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Information regarding the exercise course:

The exercise sheets are published each Wednesday on the webpage of the course, see bottom of this page, and handed out in the lecture on Thursday.

Scheinkriterium: In order to successfully pass “Theoretische Teilchenphysik II” we request you to

- visit at least 10 exercise classes (We collect signatures!).
- present parts of/one/two exercises at the blackboard (in total $\sim 45 - 60$ minutes).
- register by writing an email to “stefan.liebler@kit.edu” before 26.10.2018.

Exercise 1: $SU(N)$ representations

The Lie groups $SU(N)$ are of particular importance in particle physics as they for example form the gauge structures of the Weinberg-Salam model and quantum chromodynamics, which are associated with the Lie groups $SU(2)$ and $SU(3)$, respectively. Also $SU(3)$ explains the flavour structure in the quark model (Eightfold Way). A representation of a group is a mapping on a set of matrices that fulfill the group axioms. A particularly important representation of $SU(N)$ groups is the *fundamental* representation, an irreducible matrix representation of dimension N , which is formed by the generators

$$T_{ij}^a, \quad a = 1, \dots, N^2 - 1, \quad i, j = 1, \dots, N,$$

which are hermitian, i.e. $T^{a\dagger} = T^a$, and traceless, which implies $\text{Tr}(T^a) = 0$. They are normalized through the relation $\text{Tr}(T^a T^b) = \frac{1}{2} \delta^{ab}$. For $N = 2$ and 3 the generators of the *fundamental* representation are given by

- the Pauli matrices $T^a = \frac{\sigma^a}{2}$ for $SU(2)$
- the Gell-Mann matrices $T^a = \frac{\lambda^a}{2}$ for $SU(3)$.

For the *fundamental* representation the generators fulfill the following commutation and anti-commutation relations

$$[T^a, T^b] = i f_{abc} T^c, \quad (1)$$

$$\{T^a, T^b\} = \frac{1}{N} \delta^{ab} \mathbb{1}_N + d_{abc} T^c, \quad (2)$$

which in turn defines the totally antisymmetric structure constants f_{abc} and the totally symmetric symbols d_{abc} . The commutation relation in Eq. (1) is actually valid for any representation of $SU(N)$ and is the so-called Lie algebra of $SU(N)$.¹ In the *fundamental* representation each

¹Another matrix representation is given by the structure constants f_{abc} themselves, the *adjoint* representation.

complex $(N \times N)$ -matrix M can be written in terms of the $N^2 - 1$ generators and the unit matrix in the form

$$M = c_0 \mathbb{1} + \sum_{a=1}^{N^2-1} c_a T^a. \quad (3)$$

- (a) Show that the orthogonality relation in combination with Eq. (3) and the tracelessness of the generators in the *fundamental* representation yield the Fierz identity of $SU(N)$ given by

$$T_{ij}^a T_{kl}^a \equiv \sum_{a=1}^{N^2-1} T_{ij}^a T_{kl}^a = \frac{1}{2} \delta_{il} \delta_{jk} - \frac{1}{2N} \delta_{ij} \delta_{kl}. \quad (4)$$

- (b) Show that for any representation of $SU(N)$

$$C_2 = T^a T^a$$

with a sum over a is a Casimir invariant, i.e. $[C_2, T^a] = 0$ for all generators T^a .

- (c) Show that the structure constants f_{abc} and d_{abc} are real by using the hermiticity of the generators of the *fundamental* representation.
- (d) Show that the $(N \times N)$ matrices

$$\mathcal{T}^a = -T^{a*}$$

are also a representation of $SU(N)$, i.e. they fulfill the Lie algebra given in Eq. (1). These \mathcal{T}^a are the generators of the so-called \bar{N} (“ N -bar”) representation of $SU(N)$.

Hint: T^{a*} is the complex conjugated, not the hermitian matrix!

- (e) Calculate the value of C_2 for the two N -dimensional representations, the (N) as well as the (\bar{N}) representation.
- (f) Use the Jacobi identity for mixed commutation and anti-commutation relations

$$[A, \{B, C\}] + [B, \{C, A\}] + [C, \{A, B\}] = 0$$

to obtain a relation connecting f_{abc} and d_{abc} . *Hint:* Insert the generators.

- (g) Use the previous result to show that

$$C_3 = d_{abc} T^a T^b T^c$$

is a Casimir invariant of $SU(N)$.

- (h) Use Eq. (4), to calculate the sum over a in

$$T^a T^b T^a.$$

- (i) Determine the value of C_3 for both the (N) and the (\bar{N}) representation. The result shows, that (N) and (\bar{N}) are not equivalent representations for $N > 2$.²

²As indicated in the lecture those two representations are identified with quarks and anti-quarks, respectively.