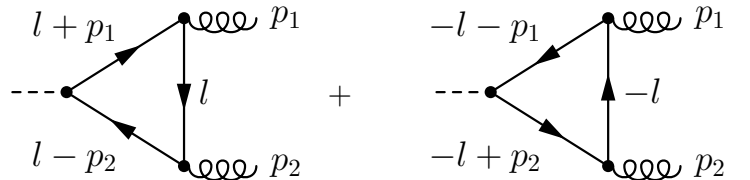


Exercises: Stefan Liebler (stefan.liebler@kit.edu) (Office 12/03 - Build. 30.23)  
Shruti Patel (shruti.patel@kit.edu) (Office 12/14 - Build. 30.23)

Even if you finished your TTP2 duties, it's worth to come by and be amazed about the beauty of these exercises.

### Exercise 1: Higgs boson decay into gluons - Part 1

The aim of this exercise is to calculate the partial decay width of the Standard Model Higgs boson into a pair of gluons,  $h^0 \rightarrow gg$ , in the first non-vanishing order. The decay is loop-mediated, i.e. the Higgs boson couples to two gluons through quark loops. The quark running in the loop with mass  $m$  couples to the Higgs boson with the Yukawa coupling  $y_q = \frac{m}{v}$  with the vacuum expectation value  $v = 1/\sqrt{\sqrt{2}G_F}$ . The relevant two Feynman diagrams for each quark, depicting the four-momenta, are given by



The two final-state gluons have outgoing momenta  $p_1$  and  $p_2$  as well as Lorentz indices  $\mu$  and  $\nu$  and colors  $a$  and  $b$ , respectively. Accordingly, the initial-state Higgs boson has momentum  $p_1 + p_2$ . We want all particles to be on-shell, i.e.  $(p_1 + p_2)^2 = m_{h^0}^2$ ,  $p_1^2 = p_2^2 = 0$ . Since there are no tree-level diagrams and thus no counterterms, the final result of the loop diagrams cannot develop an ultraviolet divergence.

- (a) Show that the amplitude in dimensional regularisation ( $d = 4 - 2\epsilon$ ) involving one quark  $q$  with mass  $m$  is of the form

$$\mathcal{M}_q = \epsilon_{1,\mu}^* \epsilon_{2,\nu}^* (ig_s)^2 (-iy_q) (-1) i^3 \mu^{4-d} \int \frac{d^d l}{(2\pi)^d} \frac{T_{ij}^a T_{ji}^b \text{Tr}[S^{\mu\nu}]}{(l^2 - m^2)((l + p_1)^2 - m^2)((l - p_2)^2 - m^2)}$$

with  $S^{\mu\nu} = \gamma^\mu (\not{l} + \not{p}_1 + m)(\not{l} - \not{p}_2 + m)\gamma^\nu (\not{l} + m) + (-\not{l} + m)\gamma^\nu (-\not{l} + \not{p}_2 + m)(-\not{l} - \not{p}_1 + m)\gamma^\mu$ .

- (b) Show that  $\text{Tr}[S^{\mu\nu}] = 8m(g^{\mu\nu}(m^2 - l^2 - p_1 \cdot p_2) + 4l^\mu l^\nu + p_2^\mu p_1^\nu)$ . Argue why the second term in  $S^{\mu\nu}$  yields the same contribution as the first term.
- (c) Introduce Feynman parameters in the form

$$\frac{1}{abc} = 2 \int dx dy dz \frac{\delta(1-x-y-z)}{(xa + yb + zc)^3} = 2 \int_0^1 dy \int_0^{1-y} dz \frac{1}{((1-y-z)a + yb + zc)^3}$$

and shift the loop momentum  $l$  such that the denominator takes the form  $(l^2 - (zp_2 - yp_1)^2 - m^2)^3 = (l^2 + yzm_{h^0}^2 - m^2)^3 =: (l^2 - M^2)^3$ . Transform the numerator accordingly.

(d) Use the tensor integrals from sheet 7 to show that

$$\mathcal{M}_q = -\epsilon_{1,\mu}^* \epsilon_{2,\nu}^* g_s^2 y_q \frac{\delta^{ab}}{24\pi^2} \frac{m_{h^0}^2}{m} \epsilon_{1,\mu}^* \epsilon_{2,\nu}^* \left( g^{\mu\nu} - \frac{2}{m_{h^0}^2} p_1^\nu p_2^\mu \right) f \left( \frac{m_{h^0}^2}{m^2} \right)$$

with

$$f(x) = 3 \int_0^1 dy \int_0^{1-y} dz \frac{1 - 4yz}{1 - xyz}.$$

*Hint:*  $g^{\mu\nu} I_d(0, 2, M^2) + 4I_d^{\mu\nu}(0, 3, M^2)$  might be a useful relation.

(e) Check gauge invariance explicitly by replacing the polarisation vector of each gluon through the corresponding momentum. Introduce a sum over different quarks  $\mathcal{M} = \sum_q \mathcal{M}_q$  with masses  $m_q$ . Square the amplitude  $\mathcal{M}$  and perform the polarisation sum over the gluon polarisation. You should obtain

$$|\mathcal{M}|^2 = \alpha_s^2 \sqrt{2} G_F \frac{4m_{h^0}^4}{9\pi^2} \left| \sum_q f \left( \frac{m_{h^0}^2}{m_q^2} \right) \right|^2.$$

(f) Finally calculate the partial decay width, which is given by

$$\Gamma(h^0 \rightarrow gg) = \frac{\alpha_s^2 G_F m_{h^0}^3}{36\pi^3 \sqrt{2}} \left| \sum_q f \left( \frac{m_{h^0}^2}{m_q^2} \right) \right|^2.$$

## Exercise 2: Higgs boson decay into gluons - Part 2

We continue with the previous exercise. *Hint:* This exercise can be performed independently. All relevant results are given in the previous exercise.

(a) Perform the integrations in the definition of  $f(x)$  by using the results obtained for the dilogarithm on sheet 11. *Hint:* Perform the integration over  $z$  and determine the roots of the argument of the remaining logarithm named  $y_\pm$ , such that  $1 - xy + xy^2 = (y_+ - y)(y_- - y)x$ . Split the logarithm, use sheet 11 and use

$$\arcsin(z) = -i \ln(iz \pm \sqrt{1 - z^2}) = -i \ln \left[ \left( \frac{\sqrt{z^2 - 1} + z}{\sqrt{z^2 - 1} - z} \right)^2 \right].$$

If you succeed, you should get for  $x < 4$

$$f(x) = \frac{6}{x} - \frac{6(4-x)}{x^2} \arcsin^2 \left( \frac{\sqrt{x}}{2} \right).$$

- (b) What value does  $f(x)$  take for heavy quarks (heavy-top limit), i.e.  $m_q \rightarrow \infty$  and  $m_{h^0}^2/m_q^2 \rightarrow 0$ ? Why does the measurement of the decay of a Higgs boson to gluons or the production of a Higgs boson from gluons allow to make a statement on the number of heavy quark generations?
- (c) If you still don't have enough, try to obtain the squared amplitude using a computer, e.g. with **FeynArts** and **FormCalc**. You might encounter the Passarino-Veltman representation of the loop integrals, for which computer codes exist that allow their numerical evaluation.