP anyway? ••••••• Conclusions

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How to detect a WIMP anyway? ••••••• Conclusions

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How to detect a WIMP anyway? ••••••• Conclusions

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- experimental challenge: overcome backgrounds

How to model WIMP interactions

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### What backrounds are we talking about?

The most important backgrounds are...

How to model WIMP interactions  ${\scriptstyle 000000000000}$ 

How to detect a WIMP anyway? 000000000

Conclusions

#### What backrounds are we talking about?

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electron recoil background

Conclusions

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Conclusions

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- neutrino floor

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How to model WIMP interactions

How to detect a WIMP anyway?

Conclusions

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Distinct behaviour from nuclear recoil.

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Electrons in the detector are produced by

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Both come from radioisotopes in surrounding rock and construction material.

Distinct behaviour from nuclear recoil. The keyword to bare in mind here is: quenching.

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#### Nuclear recoil background

Produced by

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Conclusions

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Conclusions

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muons

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Conclusions

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Conclusions

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Most importantly the neutrons are

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- sourced by primordial chains and CR muons
- luckily multiple scattering processes and thus distinguishable from WIMP interaction

How to model WIMP interactions

Conclusions

## Neutrino floor

Conclusions 00

## Neutrino floor

The neutrino floor is sometimes called the ultimate background

neutrinos will interact with the nucleus via the weak interaction

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- $\blacktriangleright$  for low WIMP masses the interesting neutrinos are solar ones and created by  $^8{\rm B}$  decay

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- However there are hypothesis how to overcome the neutrino floor by looking for annual modulations

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Conclusions



Conclusions

- use scintillation and ionization
- usually two phase detectors with liquid phase and gas phase

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- Xenon is liquid at 165.2 K and provides direct measurement of scintillation
- Argon allows for pulse shape discrimination but contains <sup>39</sup>Ar. It is liquid below 87.2 K

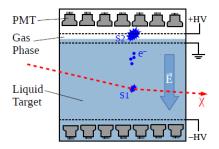
How to model WIMP interactions

How to detect a WIMP anyway?  $\tt 0000000000$ 

Conclusions

#### How to meassure a nuclear recoil in LXe?

 scintillation via photomultipliers (S1)

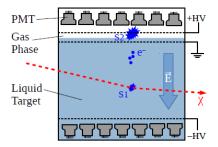


How to model WIMP interactions

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Conclusions

- scintillation via photomultipliers (S1)
- use electric field of 1 kV/cm
   to bring electrons to gas
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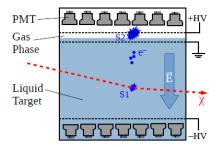


How to model WIMP interactions

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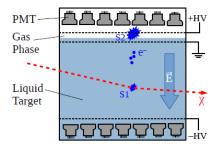


How to model WIMP interactions

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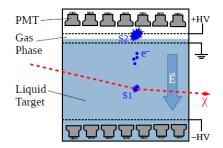


How to model WIMP interactions

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- from time difference determine location of interaction (time projection chamber)



How to model WIMP interactions

How to detect a WIMP anyway?

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#### The Xenon1T experiment

Xenon1T is a liquid noble gas detector using Xe

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How to detect a WIMP anyway?

Conclusions

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How to model WIMP interactions

How to detect a WIMP anyway?

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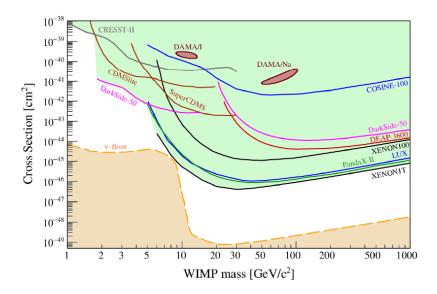
#### The Xenon1T experiment

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- many stages XENON10, XENON100, XENON1T

How to model WIMP interactions

How to detect a WIMP anyway? 000000000 Conclusions

## The current limits for SI interaction [arXiv:1903.03026]

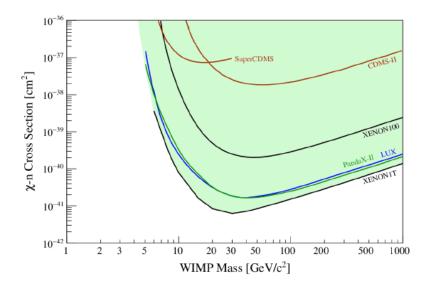


How to model WIMP interactions

How to detect a WIMP anyway?

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How to model WIMP interactions  ${\scriptstyle 000000000000}$ 

How to detect a WIMP anyway?

Conclusions •0

# Conclusions

In this talk we covered

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How to detect a WIMP anyway?

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## Conclusions

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How to detect a WIMP anyway?

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How to detect a WIMP anyway?

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## Conclusions

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- ▶ the procedure to go from the parton level to the nucleus level
- The Xenon1T experiment and its results

How to model WIMP interactions

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Thank you for your attention.

#### sources

- Marc Schumann, Direct detection of WIMP dark matter, (2019), [arXiv:1903.03026]
- Andrea De Simone, Thomas Jacques, Simplified models vs. effective field theory approaches in dark matter searches, (2016), [arXiv:1603.08002]
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