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### Information regarding the exercise course:

The exercise sheets are published each Tuesday on the webpage of the course, see below.

Scheinkriterium: In order to successfully pass “Theoretische Teilchenphysik II” we request you to

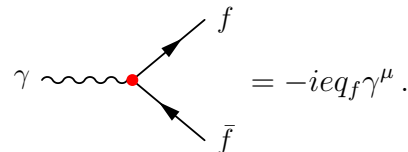
- visit at least 10 exercise classes (We collect signatures!).
- present parts of one/two exercises at the blackboard (in total  $\sim 45 - 60$  minutes).
- register by writing an email to “stefan.liebler@kit.edu” specifying your full name, your “Matrikelnummer” and your class preference (see webpage) before 20.10.2019.

### Exercise 1: Deep inelastic scattering

We consider the scattering of an electron and a quark, i.e. the process

$$e^-(p_1) + q(p_2) \rightarrow e^-(p_3) + q(p_4).$$

This process is mediated through the exchange of a virtual photon. The relevant Feynman rule of two fermions interacting with a photon is given by



$$\gamma \text{ (wavy line)} \rightarrow f \text{ (arrow)} + \bar{f} \text{ (arrow)} = -ieq_f \gamma^\mu.$$

Therein  $q_f$  is the electromagnetic charge of the fermion in units of the elementary charge  $e$ .

- Draw the relevant Feynman diagrams for this process.
- Deduce the matrix element  $\mathcal{M}$  for an arbitrary quark with charge  $q_q$ . *Hints:* Neglect the exchange of a virtual  $Z$  boson and color factors, as the average/sum over the color charge will not yield more than a factor of 1. Use Dirac spinors for the external fermions.
- Calculate the squared matrix element  $|\mathcal{M}|^2$  for massless electrons and quarks. Calculate  $\overline{\sum_{\text{spins}} |\mathcal{M}|^2}$ , i.e. average/sum over the spins of incoming/outgoing fermions, and use techniques from TTP1 to calculate the traces over the  $\gamma$  matrices.

*Hint:* Useful identities are:

$$\begin{aligned} \text{Tr}(\gamma^\mu \gamma^\nu) &= 4g^{\mu\nu} \\ \text{Tr}(\gamma^{\mu_1} \dots \gamma^{\mu_n}) &= 0 \quad \text{for odd } n \\ \text{Tr}(\gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma) &= 4 \cdot (g^{\mu\nu} g^{\rho\sigma} - g^{\mu\rho} g^{\nu\sigma} + g^{\mu\sigma} g^{\nu\rho}). \end{aligned}$$

- (d) We introduce the Mandelstam variables  $s = (p_1 + p_2)^2$  and  $t = (p_1 - p_3)^2$ . The cross section formula for the scattering of two particles with momenta  $p_1$  and  $p_2$  and masses  $m_1$  and  $m_2$  into two particles with momenta  $p_3$  and  $p_4$  and masses  $m_3$  and  $m_4$  reads

$$d\sigma = d\Phi_2 \frac{1}{4\sqrt{(p_1 \cdot p_2)^2 - m_1^2 m_2^2}} \sum_{\text{spins}} |\mathcal{M}|^2$$

with the two-particle phase space

$$d\Phi_2 = \frac{d^3 p_3}{(2\pi)^3 2E_3} \frac{d^3 p_4}{(2\pi)^3 2E_4} (2\pi)^4 \delta^{(4)}(p_1 + p_2 - p_3 - p_4).$$

Therein  $E_3$  and  $E_4$  are the energies of the two outgoing particles. In TTP1 you already showed that in the center-of-mass frame

$$d\Phi_2 = d\Omega \frac{\lambda(s, m_3^2, m_4^2)}{32\pi^2 s},$$

where  $d\Omega = d \cos \theta_3 d\phi_3$  includes the angles of particle 3 with respect to particle 1 and we neglect  $\Theta$  functions and introduced the *Källén function*

$$\lambda(a, b, c) = \sqrt{a^2 + b^2 + c^2 - 2ab - 2ac - 2bc}.$$

Show that in the center-of-mass frame this leads to

$$\frac{d\sigma}{dt} = \frac{1}{16\pi \lambda^2(s, m_1^2, m_2^2)} \sum_{\text{spins}} |\mathcal{M}|^2,$$

when the remaining integration over the angle  $\cos \theta_3$  is expressed in terms of the Mandelstam variable  $t$ . This latter expression is independent of the reference frame.

- (e) Calculate the differential cross section  $d\sigma/dt$  as a function of the Mandelstam variables  $s$  and  $t$  for deep inelastic scattering, again in the limit of massless fermions.  
*Note:* Deep inelastic scattering is actually the scattering of an electron with the quarks and gluons in the proton, as free quarks do not exist.

## Exercise 2: Massive quantum electrodynamics

Consider quantum electrodynamics with a single fermion  $\psi$  and a photon field  $A^\mu$ , that enter the Lagrangian density as follows

$$\mathcal{L}_{\text{QED}} = \bar{\psi}(i\not{D} - m)\psi - \frac{1}{4}F^{\mu\nu}F_{\mu\nu}$$

with

$$D_\mu = \partial_\mu + ieA_\mu, \quad F^{\mu\nu} = \partial^\mu A^\nu - \partial^\nu A^\mu.$$

We consider the U(1) gauge transformation

$$\psi(x) \rightarrow e^{i\alpha(x)}\psi(x), \quad A_\mu \rightarrow A_\mu - \frac{1}{e}\partial_\mu\alpha(x).$$

- (a) Show that  $\mathcal{L}_{\text{QED}}$  is invariant under the specified transformation.  
 (b) Show that on the other hand a mass term  $\mathcal{L}_{\text{mass}} = \frac{m_A^2}{2}A^\mu A_\mu$  is not invariant under the specified transformation.