

Einführung in Theoretische Teilchenphysik

Lecture: PD Dr. S. Gieseke – Exercises: Dr. Christoph Borchensky, Dr. Cody B Duncan

Exercise Sheet 2

<u>Hand-in Deadline</u>: Mo 01.11.21, 12:00. <u>Discussion</u>: Di 02.11.21, Mi 03.11.21.

1. [7 points] Lie Algebra structure of SU(N) and SO(N)

- (a) [1 point] Show that real $N \times N$ unimodular (i.e. det(R) = +1) and orthogonal $(R^{T}R = \mathbb{1}_{N})$ matrices (i.e. the elements of SO(N)) have N(N-1)/2 independent parameters.
- (b) [1 point] Show that complex $N \times N$ unimodular (i.e. det(U) = +1) and unitary $(R^{\dagger}R = \mathbb{1}_N)$ matrices (i.e. the elements of SU(N)) have $N^2 1$ independent parameters.
- (c) [2 points] Prove that any finite transformation $U \in SU(N)$ may be written as a combination of infinitesimal transformations $U(\theta) = \exp(i\theta_a J^a)$ where $a = 1, 2, ..., N^2 1$, and the basis elements of the Lie algebra J_a are: (i) hermitian, and (ii) traceless matrices, $\operatorname{Tr}(J_a) = 0 \forall a$

<u>Hint</u>: Take the limit $\theta^a \to 0 \forall a$. To check that $Tr(J_a) = 0$, recall that the identity $det(e^A) = e^{TrA}$.

We will from now on focus only on the SU(2) case.

(d) [1 points] Convince yourself that a generic SU(2) transformation can be parameterized as:

$$A = \begin{pmatrix} a_0 + ia_3 & a_2 + ia_1 \\ -a_2 + ia_1 & a_0 - ia_3 \end{pmatrix} = a_0 + i\vec{a} \cdot \vec{\sigma}$$

where $(a_0, \vec{a}) = (a_0, a_1, a_2, a_3)$ are real-valued parameters satisfying $\sum_i a_i^2 = 1$, and the vector $\vec{\sigma} = (\sigma_1, \sigma_2, \sigma_2)$ contains the Pauli matrices:

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad \sigma_2 = \begin{pmatrix} 0 & i \\ -i & 0 \end{pmatrix} \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

(e) [2 points] Prove that the parameterization of A above is equivalent to:

$$A = \exp\left[i\frac{\theta}{2}\left(\vec{\sigma}\cdot\vec{n}\right)\right],\,$$

where the rotation angle θ and the unit vector \vec{n} are related to (a_0, \vec{a}) . Make these relations explicit.

- 2. [6 points] A closer look at SO(3)
 - (a) [1 point] Given a generic SO(3) transformation represented in the form $R(\theta) = \exp(\theta_a K_a)$, show that the generators K_a are real and antisymmetric matrices.
 - (b) [1 point] If instead the transformation is represented as $R(\theta) = \exp(i\theta_a J_a)$, show that the generators J_a are imaginary and hermitian.

The following matrices implement independent rotations of angles θ_a , a = 1, 2, 3 around the x, y, z axis in 3-dimensional Euclidean space:

$$R_{x}(\theta_{1}) = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \theta_{1} & -\sin \theta_{1} \\ 0 & \sin \theta_{1} & \cos \theta_{1} \end{pmatrix}$$
$$R_{y}(\theta_{2}) = \begin{pmatrix} \cos \theta_{2} & 0 & \sin \theta_{2} \\ 0 & 1 & 0 \\ -\sin \theta_{2} & 0 & \cos \theta_{2} \end{pmatrix}$$
$$R_{z}(\theta_{3}) = \begin{pmatrix} \cos \theta_{3} & -\sin \theta_{3} & 0 \\ \sin \theta_{3} & \cos \theta_{3} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

(c) [2 points] Taking the limit $R(\delta\theta_a) = 1 + i\delta\theta_a J^a + \mathcal{O}(\theta^2)$, obtain the matrix form of $J_a, a = 1, 2, 3$. Check that your result is equivalent to the systematic recipe for computing the group generators:

$$J_a = \frac{1}{i} \left. \frac{\partial U(\theta_a)}{\partial \theta_a} \right|_{\theta_a = 0}$$

(d) [2 points] Using their explicit matrix form, check that J_a satisfy the commutation relation:

 $[J_a, J_b] = f_{ab}^c J_c.$

Determine the structure constants of the Lie algebra f_{ab}^c .

3. [7 points] Tensors under Lorentz transformations

Taking $x \to x'^{\mu} = \Lambda^{\mu}{}_{\nu}x^{\nu}$ and the orthogonality condition, $\Lambda^{\mu}{}_{\nu}\Lambda^{\rho}{}_{\sigma}g_{\mu\rho} = g_{\nu\sigma}$,

- (a) [1 point] Show that $a^{\mu}b_{\mu}$ is invariant under Lorentz transformations.
- (b) [2 points] Let $(\Lambda^{-1})^{\mu}{}_{\nu}$ be the inverse of $\Lambda_{\nu}{}^{\mu}$, meaning that

$$(\Lambda^{-1})^{\mu}{}_{\nu}\Lambda^{\nu}{}_{\rho} = \delta^{\mu}_{\rho}$$

Show that $(\Lambda^{-1})^{\mu}{}_{\nu} = \Lambda_{\nu}{}^{\mu}$.

- (c) [2 points] Show that $\partial_{\mu}{}' = (\Lambda^{-1})^{\nu}{}_{\mu}\partial_{\nu}$ where $\partial'_{\mu} = \frac{\partial}{\partial x'^{\mu}}$.
- (d) [2 points] Show that for a general 2-index tensor $M_{\mu\nu}$, that $M^{\mu}{}_{\nu}$ and $M_{\nu}{}^{\mu}$ are not necessarily equal. On the other other hand, show that $\delta^{\mu}{}_{\nu} = \delta_{\nu}{}^{\mu}$. Why is the ordering not relevant in the latter case?