
Prospects of Higgs Physics at the LHC

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Outline

I Introduction

- Higgs mechanism
- Supersymmetry

II Test of the Higgs mechanism at the LHC

- Higgs search at the LHC
- Higgs couplings to fermions, bosons
- Higgs boson quantum numbers
- Higgs self-couplings

III The Composite Higgs Boson

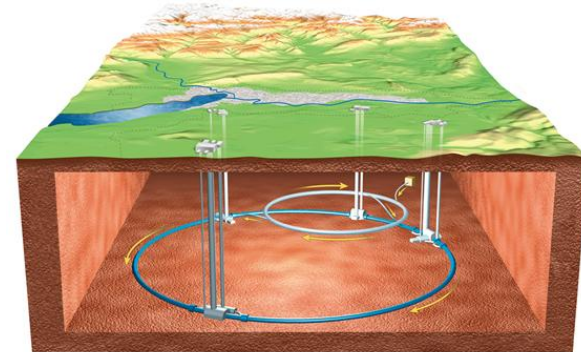
IV Conclusion

Research at the Large Hadron Collider LHC

Research at the LHC

Discoveries \Rightarrow

Understanding of matter and its interactions:

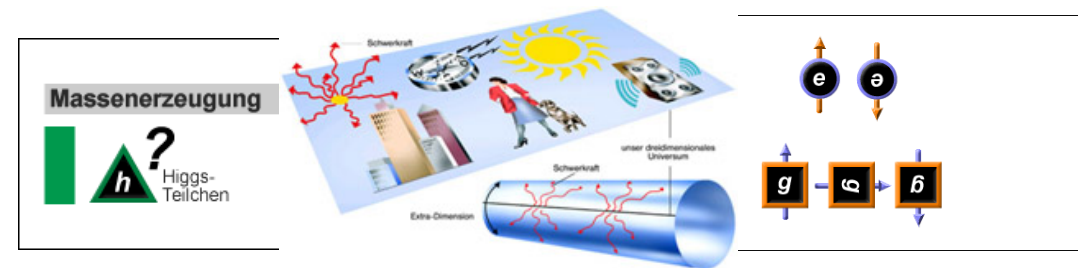


\rightarrow Verification of the Higgs mechanism

\rightarrow Search for supersymmetric particles

\rightarrow Search for extra dimensions

\vdots



The Standard Model of Particle Physics

Symmetry group $SU(3) \times SU(2)_L \times U(1)_Y$

I Particle content




Matter particles:

u	c	t	} Quarks
d	s	b	
ν_e	ν_μ	ν_τ	} Leptons
e	μ	τ	
1.	2.	3.	Family

Interaction particles:

γ	} Bosons
g	
Z	
W^\pm	

II Fundamental Forces

Electromagnetic	Photon	
Strong	Gluon	
Weak	W, Z	

III Higgs mechanism

Masses of the fundamental particles

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Masses of the fundamental particles



The Higgs mechanism

Why?

Explain the existence of massive particles consistently with the underlying symmetries of the Standard Model

Solution

Mechanism, which “breaks” the gauge symmetry in a specific way

Realisation

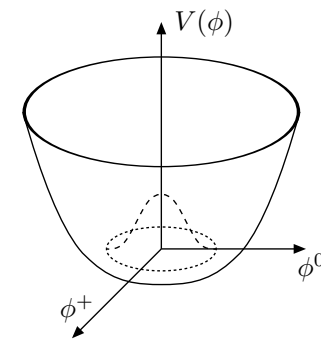
Higgs mechanism \rightsquigarrow Higgs particle



How it works

 Mass generation through spontaneous symmetry breaking (SSB)

- Self-interaction of the scalar field \rightsquigarrow ∞ number of degenerate ground states with non-vanishing field strength
- Choice of one ground state as the physical ground state \rightsquigarrow $SU(2)_L \times U(1)_Y$ symmetry hidden, $U(1)_{em}$ symmetry left: SSB
- Particles acquire mass through the interaction with the scalar field in the ground state
- Non-vanishing field strength $v = 246$ GeV \leftarrow typical minimax form of the Higgs potential



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Test of the
Higgs mechanism?
Accelerator
experiments!

Why Supersymmetry?

Standard Model: incomplete picture of the universe

- Common origin of all three forces of the Standard Model?
- How can we incorporate gravity?
- Candidate for Dark Matter? ...



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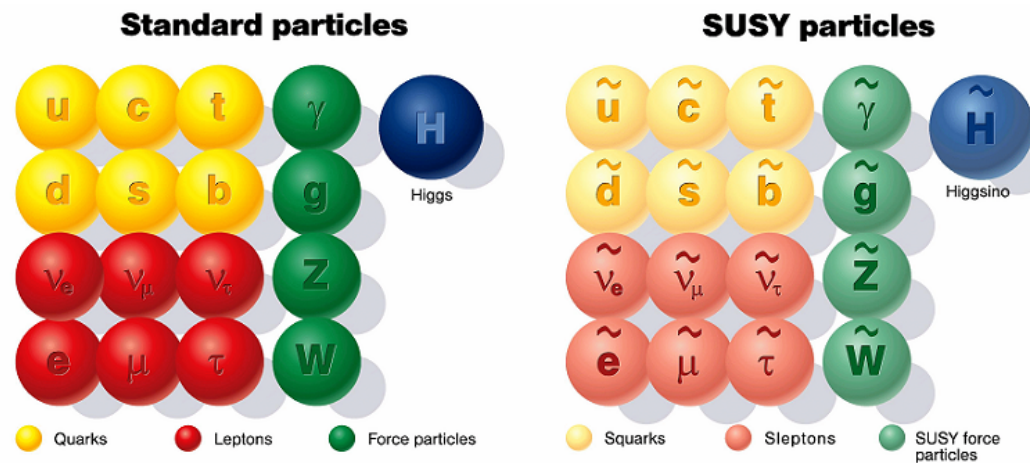
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Supersymmetry: provides answers

Fermions \leftrightarrow Bosons

Price: doubling of the particle spectrum

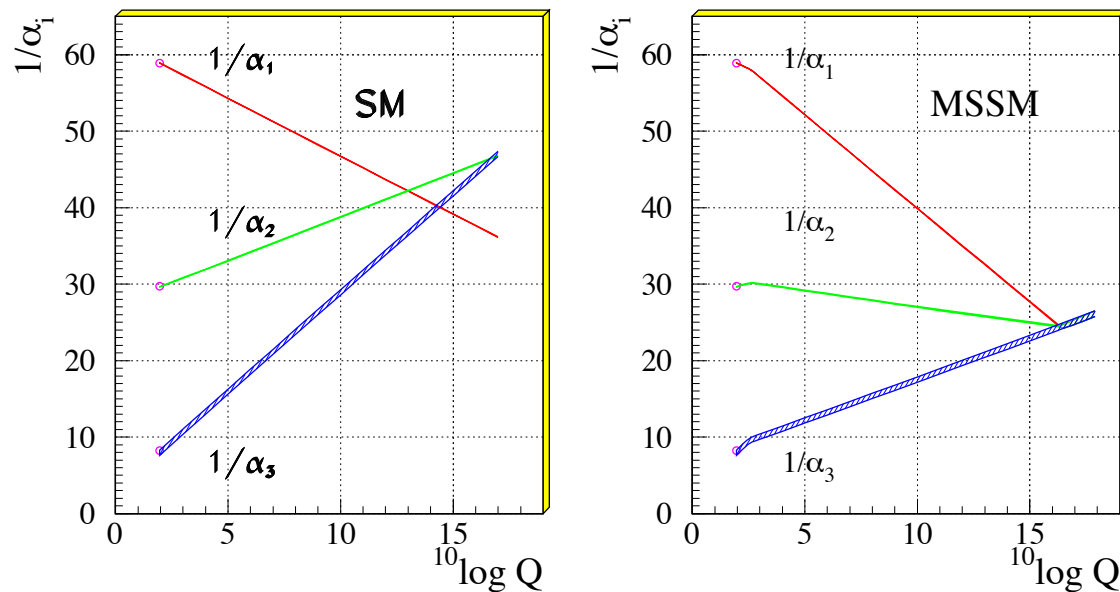


SUSY's answers

- ◇ **Unification of the coupling constants** electromagnetic - weak - strong

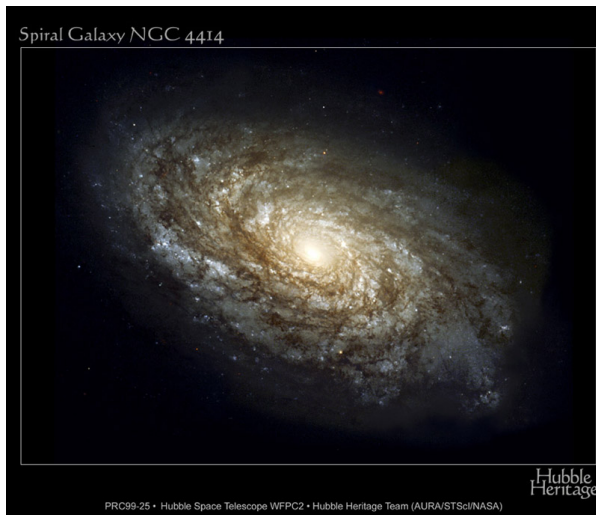
Amaldi, de Boer, Fürstenauf

Unification of the Coupling Constants in the SM and the minimal MSSM



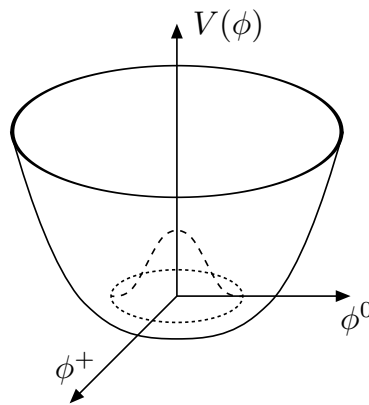
SUSY's Antworten

- ◇ **Unification of the couplings constants** elektromagnetic - weak - strong
- ◇ **Solution of the hierarchy problem** bosonic masses (\rightarrow Higgs mass) kept small in a natural way \leftarrow fermions \leftrightarrow bosons
- ◇ **Candidate for Cold Dark Matter** SUSY with R parity (DM $\sim 25\%$ of the universe)



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- ◇ **Candidate for cold Dark Matter** SUSY with R -Parity (DM $\sim 25\%$ of the universe)
- ◇ **Local supersymmetry** enforces gravity
- ◇ **Higgs mechanism** generated through radiative corrections



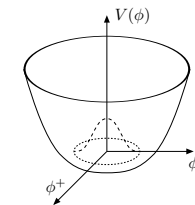
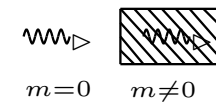
Experimental verification of the Higgs mechanism

Higgs mechanism:

Creation of particle masses without violating gauge principles

Test of the Higgs mechanism

- Discovery – m
- Interaction with a scalar Higgs with $v = 246 \text{ GeV} \neq 0$ $\rightsquigarrow g_{HXX} \sim m_X$
- Spin- and parity quantum numbers – J^{PC}
- EWSB requires Higgs potential – $\lambda_{HHH}, \lambda_{HHHH}$



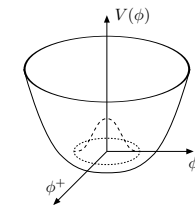
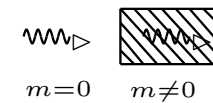
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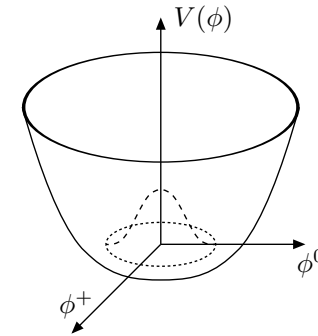


The SM Higgs Sector

The Higgs potential: [$v = 246$ GeV]

$$V(\Phi) = \lambda[\Phi^\dagger\Phi - \frac{v^2}{2}]^2 \quad \Phi = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v+H \end{pmatrix} \rightarrow$$

$$V(H) = \frac{1}{2}M_H^2 H^2 + \frac{M_H^2}{2v} H^3 + \frac{M_H^2}{8v^2} H^4$$



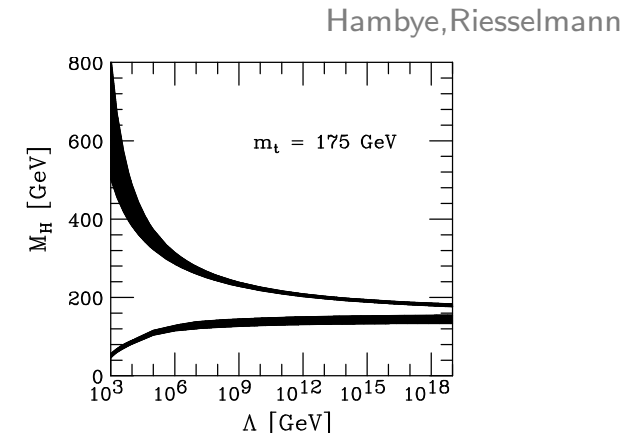
Higgs boson mass	$M_H = \sqrt{2\lambda}v$	
Gauge couplings	$g_{VVH} = \frac{2M_V^2}{v}$	
Yukawa couplings	$g_{ffH} = \frac{m_f}{v}$	
Trilinear coupling [units $\lambda_0 = 33.8$ GeV]	$\lambda_{HHH} = 3\frac{M_H^2}{M_Z^2}$	
Quartic coupling [units λ_0^2]	$\lambda_{HHHH} = 3\frac{M_H^2}{M_Z^4}$	

The only unknown parameter in the SM is the Higgs boson mass!

SM Higgs Mass Limits

- Triviality → upper bound Cabibbo,...; Sher;
Lindner; Hasenfratz,...;
- Vacuum stability → lower bound Lüscher, Weisz;
Hambye,...;...

$\Lambda = 1 \text{ TeV} :$ $55 \text{ GeV} \lesssim M_H \lesssim 700 \text{ GeV}$
 $\Lambda_{GUT} = 10^{16} \text{ GeV} :$ $130 \text{ GeV} \lesssim M_H \lesssim 190 \text{ GeV}$



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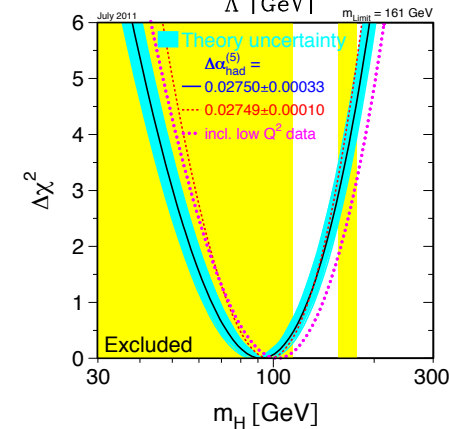
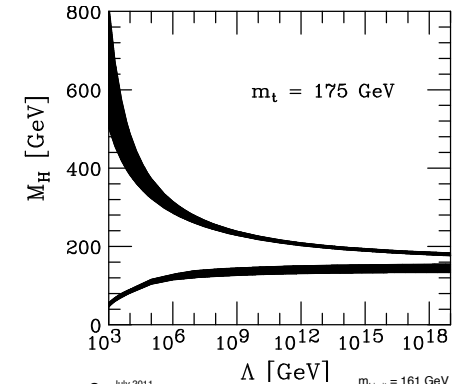
$$\Lambda_{GUT} = 10^{16} \text{ GeV} : \quad 130 \text{ GeV} \lesssim M_H \lesssim 190 \text{ GeV}$$

- Fits to electroweak precision data

$$M_H = 92_{-26}^{+34} \text{ GeV}, \quad M_H \lesssim 185 \text{ GeV} @ 95\% \text{ CL}$$

EWWG

Hambye, Riesselmann



LEP Coll.

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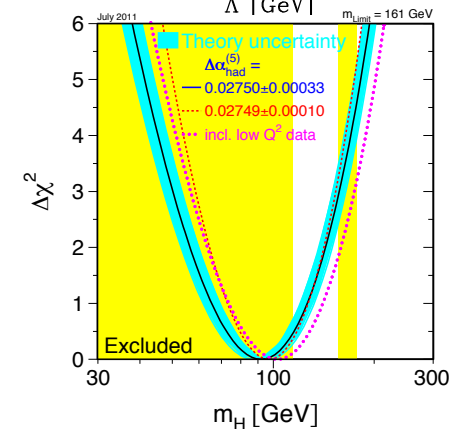
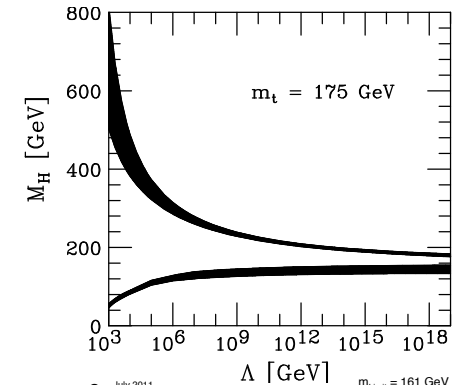
EWWG

- Direct search @ LEP: [$M_H = 115.3 \text{ GeV}$]

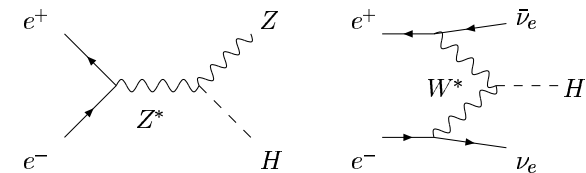
$$M_H > 114.4 \text{ GeV} @ 95\% \text{ CL}$$

LEP Coll.

Hambye, Riesselmann

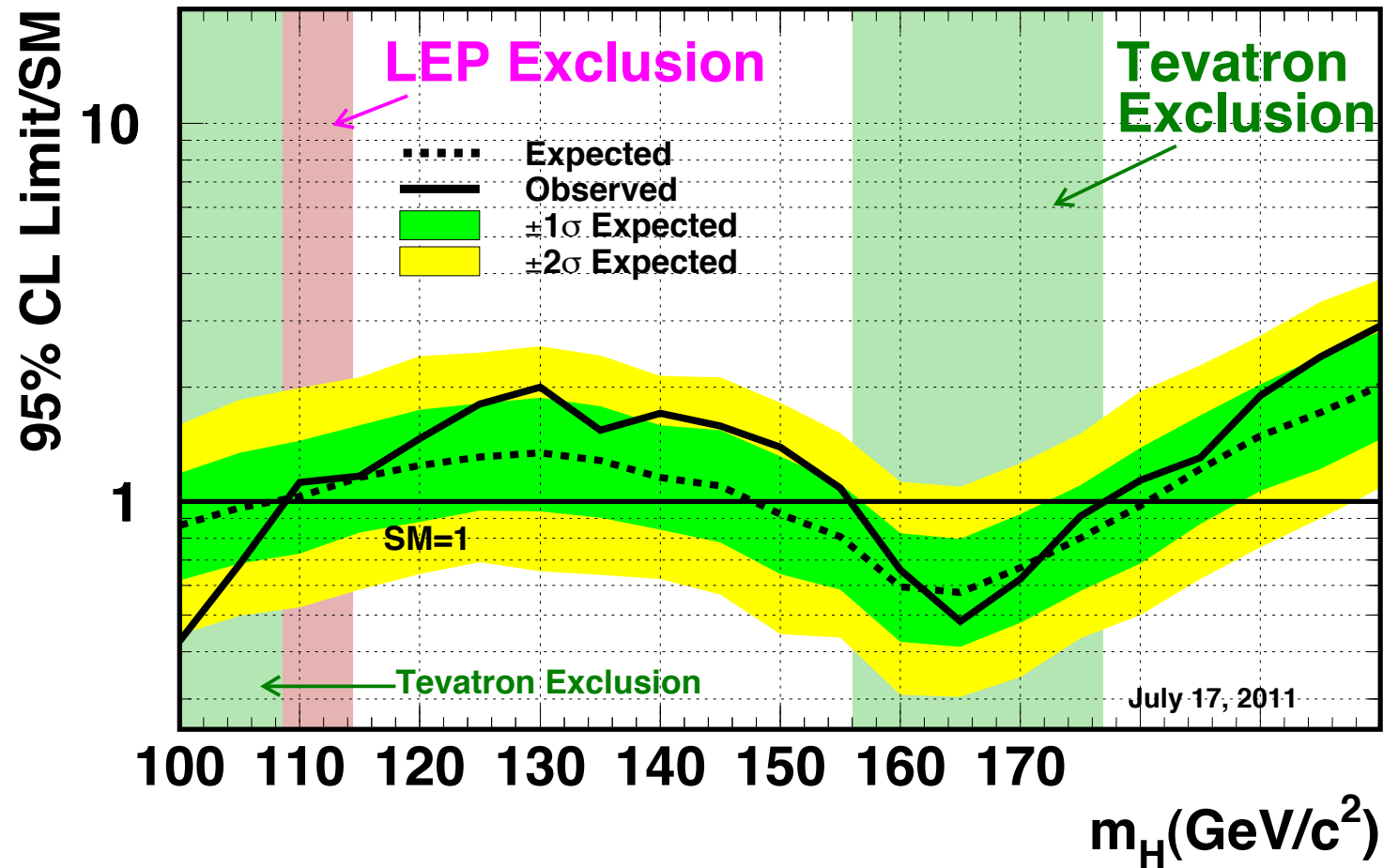


LEP Coll.



Tevatron Exclusion

Tevatron Run II Preliminary, $L \leq 8.6 \text{ fb}^{-1}$

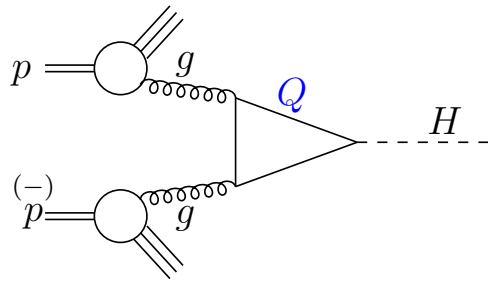


Higgs Suche am LHC

Higgsboson Produktion im SM

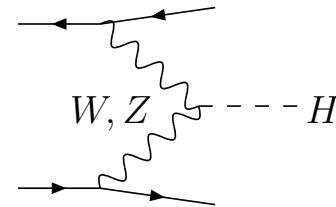
- Gluon Gluon Fusion

$$pp \rightarrow gg \rightarrow H$$



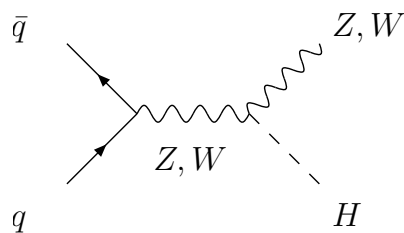
- W/Z Fusion

$$pp \rightarrow qq \rightarrow qq + WW/ZZ \rightarrow qq + H$$



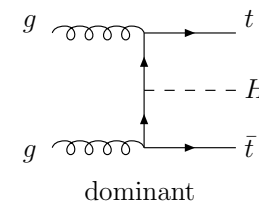
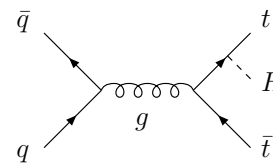
- Higgs-strahlung

$$pp \rightarrow W^*/Z^* \rightarrow W/Z + H$$



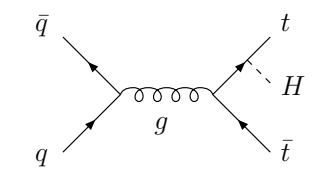
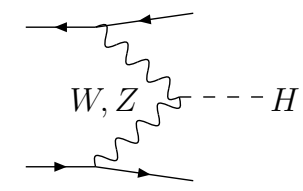
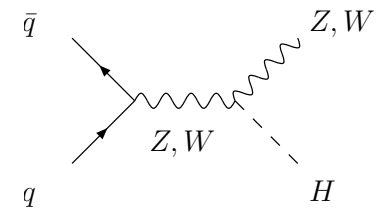
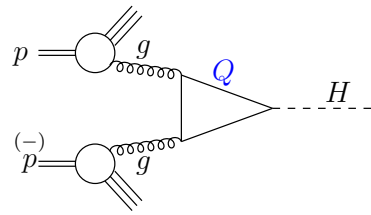
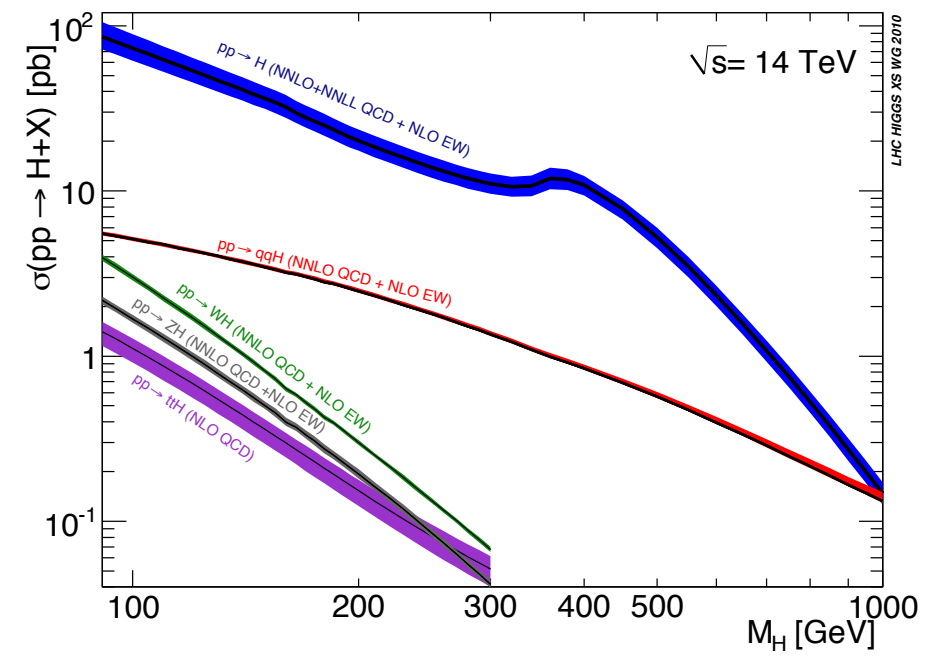
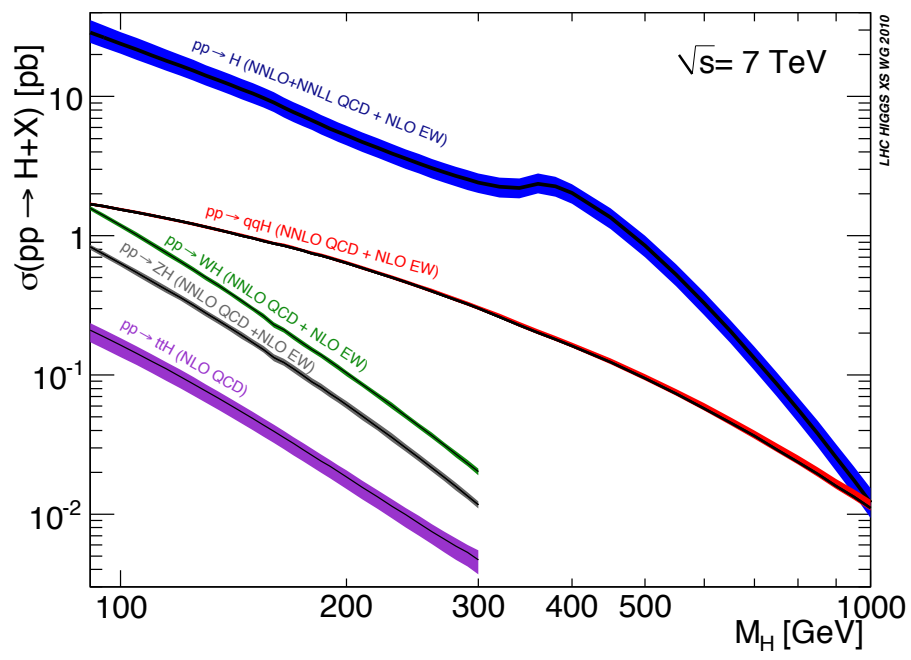
- Bremsstrahlung

$$pp \rightarrow t\bar{t} + H$$

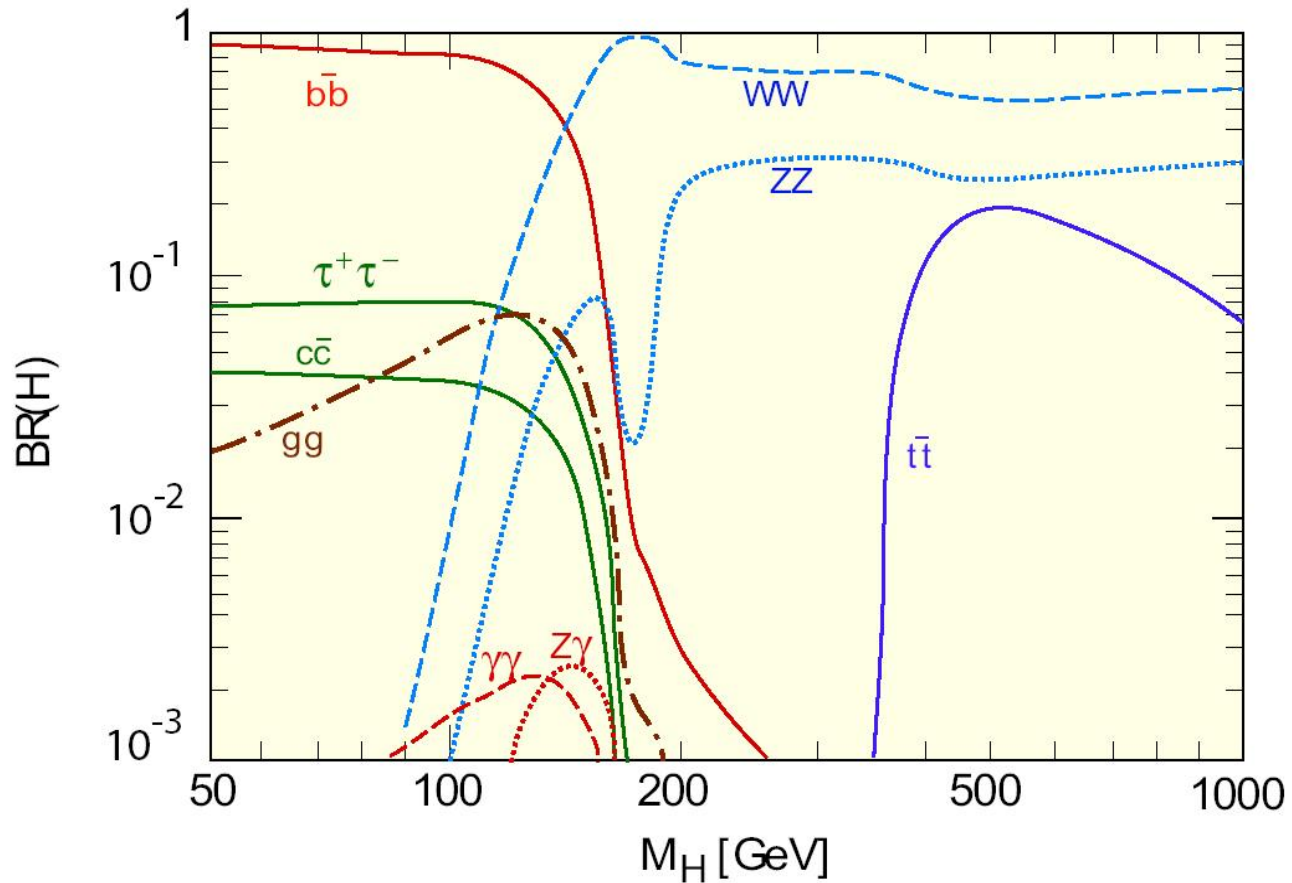


SM Higgs Boson Production at the LHC

LHC Higgs XS WG, arXiv:1101.0593



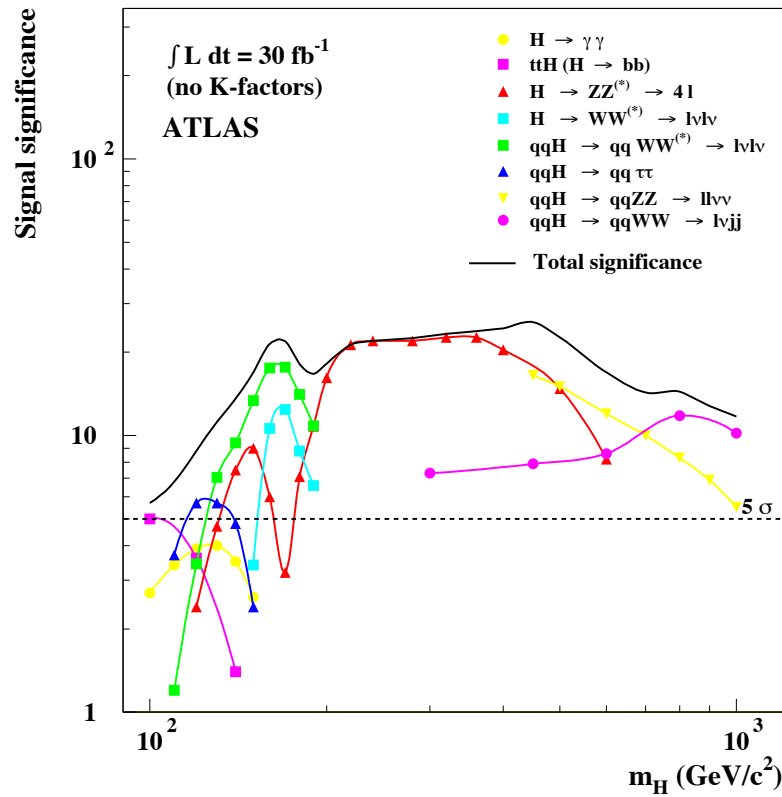
Higgsboson-Verzweigungsverhältnisse



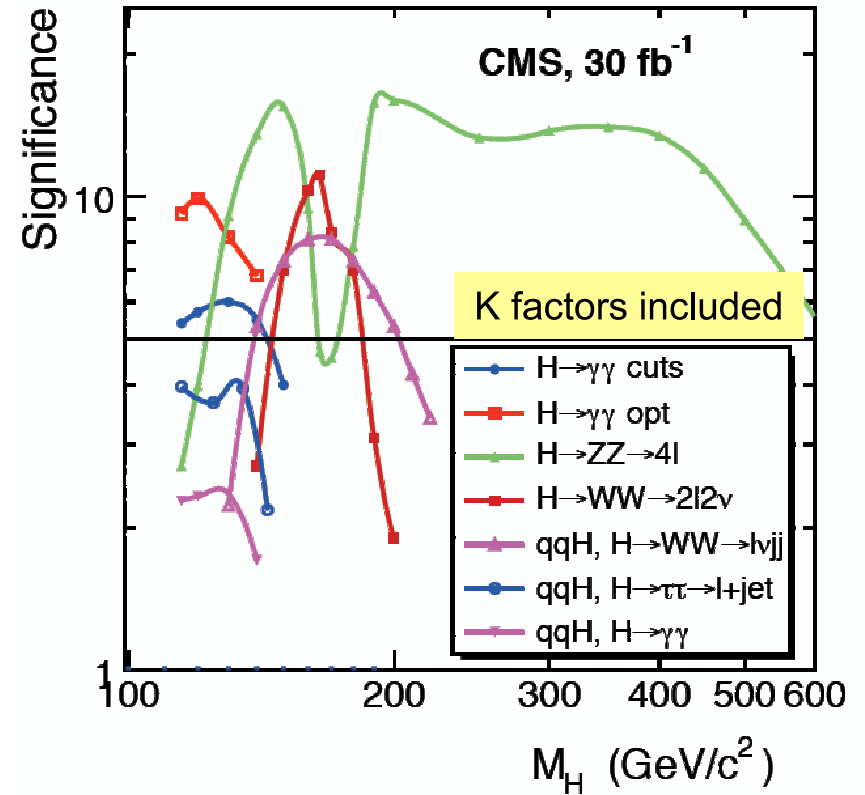
Higgs koppelt proportional zur Masse des Teilchen \rightsquigarrow Zerfälle in schwere Teilchen dominant (falls kinematisch erlaubt)

SM Higgsboson Suche am LHC

ATLAS

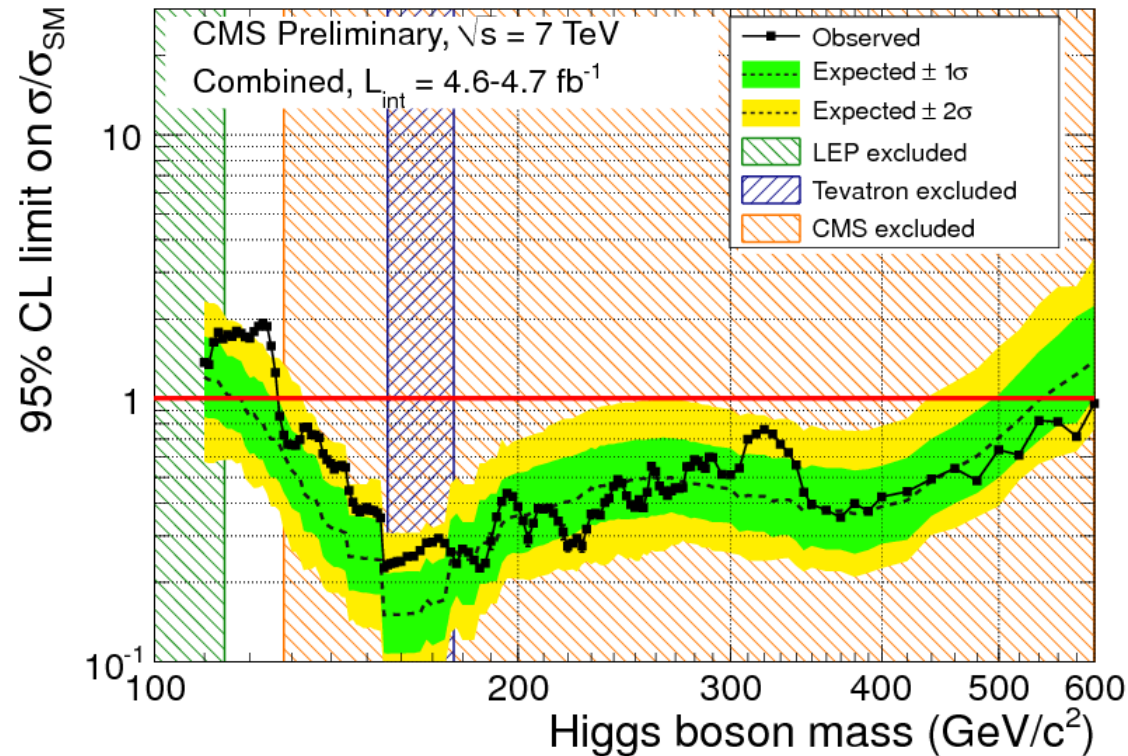


CMS



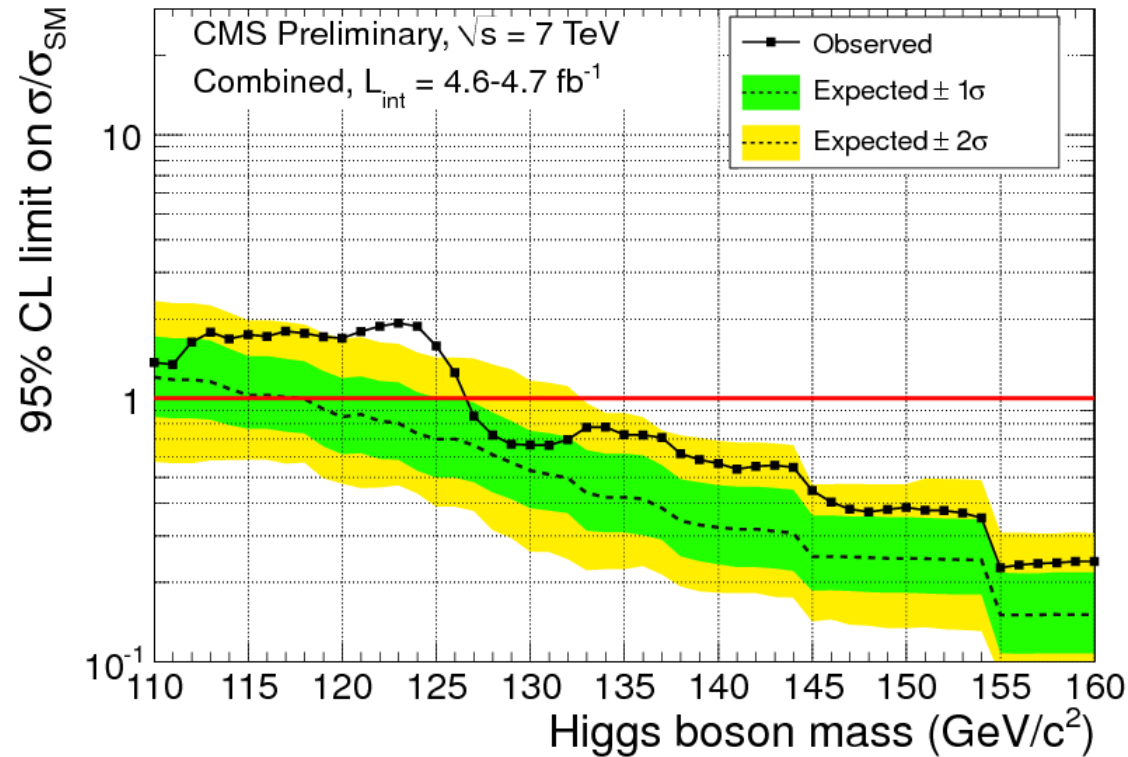
Genauigkeit: $\delta M_H / M_H \sim 10^{-3}$

CMS Ergebnisse



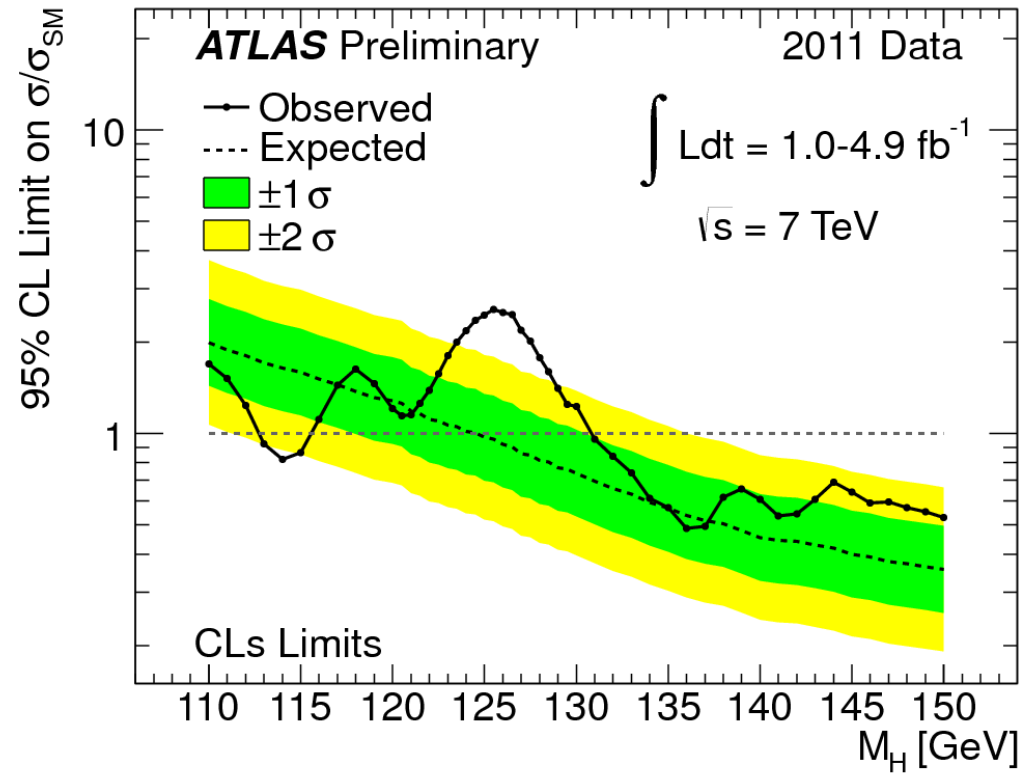
- Daten entsprechend 4.7 fb^{-1} integrierter Luminosität ausgewertet
- CMS kann den gesamten Massenbereich 114 GeV - 600 GeV untersuchen
- Kombination der Higgsuchen in den Higgszerfallskanälen in $WW^{(*)}$, $ZZ^{(*)}$, bb , $\tau\tau$, $\gamma\gamma$
- Ausschluß: 127-600 GeV auf 95% CL und 128-525 GeV auf 99% CL

CMS Ergebnisse



- Im Vergleich zum SM ein Überschuß an Daten in der Region 115-127 GeV konsistent in 5 unabh. Kanälen
- Mit den bisherigen Daten schwierig, zu unterscheiden zwischen den beiden Hypothesen Existenz bzw. Nicht-Existenz in der niedrigen Massenregion
- Überschuß ist kompatibel mit SM Higgs Hypothese bei etwa 124 GeV mit weniger als 2σ bei Einschluß von LEE; ohne LEE 2.6σ .

ATLAS Ergebnisse



- Higgsboson am wahrscheinlichsten in der Region 115-130 GeV
- Kombination der Higgsuchen in den Higgszerfallskanälen in $WW^{(*)} \rightarrow l\nu l\nu$, $ZZ^{(*)} \rightarrow 4l, \gamma\gamma$
- Ausschluß: 112.7-115.5 GeV, 131-453 GeV, außer 237-251 GeV auf 95% CL
(hohe Massenregion noch nicht untersucht)
- Überschuß ist kompatibel mit SM Higgs Hypothese bei etwa 126 GeV mit 2.6σ bei Einschluß von LEE; ohne LEE 3.6σ .

Minimal Supersymmetric Extension of the SM (MSSM)

MSSM Higgs sector – supersymmetry & anomaly-free theory \Rightarrow 2 complex Higgs doublets

$\xrightarrow{\text{EWSB}}$

neutral, CP-even h, H

neutral, CP-odd A

charged H^+, H^-

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Higgs masses

$$M_h \lesssim 140 \text{ GeV}$$
$$M_{A,H,H^\pm} \sim \mathcal{O}(v) \dots 1 \text{ TeV}$$

Ellis et al; Okada et al; Haber, Hempfling;
Hoang et al; Carena et al; Heinemeyer et al;
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Decoupling limit:

$$M_A \sim M_H \sim M_{H^\pm} \gg v$$
$$M_h \rightarrow \text{max. value, } \tan \beta \text{ fixed; } h \text{ SM-like}$$

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Modified couplings w/ respect to the SM: (decoupling limit Gunion, Haber)

Φ	$g_{\Phi u\bar{u}}$	$g_{\Phi d\bar{d}}$	$g_{\Phi VV}$
h	$c_\alpha / s_\beta \rightarrow 1$	$-s_\alpha / c_\beta \rightarrow 1$	$s_{\beta-\alpha} \rightarrow 1$
H	$s_\alpha / s_\beta \rightarrow 1 / \tan \beta$	$c_\alpha / c_\beta \rightarrow \tan \beta$	$c_{\beta-\alpha} \rightarrow 0$
A	$1 / \tan \beta$	$\tan \beta$	0

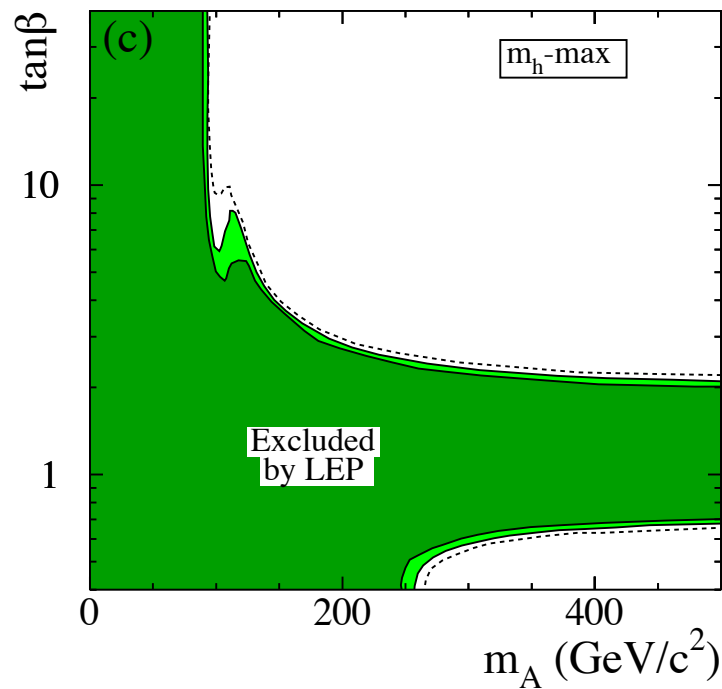
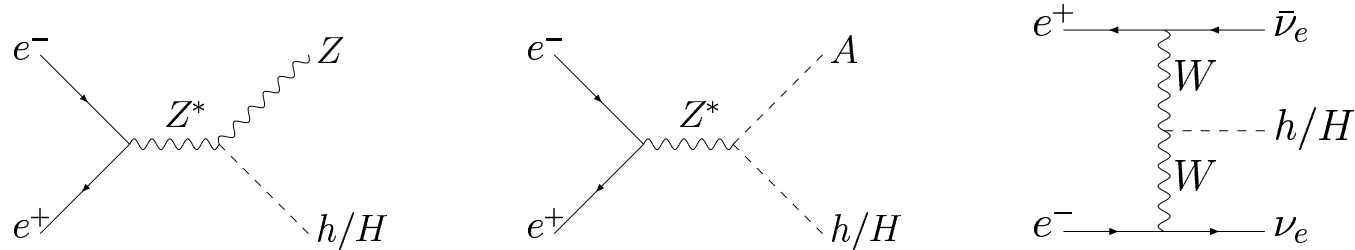
$$\tan \beta \uparrow \Rightarrow g_{\Phi uu} \downarrow$$

$$g_{\Phi dd} \uparrow$$

$$g_{\Phi VV}^{\text{MSSM}} \lesssim g_{\Phi VV}^{\text{SM}}$$

MSSM Higgs Mass Limits

▷ Direct Search at LEP $e^+e^- \rightarrow Z + h/H, A + h/H, \nu_e\bar{\nu}_e + h/H$



$$M_{h/H} \gtrsim 92.6 \text{ GeV}$$

$$M_A \gtrsim 93.4 \text{ GeV}$$

$$M_{H^\pm} > 78.6 \text{ GeV}$$

$0.6 < \tan \beta < 2.5$ excluded

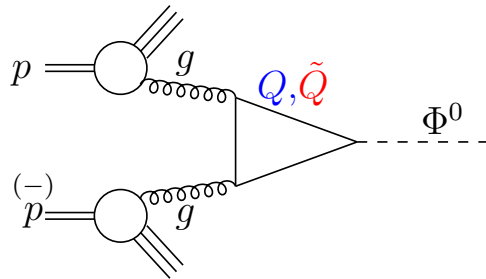
(only in this scenario, $m_t = 174.3 \text{ GeV}$!)

Higgs Search at the \mathcal{LHC}

Higgs boson production in SM/MSSM

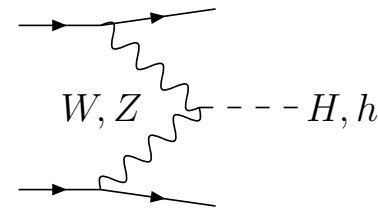
- Glue Gluon Fusion**

$$pp \rightarrow gg \rightarrow H^{SM} / h, H, A$$



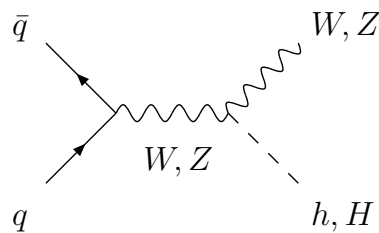
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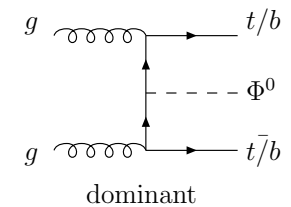
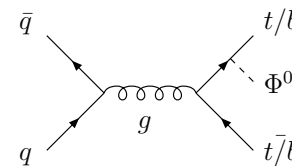
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$$pp \rightarrow W^*/Z^* \rightarrow W/Z + H^{SM} / h, H$$

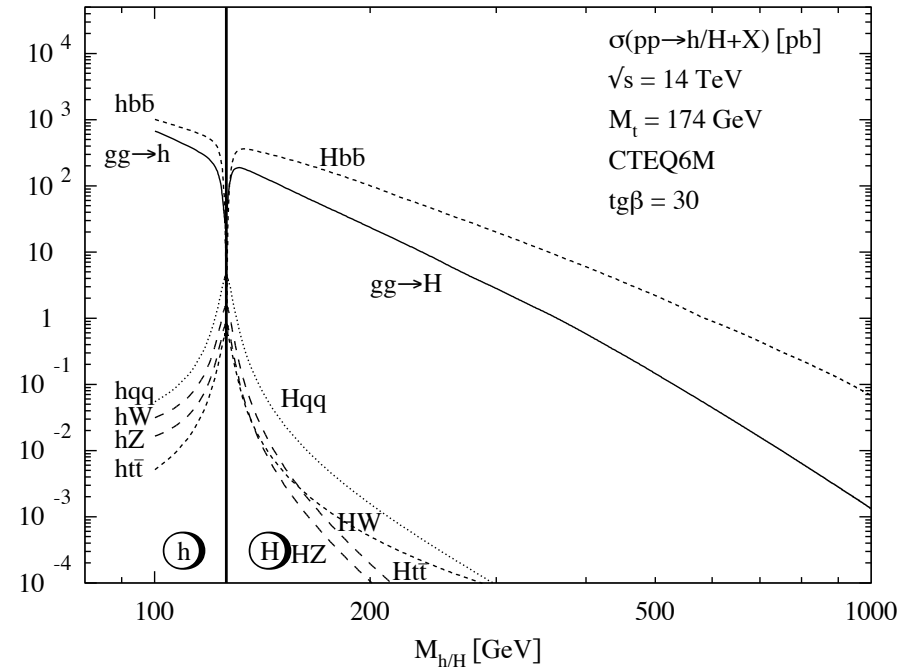
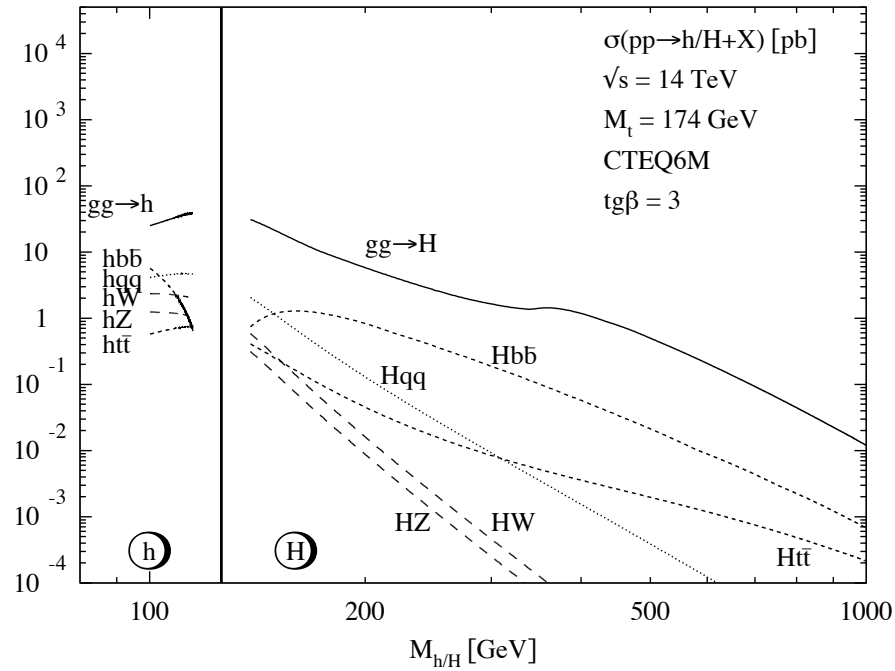


- Associated Production**

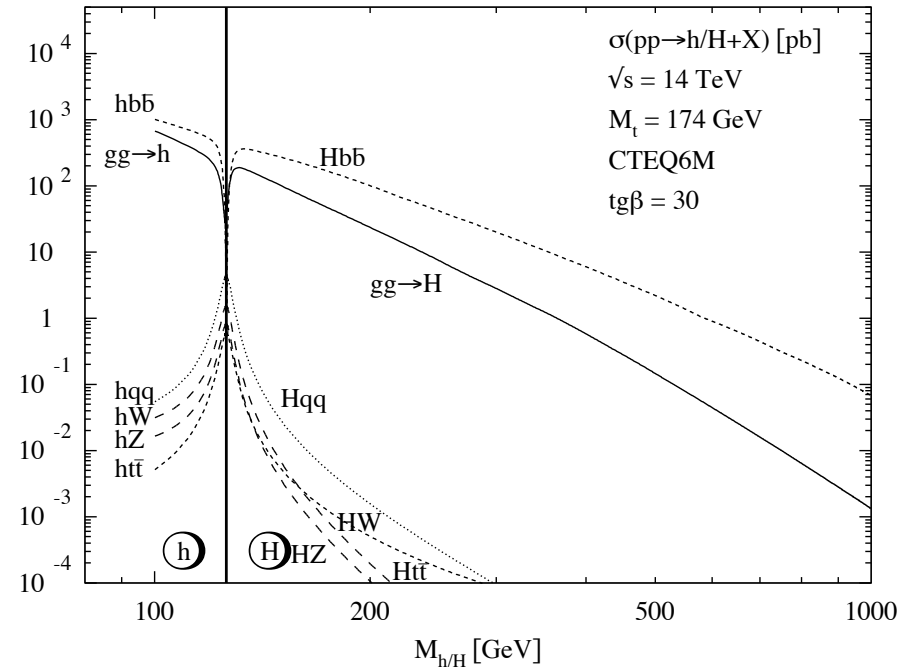
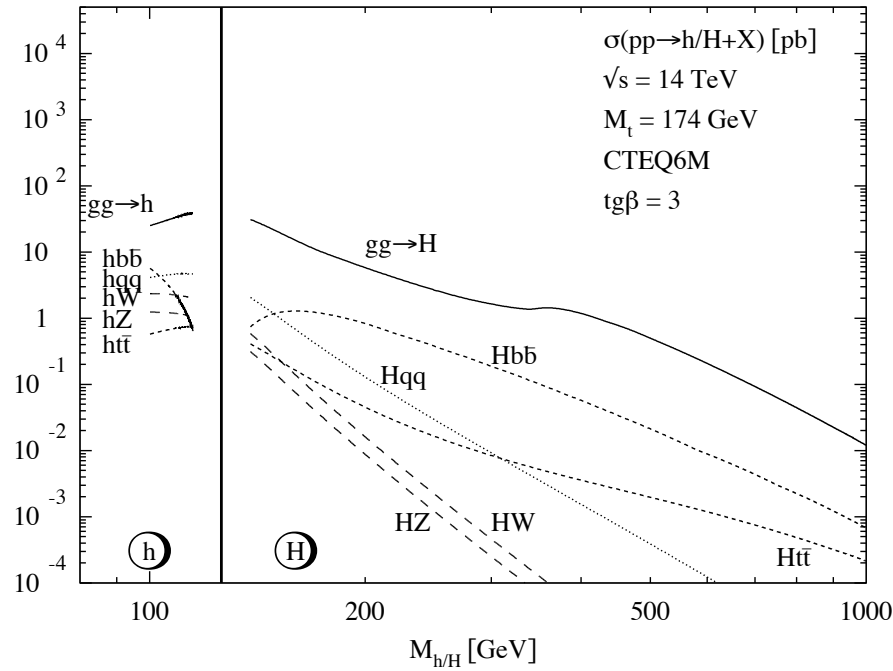
$$pp \rightarrow t\bar{t}/b\bar{b} + (H^{SM}) / h, H, A$$



MSSM Higgs Boson Production at the LHC



MSSM Higgs Boson Production at the LHC

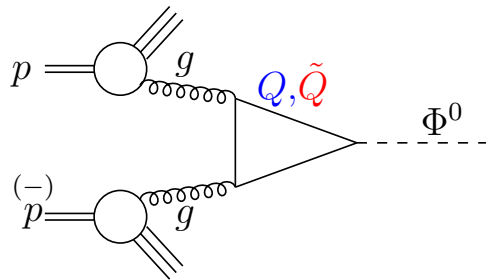


Reminder

$$\begin{aligned}
 \tan\beta \uparrow &\Rightarrow g_{\Phi uu} \downarrow \\
 &\qquad\qquad\qquad g_{\Phi dd} \uparrow \\
 g_{\Phi VV}^{MSSM} &\lesssim g_{\Phi VV}^{SM}
 \end{aligned}$$

Higgs Boson Production in gluon fusion

(i) **Dominant: Gluon Fusion** $pp \rightarrow gg \rightarrow H^{SM} / h, H, A$ (small & moderate $\tan \beta$)



Georgi et al; Gamberini et al

QCD corrections to top & bottom loops

- ▷ NLO (SM, MSSM): **increase σ by $\sim 10\text{...}100\%$**
[moderate for large $\tan \beta \leftarrow b$ -loop]
- ▷ SM; $\text{tg}\beta \lesssim 5$: limit $M_\Phi \ll m_t$ - **approximation $\sim 20\text{-}30\%$**
- ▷ NNLO @ $M_\Phi \ll m_t \Rightarrow$ **further increase by 20-30%**
- ▷ Estimate of NNNLO effects @ $M_\Phi \ll m_t \rightsquigarrow$ **scale stabilisation**
scale dependence: **$\Delta \lesssim 10 - 15\%$**
- ▷ NNLL resummation: **$+ \sim 10\%$**
- ▷ resummation of soft gluons @ N³LL and of π^2 enhanced terms

Spira, Djouadi, Graudenz, Zerwas
Dawson; Kauffman, Schaffer

Krämer, Laenen, Spira

Harlander, Kilgore
Anastasiou, Melnikov
Ravindran, Smith, van Neerven
Moch, Vogt
Ravindran

Catani, de Florian, Grazzini, Nason

Ahrens, Neubert, Becher, Yang

Higgs Boson Production in gluon fusion

Corrections to top & bottom loops

- ▷ NNLO mass effects (t loops) Harlander, Ozeren; Pak, Rogal, Steinhauser; Marzani et al.
for $M_H \lesssim 300 \text{ GeV} \Rightarrow \mathcal{O}(0.5\%)$
- ▷ NLO electroweak corrections $\sim -4\% - 6\%$ (SM) Aglietti et al.; Degrassi, Maltoni; Actis et al
- ▷ mixed QCD and EW corrections Anastasiou, Boughezal, Petriello

NLO corrections to squark loops

- ▷ in the heavy mass limit Dawson, Djouadi, Spira
- ▷ full SUSY-QCD corrections in heavy mass limit Harlander, Steinhauser; Harlander, Hofmann; Degrassi, Slavich '11
- ▷ bottom/sbottom contributions Degrassi, Slavich '11
asymptotic expansion in $\tilde{M} \gg m_b, M_\phi$ Harlander, Hofmann, Mantler '11

$m_{\tilde{Q}} \lesssim 400 \text{ GeV}$:

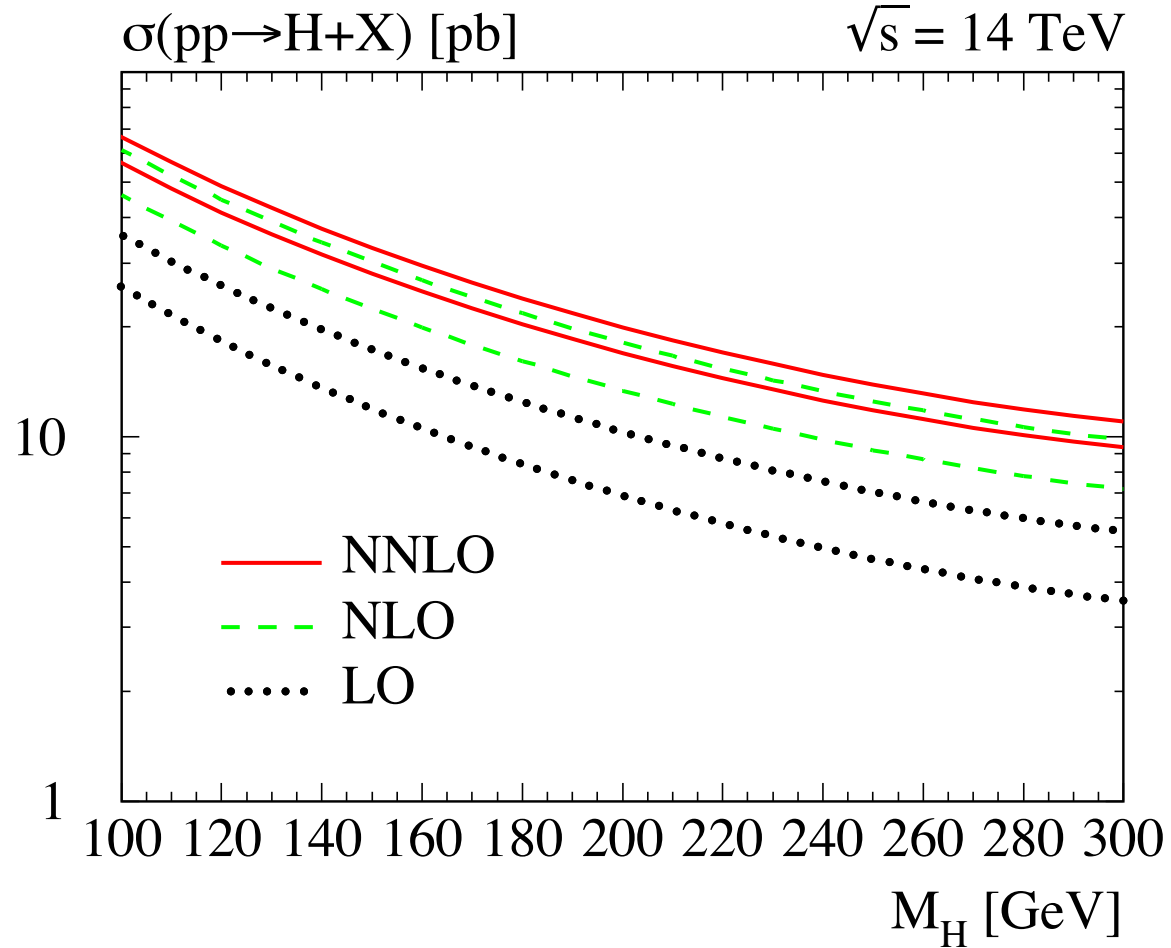
- ▷ NLO squark mass effects $\sim 15\%$ MMM, Spira; Anastasiou, Beerli, Bucherer, Daleo, Kunszt; Aglietti, Bonciani, Degrassi, Vicini
- ▷ full NLO SUSY QCD calculation Anastasiou, Beerli, Daleo; MMM, Rzehak, Spira

NNLO SUSY-QCD corrections from t/\tilde{t} sector

Pak, Steinhauser, Zerf '10

$gg \rightarrow H$ \mathcal{NNLO}

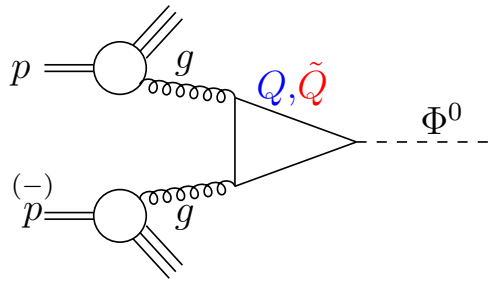
Harlander, Kilgore



gg → H, h at leading order

Lowest order - 1 loop

Georgi,...;Gamberini,...



$$\sigma(pp \rightarrow \Phi + X) = \sigma_0^\Phi \tau_\Phi \frac{d\mathcal{L}^{gg}}{d\tau_\Phi}$$

$$\sigma_0^{h/H} = \frac{G_F \alpha_S^2(\mu_R)}{288\sqrt{2}\pi} \left| \sum_Q g_Q^{h/H} F_Q^{h/H}(\tau_Q) + \sum_{\tilde{Q}} g_{\tilde{Q}}^{h/H} F_{\tilde{Q}}^{h/H}(\tau_{\tilde{Q}}) \right|^2$$

$$\sigma_0^A = \frac{G_F \alpha_s^2}{128\sqrt{2}\pi} \left| \sum_Q g_Q^A F_Q^A(\tau_Q) \right|^2$$

$$F_Q^{h/H}(\tau_Q) = \frac{3}{2} \tau_Q \left[1 + (1 - \tau_Q) f(\tau_Q) \right]$$

$$F_Q^A(\tau_Q) = \tau_Q f(\tau_Q)$$

$$F_{\tilde{Q}}^{h/H}(\tau_{\tilde{Q}}) = -\frac{3}{4} \tau_{\tilde{Q}} \left[1 - \tau_{\tilde{Q}} f(\tau_{\tilde{Q}}) \right]$$

$$f(\tau) = \begin{cases} \arcsin^2 \frac{1}{\sqrt{\tau}} & \tau \geq 1 \\ -\frac{1}{4} \left[\log \left(\frac{1+\sqrt{1-\tau}}{1-\sqrt{1-\tau}} \right) - i\pi \right]^2 & \tau < 1 \end{cases}$$

$$\tau_\Phi = \frac{M_\Phi^2}{s}, \quad \tau_{Q,\tilde{Q}} = \frac{4m_{Q,\tilde{Q}}^2}{M_\Phi^2}$$

First Step: QCD corrections

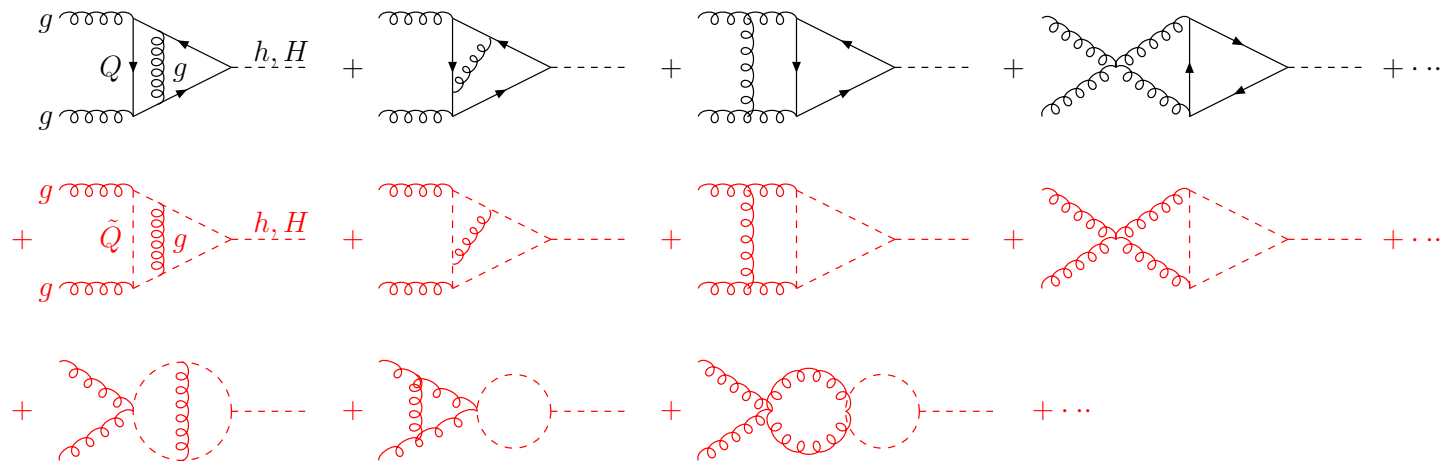
$$\Delta\hat{\sigma}_{ij} = \sigma_0 \left\{ C_{ij} \delta(1 - \hat{\tau}) + D_{ij} \Theta(1 - \hat{\tau}) \right\} \frac{\alpha_s}{\pi}$$

$$\hat{\tau} = \frac{M_\Phi^2}{\hat{s}}$$

↑
virtual+soft
corrections

↑
real corrections

Virtual corrections [2 loops, first step: no gluino contributions]



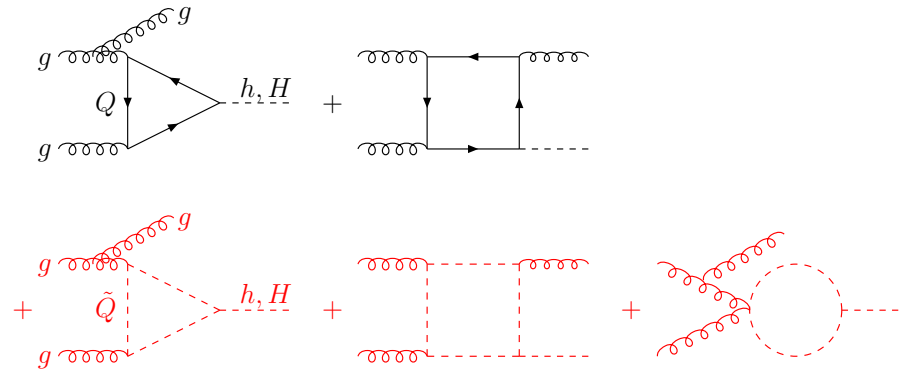
UV-,IR-,Coll-singularities in $n = 4 - 2\epsilon$ dimensions.

Real Corrections

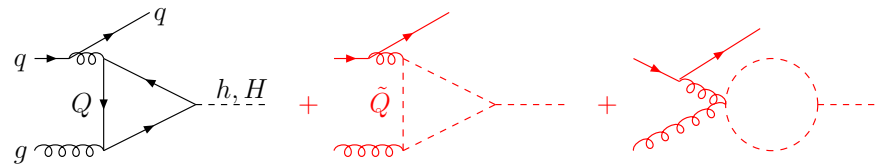
After renormalization: IR & coll. singularities \rightsquigarrow real corrections have to be added.

3 incoherent processes:

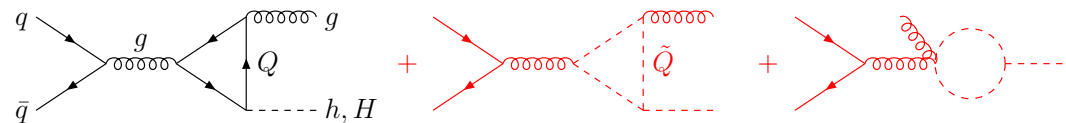
$gg \rightarrow Hg$:



$gq \rightarrow Hq$:



$q\bar{q} \rightarrow Hg$:



Phase space integration in $n = 4 - 2\epsilon$ dimensions \rightsquigarrow IR, Coll. singularities: poles in ϵ

Result

- α_S : $\overline{\text{MS}}$ scheme, 5 active flavours - μ =Ren. scale, Q =Fact. scale, $\mu^2 = Q^2 = M_\phi^2$

$$\sigma(pp \rightarrow \phi + X) = \sigma_0^\phi [1 + C^\phi \frac{\alpha_S}{\pi}] \tau_\phi \frac{d\mathcal{L}_{gg}}{d\tau_\phi} + \Delta\sigma_{gg}^\phi + \Delta\sigma_{gq}^\phi + \Delta\sigma_{q\bar{q}}^\phi$$

$$C^\phi(\tau_Q, \tau_{\tilde{Q}}) = \pi^2 + C_1^\phi(\tau_Q, \tau_{\tilde{Q}}) + \frac{33-2N_F}{6} \log \frac{\mu^2}{M_\phi^2}$$

$$\Delta\sigma_{gg}^\phi = \int_{\tau_\phi}^1 d\tau \frac{d\mathcal{L}_{gg}}{d\tau} \frac{\alpha_S}{\pi} \sigma_0^\phi \left\{ -\hat{\tau} P_{gg}(\hat{\tau}) \log \frac{Q^2}{\hat{s}} + d_{gg}^\phi(\hat{\tau}, \tau_Q, \tau_{\tilde{Q}}) \right. \\ \left. + 12 \left[\left(\frac{\log(1-\hat{\tau})}{1-\hat{\tau}} \right)_+ - \hat{\tau} [2 - \hat{\tau}(1-\hat{\tau})] \log(1-\hat{\tau}) \right] \right\}$$

$$\Delta\sigma_{gq}^\phi = \int_{\tau_\phi}^1 d\tau \sum_{q, \bar{q}} \frac{d\mathcal{L}_{gq}}{d\tau} \frac{\alpha_S}{\pi} \sigma_0^\phi \left\{ -\frac{\hat{\tau}}{2} P_{gq}(\hat{\tau}) \left[\log \frac{Q^2}{\hat{s}(1-\hat{\tau})^2} \right] d_{gq}^\phi(\hat{\tau}, \tau_Q, \tau_{\tilde{Q}}) \right\}$$

$$\Delta\sigma_{q\bar{q}}^\phi = \int_{\tau_\phi}^1 d\tau \sum_q \frac{d\mathcal{L}_{q\bar{q}}}{d\tau} \frac{\alpha_S}{\pi} \sigma_0^\phi d_{q\bar{q}}^\phi(\hat{\tau}, \tau_Q, \tau_{\tilde{Q}})$$

$$- \tau_{Q, \tilde{Q}} = \frac{4m_{Q, \tilde{Q}}^2}{M_\Phi^2}, \quad \hat{\tau} = \frac{m_\phi^2}{\hat{s}}$$

The Scenario

The gluophobic Higgs scenario [$m_t = 174.3$ GeV]

Carena, Heinemeyer, Wagner, Weiglein

$$M_{SUSY} = 350 \text{ GeV}, \mu = M_2 = 300 \text{ GeV}, X_t = -770 \text{ GeV}, A_b = A_t, m_{\tilde{g}} = 500 \text{ GeV}$$

$$\tan \beta = 3$$

$$m_{\tilde{t}_1} = 156 \text{ GeV} \quad m_{\tilde{t}_2} = 517 \text{ GeV}$$

$$m_{\tilde{b}_1} = 346 \text{ GeV} \quad m_{\tilde{b}_2} = 358 \text{ GeV}$$

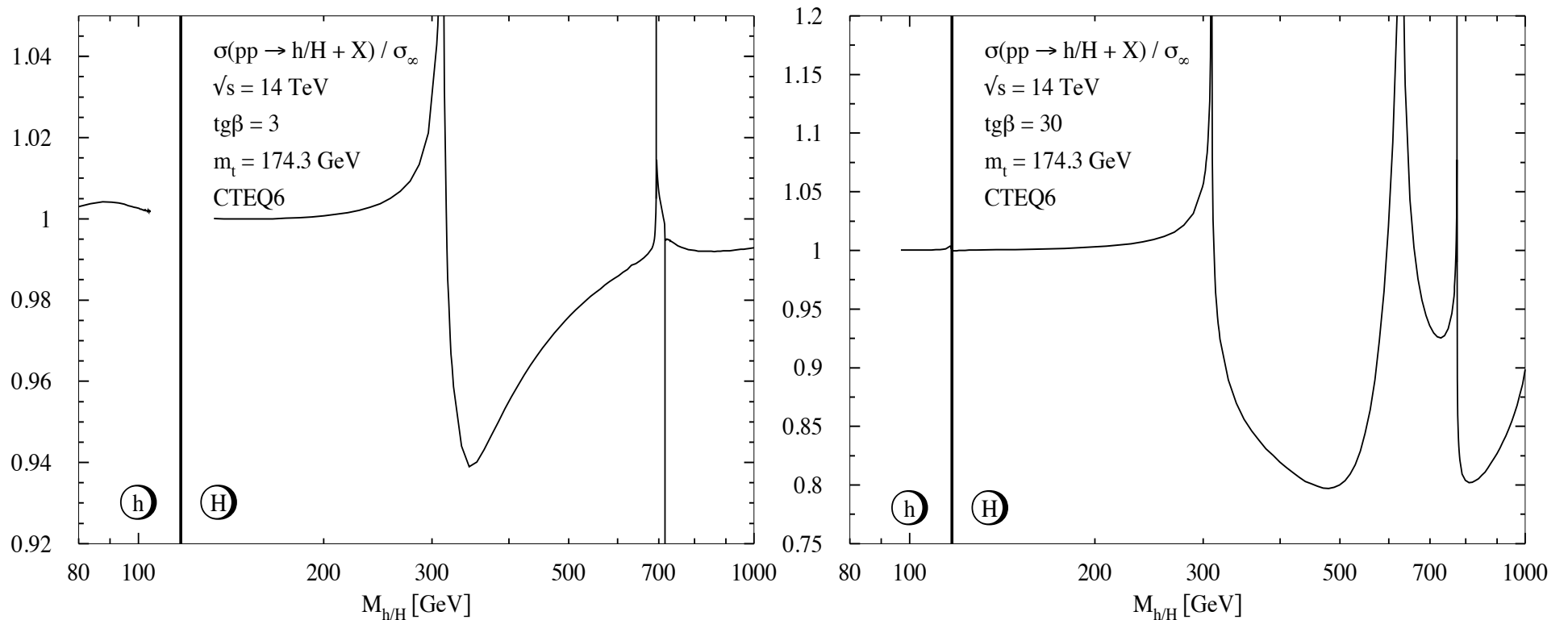
$$\tan \beta = 30$$

$$m_{\tilde{t}_1} = 155 \text{ GeV} \quad m_{\tilde{t}_2} = 516 \text{ GeV}$$

$$m_{\tilde{b}_1} = 314 \text{ GeV} \quad m_{\tilde{b}_2} = 388 \text{ GeV}$$

NLO cross section →

σ_{NLO} w/ full squark mass dependence / σ_{NLO} in the heavy squark limit



$\sigma(pp \rightarrow h/H + X) / \sigma_\infty$ up to 20%

Kinks, bumps, spikes: $\tilde{t}_1\tilde{t}_1, \tilde{b}_1\tilde{b}_1, \tilde{b}_2\tilde{b}_2$ thresholds in consecutive order with rising Higgs mass.

Coulomb singularities

$\tilde{Q}\bar{Q}$ thresholds: Formation of 0^{++} states \rightsquigarrow Coulomb singularities

Singular behaviour can be derived from the Sommerfeld rescattering corrections \rightsquigarrow

At each specific $\tilde{Q}_0\bar{Q}_0$ threshold:

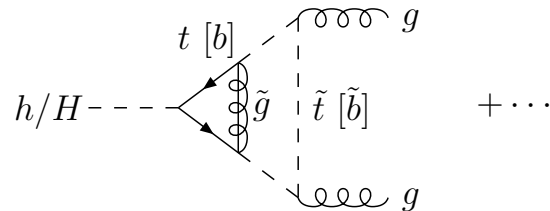
$$C_1(\tau_Q, \tau_{\tilde{Q}}) \rightarrow \text{Re} \left\{ \frac{g_{\tilde{Q}_0}^{\Phi} \tilde{F}(\tilde{Q}_0) \frac{16\pi^2}{3(\pi^2-4)} \left[-\ln(\tau_{\tilde{Q}_0}^{-1}-1) + i\pi + \text{const} \right]}{\sum_Q g_Q^{\Phi} F(\tau_Q) + \sum_{\tilde{Q}} g_{\tilde{Q}}^{\Phi} \tilde{F}(\tau_{\tilde{Q}})} \right\}$$

Agrees quantitatively with numerical results.

Genuine SUSY-QCD corrections

- Limit heavy SUSY masses $\rightarrow \mathcal{O}(10\%)$

Harlander, Steinhauser, Hofmann



Anastasiou, Beerli, Daleo
MMM, Rzehak, Spira

- Small α_{eff} scenario [modified]

$$\tan \beta = 30$$

$$M_{\tilde{Q}} = 800 \text{ GeV}$$

$$M_{\tilde{g}} = 1000 \text{ GeV} \quad \leftarrow$$

$$M_2 = 500 \text{ GeV}$$

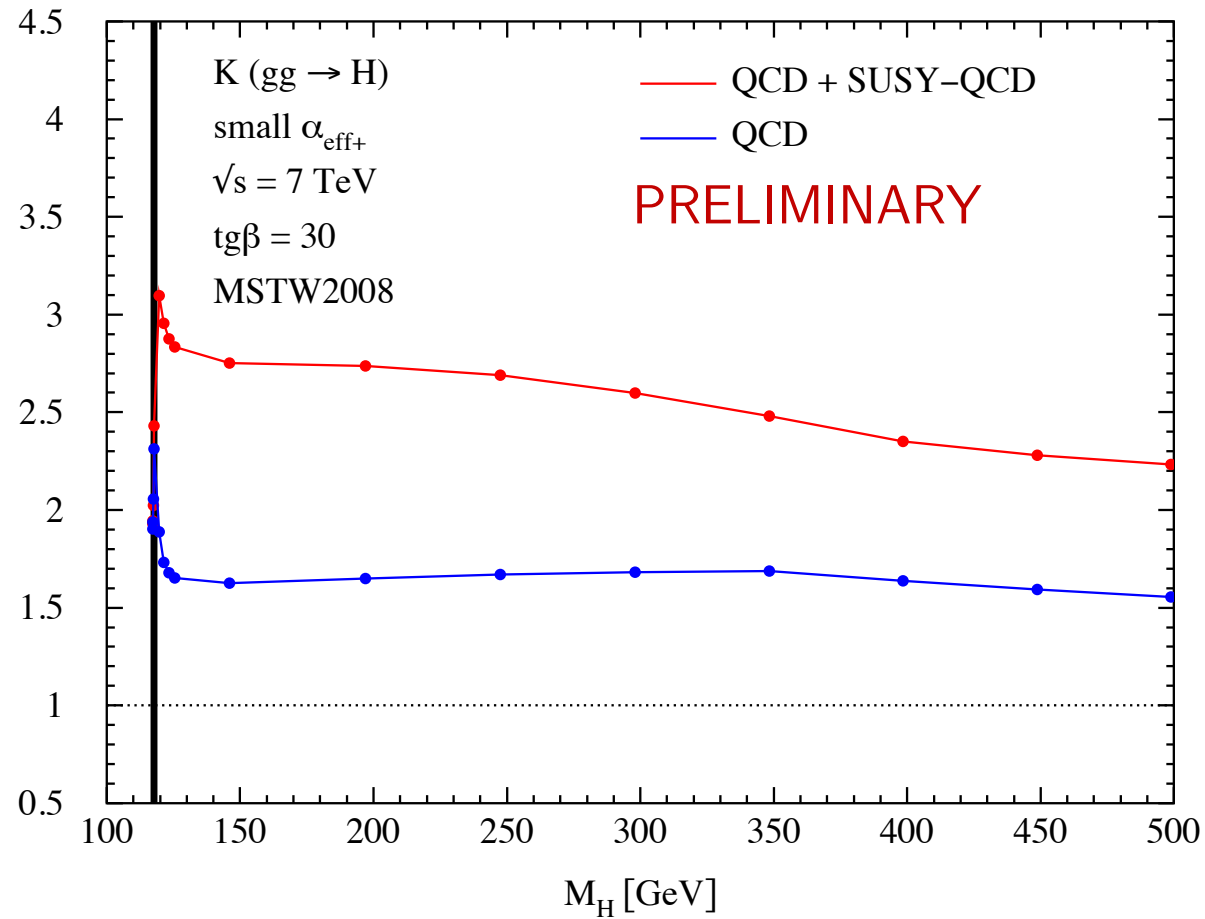
$$A_b = A_t = -1.133 \text{ TeV}$$

$$\mu = 2 \text{ TeV}$$

$$m_{\tilde{t}_1} = 679 \text{ GeV} \quad m_{\tilde{t}_2} = 935 \text{ GeV}$$

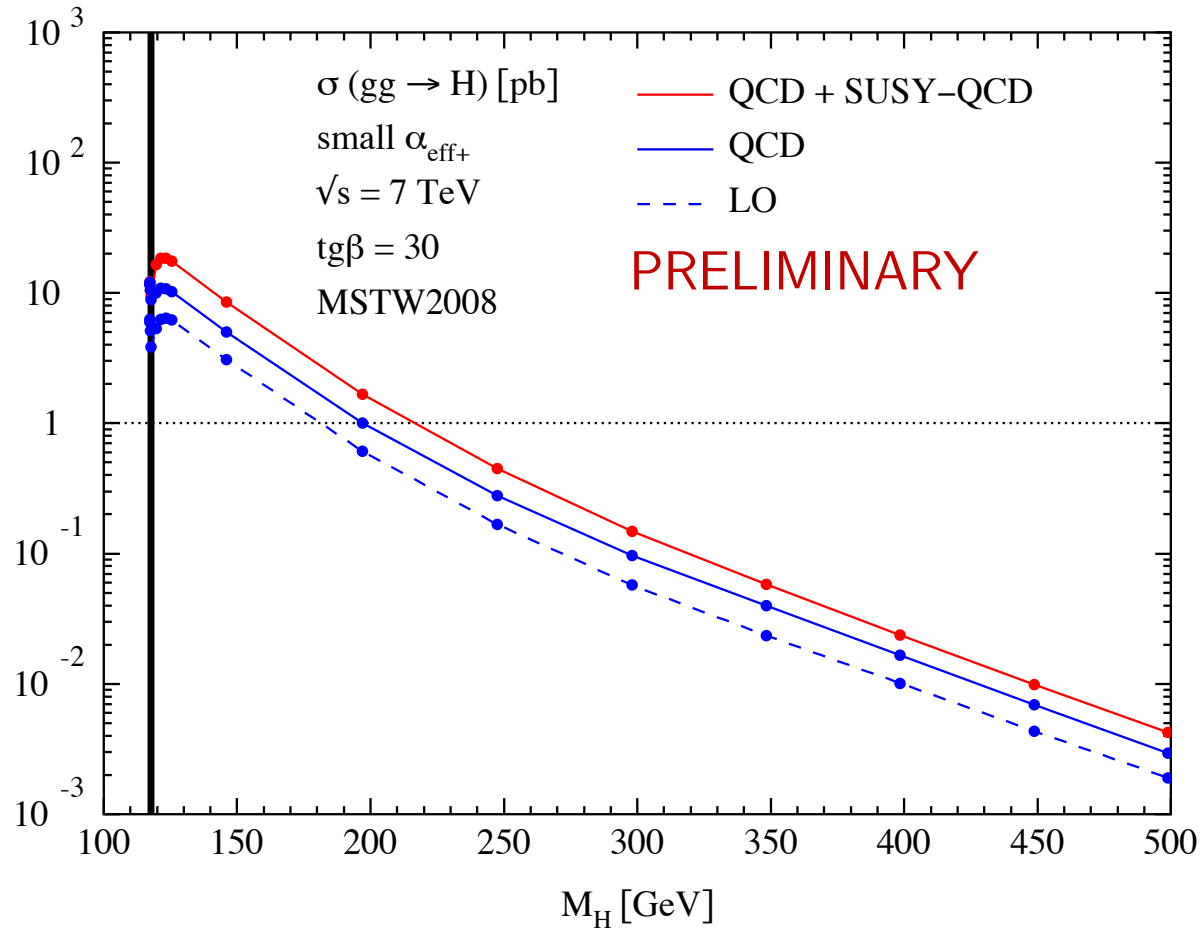
$$m_{\tilde{b}_1} = 601 \text{ GeV} \quad m_{\tilde{b}_2} = 961 \text{ GeV}$$

Genuine SUSY-QCD Corrections - K -factor



MMM, Rzehak, Spira

Genuine SUSY-QCD Corrections - Total Cross Section



MMM, Rzehak, Spira

Higher order corrections to Higgs boson production at the LHC

• W/Z Fusion

NLO QCD σ (SM/MSSM)	~ 5 bis 10%	Han, Valencia, Willenbrock	EW & QCD	$\sim 5\%$	Ciccolini, Denner Dittmaier
distributions	$\sim 10\%$	Figy, Oleari, Zeppenfeld; Berger, Campbell	SUSY QCD	klein	Djouadi, Spira
NLO QCD H+3j		Figy, Hankele, Zeppenfeld	SUSY QCD&EW	klein	Hollik et al. Figy et al.

• Higgs-strahlung

NLO QCD (SM/MSSM)	$\sim +30\%$ (Drell-Yan)	Han, Willenbrock		
NNLO QCD (SM/MSSM)	$\sim +5 - 10\%$	Harlander, Kilgore; Hamberg et al. Brein, Djouadi, Harlander		$\Delta_{\text{theor}} \sim 5\%$
SUSY QCD: klein	Djouadi, Spira	Volle EW(SM): $-5-10\%$	Ciccolini, Dittmaier, Krämer	

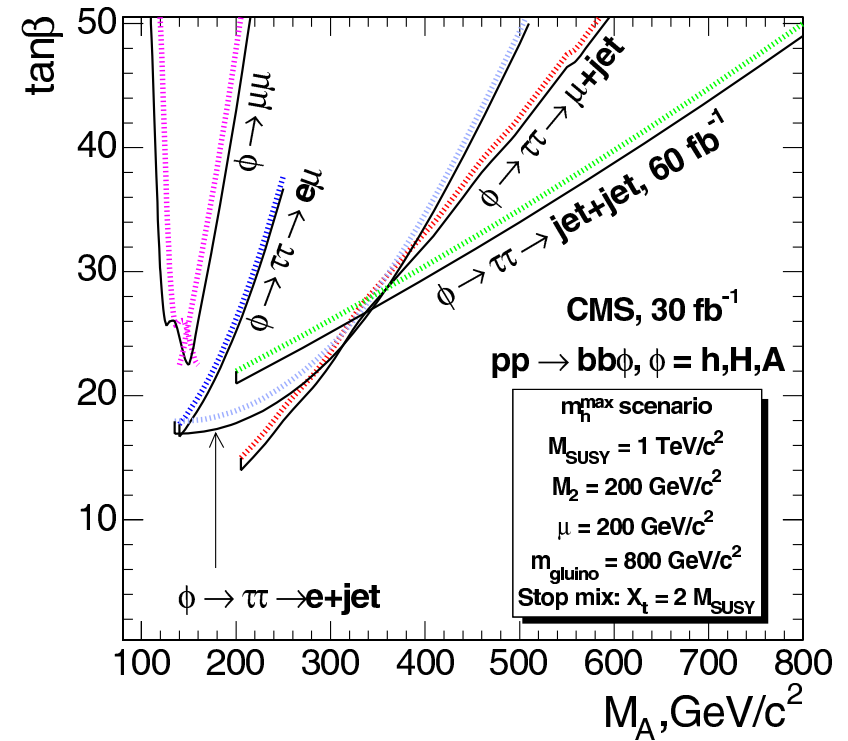
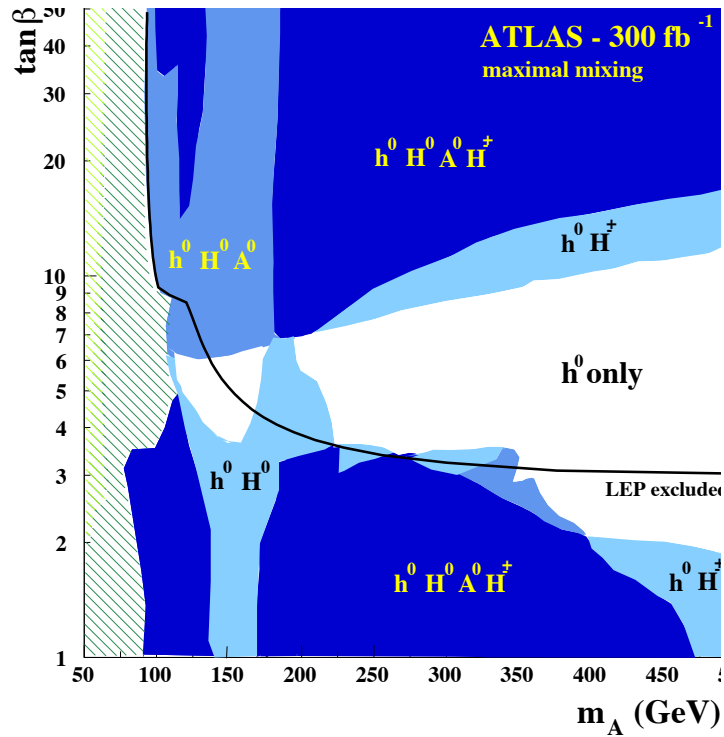
• Associated production

$b\bar{b}\Phi^0$	NLO	Dittmaier, Krämer, Spira, Walser; Dawson, Jackson, Reina, Wackerroth	$b\bar{b} \rightarrow \Phi^0$	NLO, NNLO	Dicus, Willenbrock; Stelzer et al.; Balazs et al.; Campbell et al.; Harlander, Kilgore; Kidonakis
t Beitr.	NNLO	Boudjema, Ninh		EW	Dittmaier, Krämer, Mück, Schlüter
	SUSY QCD	Gao et al.; Hollik, Rauch	$b\bar{b} \rightarrow \Phi^0$ $bg \rightarrow b\Phi^0$	SUSY QCD SUSY EW	Dawson, Jackson (also $bg \rightarrow b\Phi^0$) Dawson, Jaiswal; Beccaria et al. et al.
$t\bar{t}\Phi^0$	NLO QCD +20%	Beenakker et al.; Dawson et al.	SUSY QCD	20-30%	Peng et al.; Dittmaier et al.

MSSM Higgs Boson Search at the LHC

ATLAS

CMS



accuracy: $\delta M_H / M_H (\text{SM}/\text{MSSM}) \sim 10^{-3}$

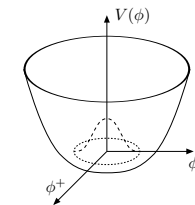
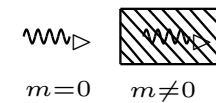
Experimental verification of the Higgs mechanism

Higgs mechanism:

Creation of particle masses without violating gauge principles

Test of the Higgs mechanism

- Discovery – m
- Interaction with a scalar Higgs with $v = 246 \text{ GeV} \neq 0$ $\rightsquigarrow g_{HXX} \sim m_X$
- Spin- and parity quantum numbers – J^{PC}
- EWSB requires Higgs potential – $\lambda_{HHH}, \lambda_{HHHH}$

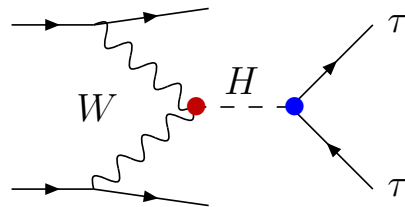


Determination of the Higgs Couplings

Strategy

Combination of the **Higgs production** and **decay channels** \Rightarrow Higgs decay rates, absolute couplings

E.g.:



$$\sim \Gamma_{WW} \text{BR}(H \rightarrow \tau\tau)$$

Problem

- total Higgs production cross section not measurable
 - some Higgs decay channels not observable
- \Rightarrow only ratios of couplings are measurable

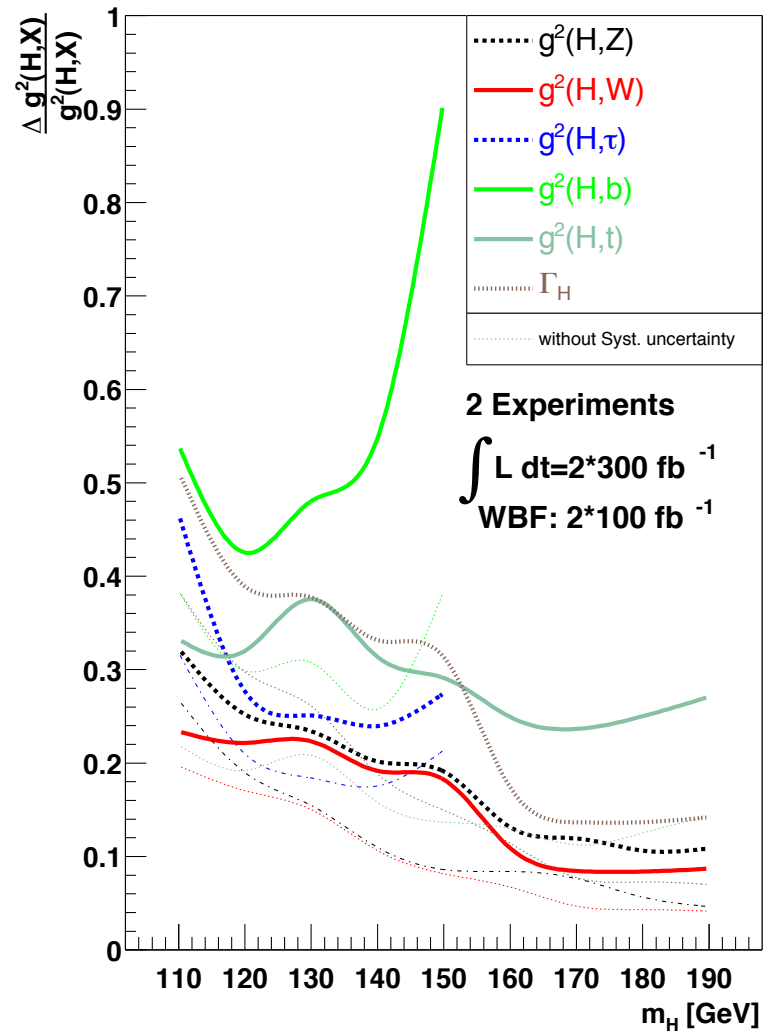
Ansatz

Mild theoretical assumptions \Rightarrow total Higgs width and absolute couplings

- Light Higgs with SM-like couplings Kinnunen, Nikitenko, Richter-Was, Zeppenfeld
- General two-Higgs doublet model Dührssen, Heinemeyer, Logan, Rainwater, Weiglein, Zeppenfeld

Determination of the Higgs Couplings

Dührssen, Heinemeyer, Logan, Rainwater, Weiglein, Zeppenfeld



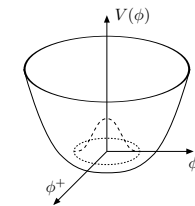
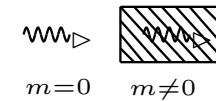
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- Spin- and parity quantum numbers – J^{PC}
- EWSB requires Higgs potential – $\lambda_{HHH}, \lambda_{HHHH}$



Higgs boson quantum numbers

Quantum numbers of the Higgs boson: J^{PC}

J spin
 P parity
 C charge conjugation

◇ $\gamma\gamma \rightarrow H$ or $H \rightarrow \gamma\gamma \rightsquigarrow J \neq 1$.

Spin and CP quantum numbers: angular correlations

- angular correlations in production: Hjj in vector boson fusion, gluon gluon fusion

Plehn, Rainwater, Zeppenfeld;
Hankele, Klämke, Zeppenfeld

- angular correlations in Higgs decays, e.g. $H \rightarrow ZZ \rightarrow l^+l^-l^+l^-$

Dell'Aquila, Nelson;
Kramer, Kühn, Stong, Zerwas;
Choi, Miller, MMM, Zerwas; Bluj;
Buszello, Fleck, Marquard, van der Bij

observables sensitive to CP -violation

Godbole, Miller, MMM

- below ZZ threshold: angular correlations, threshold effects

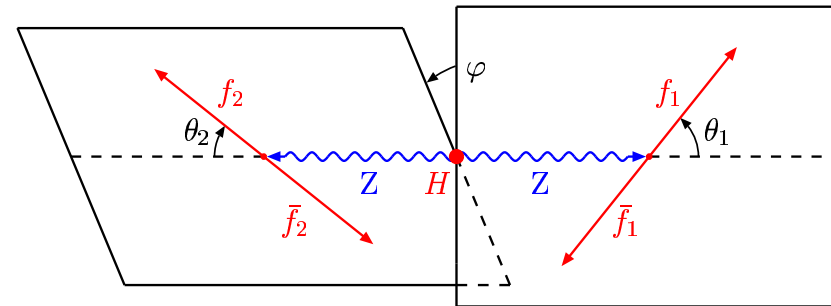
Choi, Miller, MMM, Zerwas
Buszello, Marquard

Higgs boson quantum numbers

Miller, Choi, Eberle, MMM, Zerwas; Choi, Miller, MMM, Zerwas

- ◇ Determination of spin and parity in

$$gg \rightarrow H \rightarrow ZZ^{(*)} \rightarrow (f_1\bar{f}_1)(f_2\bar{f}_2)$$



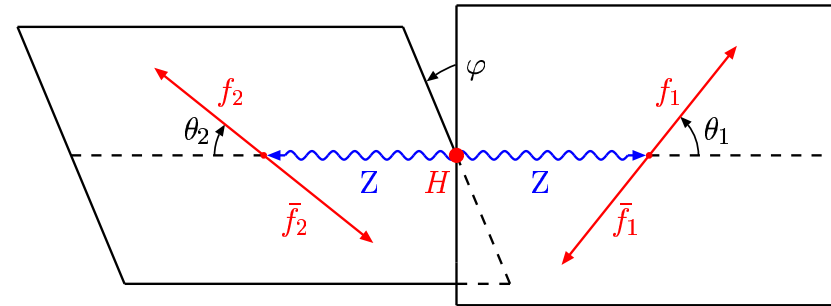
- ◇ Helicity methods: general HZZ coupling for arbitrary spin and parity

$$\langle Z(\lambda_1)Z(\lambda_2)|H_{\mathcal{J}}(m)\rangle = \frac{g_W M_Z}{\cos \theta_W} \mathcal{T}_{\lambda_1\lambda_2} d_{m,\lambda_1-\lambda_2}^{\mathcal{J}}(\Theta) e^{-i(m-\lambda_1+\lambda_2)\Phi}$$

Higgs boson quantum numbers

◇ Determination of spin and parity in

$$gg \rightarrow H \rightarrow ZZ^{(*)} \rightarrow (f_1\bar{f}_1)(f_2\bar{f}_2)$$



◇ Helicity methods: general HZZ coupling for arbitrary spin and parity

◇ Threshold behaviour and angular correlations \rightsquigarrow determination of $\mathcal{J}^{\mathcal{P}}$

• $M_H < 2M_Z$: $d\Gamma/dM_*^2 \sim \beta$ for $\mathcal{J}^{\mathcal{P}} = 0^+$

Choi, Miller, MM, Zerwas

◇ $d\Gamma/dM_*^2$ rules out $\mathcal{J}^{\mathcal{P}} = 0^-, 1^-, 2^-, 3^\pm, 4^\pm$

◇ $d\Gamma/dM_*^2$ and no $[1 + \cos^2 \theta_1] \sin^2 \theta_2$

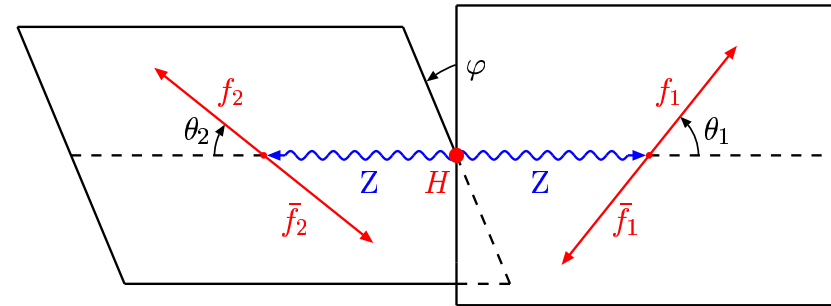
$[1 + \cos^2 \theta_2] \sin^2 \theta_1$

rules out $\mathcal{J}^{\mathcal{P}} = 1^+, 2^+$

Higgs boson quantum numbers

◇ Determination of spin and parity in

$$gg \rightarrow H \rightarrow ZZ^{(*)} \rightarrow (f_1\bar{f}_1)(f_2\bar{f}_2)$$



◇ Helicity methods: general HZZ coupling for arbitrary spin and parity

◇ Threshold behaviour and angular correlations \rightsquigarrow determination of $\mathcal{J}^{\mathcal{P}}$

• $M_H > 2M_Z$:

Choi, Miller, MM, Zerwas

◇ odd normality: $\mathcal{J}^{\mathcal{P}} = 0^-, 1^+, 2^-, 3^+, \dots$ excluded by non-zero $\sin^2 \theta_1 \sin^2 \theta_2$

◇ even normality: $\mathcal{J}^{\mathcal{P}} = 1^-, 3^-, \dots$ excluded by non-zero $\sin^2 \theta_1 \sin^2 \theta_2$

◇ rule out $\mathcal{J}^{\mathcal{P}} = 2^+, 4^+$ with:

$$\frac{d\sigma}{d\cos\theta} [gg/\gamma\gamma \rightarrow H \rightarrow ZZ] \quad \text{only isotropic for spin 0}$$

\mathcal{CP} Violation

- \mathcal{CP} Violation: Examine behaviour with

most general vertex =
sum of even and odd normality tensors

- Case spin 0: $p = p_{Z_1} + p_{Z_2}, k = p_{Z_1} - p_{Z_2}$

$$\text{Vertex } HZZ : \frac{igM_Z}{\cos\theta_W} \left[a g_{\mu\nu} + \frac{b}{M_Z^2} p_\mu p_\nu + i \frac{c}{M_Z^2} \epsilon_{\mu\nu\alpha\beta} p^\alpha k^\beta \right]$$

- ◇ $a = 1, b = c = 0$: SM
- ◇ $(a \neq 0 \wedge c \neq 0) \vee (b \neq 0 \wedge c \neq 0)$: \mathcal{CP} -violation

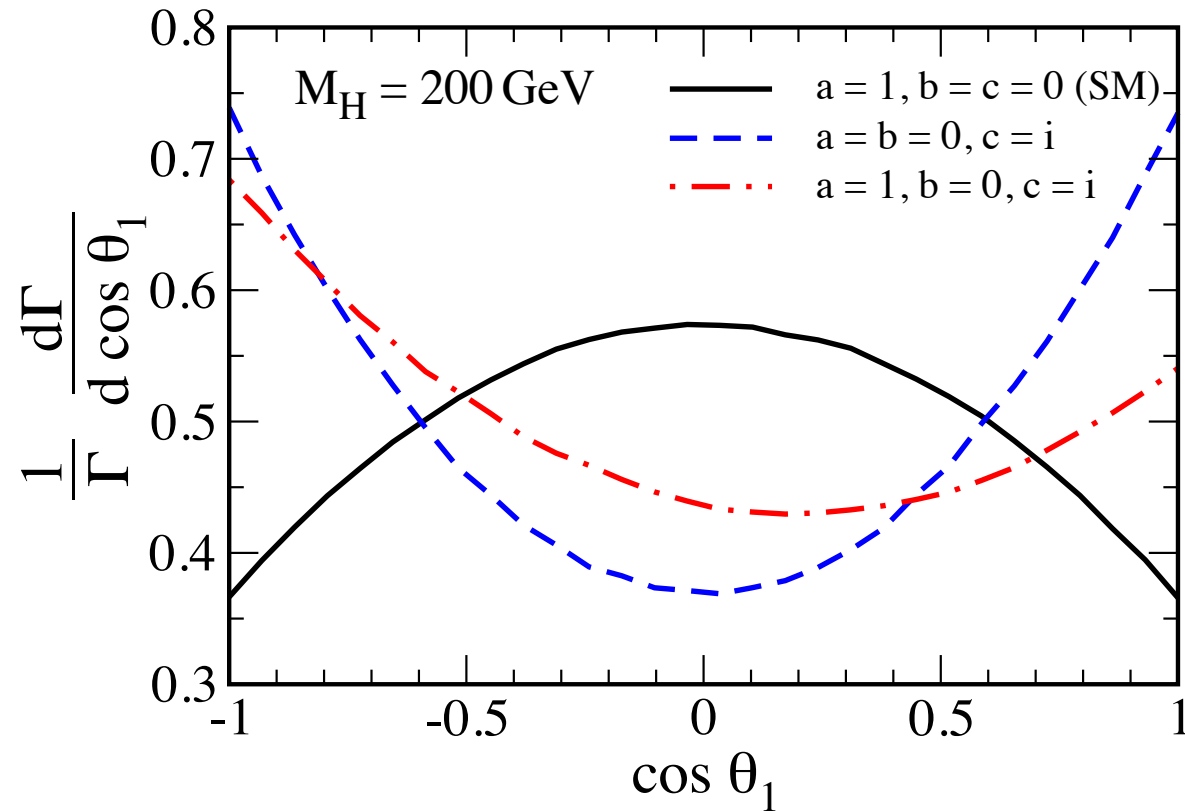
- Observables sensitive to \mathcal{CP}

- ◇ angle ϕ between oriented Z decay planes in the Higgs rest frame
- ◇ \cos of the fermion polar angle θ in the Z rest frame

Higgs boson quantum numbers

angular distribution in $\cos \theta$

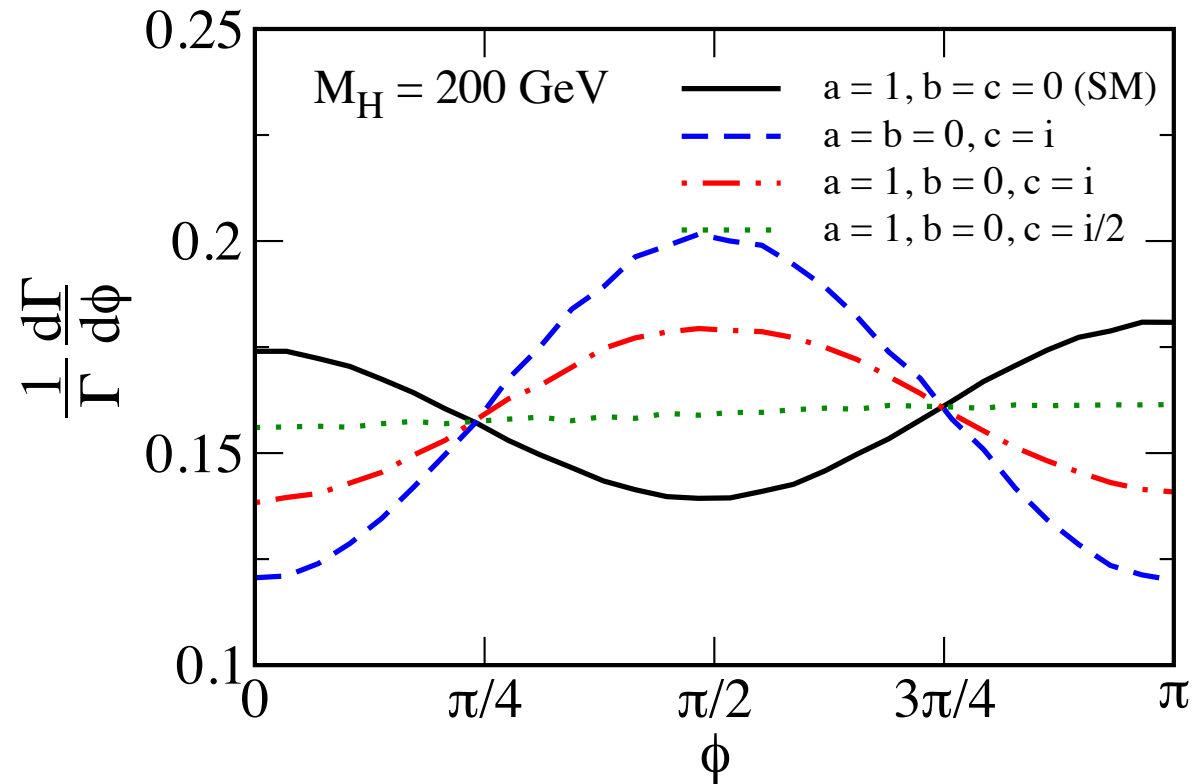
Godbole, Miller, MMM



Higgs boson quantum numbers

angular distribution in ϕ

Godbole, Miller, MMM



Higgs boson quantum numbers

Asymmetries sensitive to \mathcal{CP}

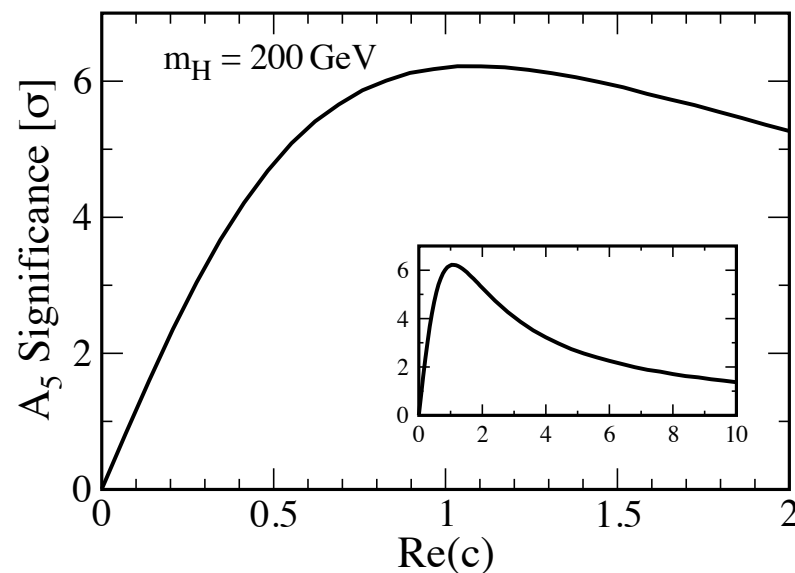
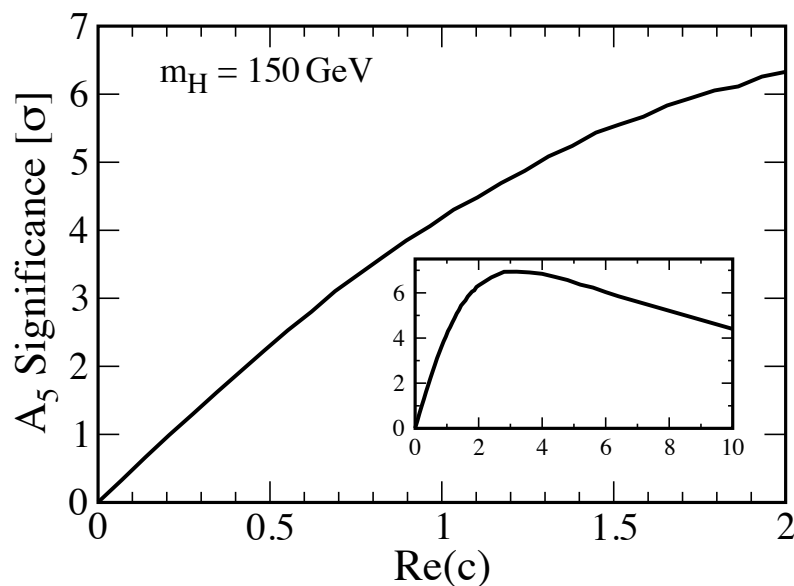
Godbole, Miller, MMM

◇ Example:

$$O_5 = \sin \theta_1 \sin \theta_2 \sin \phi [\sin \theta_1 \sin \theta_2 \cos \phi - \cos \theta_1 \cos \theta_2]$$

$$O_5 = \frac{[(\vec{p}_{4H} \times \vec{p}_{3H}) \cdot \vec{p}_{1H}][(\vec{p}_{1Z} - \vec{p}_{2Z}) \cdot \vec{p}_{3Z}]}{|\vec{p}_{3H} + \vec{p}_{4H}| |\vec{p}_{3Z} - \vec{p}_{4Z}|^2 |\vec{p}_{1Z} - \vec{p}_{2Z}|^2 / 8}$$

$$\mathcal{A}_5 = \frac{\Gamma(O_5 > 0) - \Gamma(O_5 < 0)}{\Gamma(O_5 > 0) + \Gamma(O_5 < 0)}$$



Higgs boson quantum numbers

gluon gluon fusion

CP-even Htt can be distinguished from CP-odd at $> 5\sigma$ ($M_H = 160$ GeV)

Klämke, Zeppenfeld

$H \rightarrow ZZ \rightarrow 4l$

consistency with SM; $0^-, 1^\pm$ excluded
($\int \mathcal{L} = 100\text{fb}^{-1}$, $M_H = 200$ GeV)

Buszello, Fleck,
Marquard, van der Bij

CMS: $H \rightarrow ZZ \rightarrow 4l$

scalar, pseudoscalar can be distinguished at 3σ
($\int \mathcal{L} = 60\text{fb}^{-1}$, $M_H = 300$ GeV)

CMS

ALTAIS: $H \rightarrow ZZ \rightarrow 4l$

strong limits to anomalous couplings

Buszello, Fleck, Marquard,
van der Bij

▷ CP-odd excluded at 8.7σ (2.9σ) for
 $M_H = 200$ GeV (130 GeV), $\int \mathcal{L} = 100\text{fb}^{-1}$

Strässner

▷ b/c w/ precision 1 (0.35/0.2)

for $M_H = 130$ GeV (higher M_H)

▷ include info from measured signal cxn

↔ further increase in precision

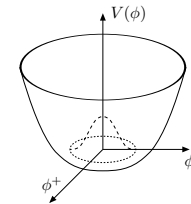
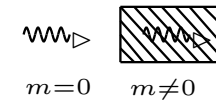
Experimental verification of the Higgs mechanism

Higgs mechanism:

Creation of particle masses without violating gauge principles

Test of the Higgs mechanism

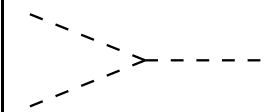
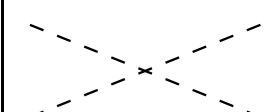
- Discovery – m
- Interaction with a scalar Higgs with $v = 246 \text{ GeV} \neq 0$ $\rightsquigarrow g_{HXX} \sim m_X$
- Spin- and parity quantum numbers – J^{PC}
- EWSB requires Higgs potential – $\lambda_{HHH}, \lambda_{HHHH}$

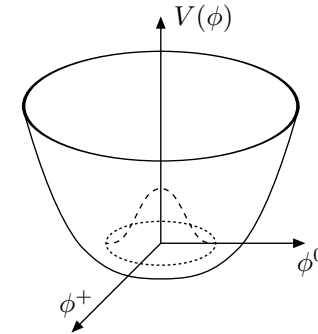


Determination of the Higgs Self-Couplings

The Higgs potential:

$$V(H) = \frac{1}{2!} \lambda_{HH} H^2 + \frac{1}{3!} \lambda_{HHH} H^3 + \frac{1}{4!} \lambda_{HHHH} H^4$$

Trilinear coupling	$\lambda_{HHH} = 3 \frac{M_H^2}{v}$	
Quartic coupling	$\lambda_{HHHH} = 3 \frac{M_H^2}{v^2}$	



Measurement of the Higgs self-couplings and Reconstruction of the Higgs potential	} Experimental verification Of the scalar sector of the Higgs mechanism
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Determination of the Higgs self-couplings at colliders:

λ_{HHH} via Higgs pair production

λ_{HHHH} via triple Higgs production

Higgs-strahlung, WW/ZZ fusion, gg fusion

The Trilinear Higgs Self-Coupling at the LHC

Determination of λ_{HHH} at the LHC

Djouadi, Kilian, MMM, Zerwas;
Lafaye, Miller, Moretti, MMM

double Higgs-strahlung: $q\bar{q} \rightarrow W/Z + HH$

Barger, Han, Phillips

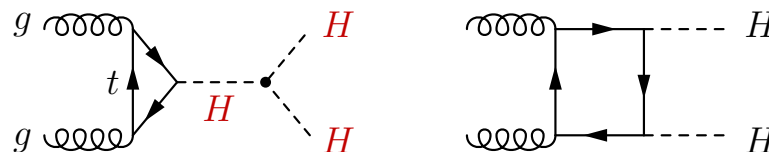
WW/ZZ fusion: $qq \rightarrow qq + HH$

Dicus, Kallianpur, Willenbrock
Abbasabadi, Repko, Dicus, Vega
Dobrovolskaya, Novikov
Eboli, Marques, Novaes, Natale

gluon gluon fusion: $gg \rightarrow HH$

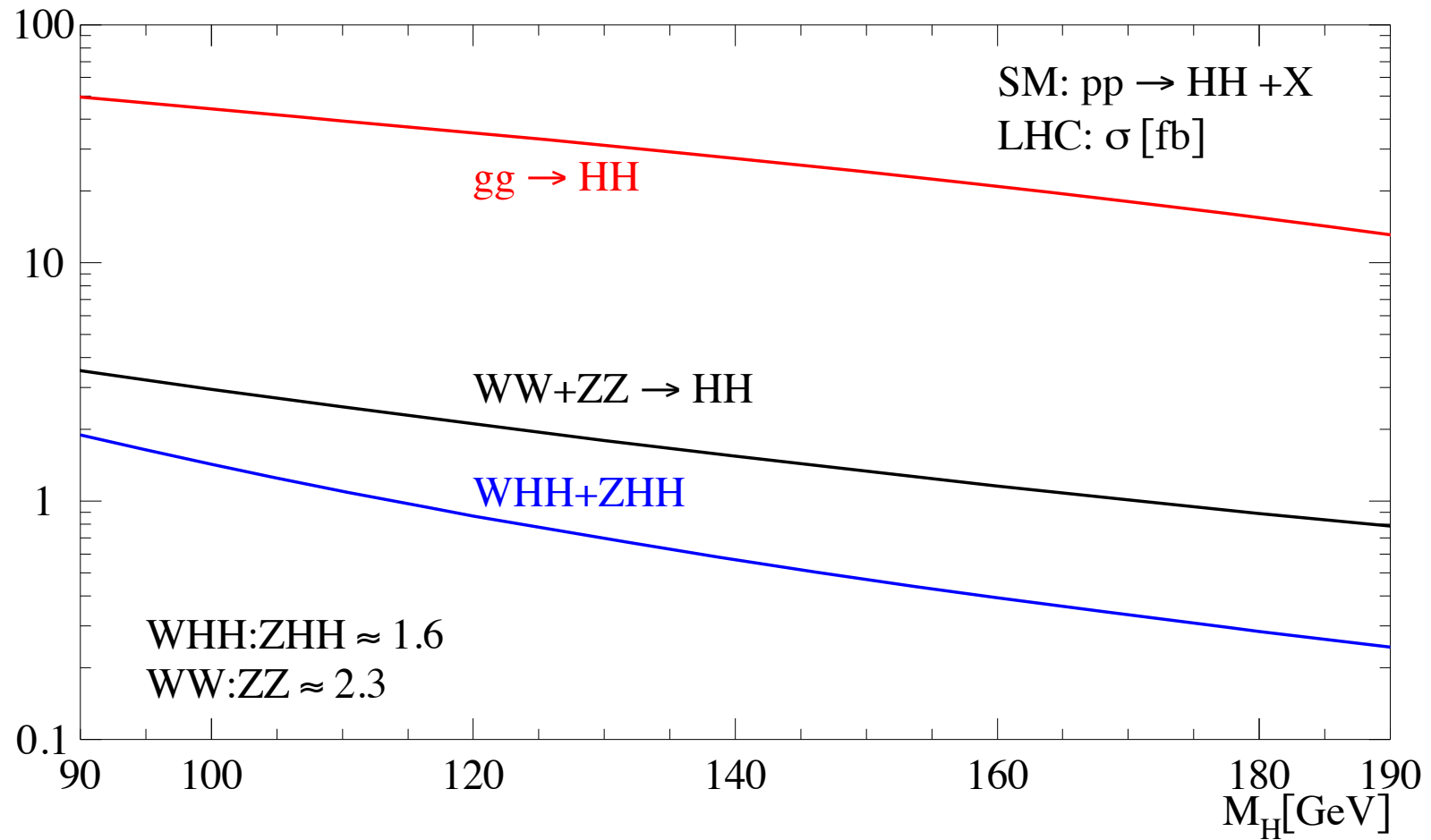
Glover, van der Bij
Plehn, Spira, Zerwas
Dawson, Dittmaier, Spira

gluon gluon fusion - dominant process



Double SM Higgs Production at the LHC

Djouadi, Kilian, MMM, Zerwas



Expected Accuracies in λ_{HHH} at the \mathcal{LHC}

small signal + large QCD background \rightsquigarrow challenge!

$M_H < 140$ GeV: $gg \rightarrow HH \rightarrow b\bar{b}\gamma\gamma$:

Baur,Plehn,Rainwater

○ SLHC [$\int \mathcal{L} = 6 \text{ ab}^{-1}$]:

$M_H = 120$ GeV $\lambda_{HHH} = 0$ exclusion at 90% CL

○ VLHC [$\sqrt{s} = 200$ TeV]:

$M_H = 120$ GeV: $\delta\lambda_{HHH}/\lambda_{HHH} = 20 - 40\%$ at 1σ

$M_H > 140$ GeV: $gg \rightarrow HH \rightarrow W^+W^-W^+W^-$:

Gianotti et al.;Blondel,Clark,Mazzucato
Baur,Plehn,Rainwater
Dahlhoff

○ LHC [$\int \mathcal{L} = 300 \text{ fb}^{-1}$]:

$150 \lesssim M_H \lesssim 200$ GeV: $\lambda_{HHH} = 0$ exclusion at 95% CL

○ SLHC [$\int \mathcal{L} = 3 \text{ ab}^{-1}$]:

$150 < M_H < 200$ GeV $\delta\lambda_{HHH}/\lambda_{HHH} = 20 - 30\%$ at 1σ

Higgs Physics - Beyond

UnHiggs
Gaugephobic Higgs
Composite Higgs
Gauge Higgs
Simplest Higgs
Private Higgs
Intermediate Higgs
Fat Higgs
Twin Higgs
Phantom Higgs
Little Higgs
Littlest Higgs
Slim Higgs
Higgsless
Portal Higgs
Lone Higgs

Composite Higgs Boson - Introduction

- **Higgs: bound state from a strongly interacting sector** Kaplan,Georgi;Dimopoulos eal;Dugan eal
- **SILH** effective low energy description, Higgs couplings modified by $\xi = \frac{v^2}{f^2}$ Giudice,Grojean
Pomarol,Rattazzi
- **Fermion couplings** depend on embedding into representations of the bulk symmetry Contino eal;
Agashe eal

spinorial representations of $SO(5)$

fundamental representations of $SO(5)$

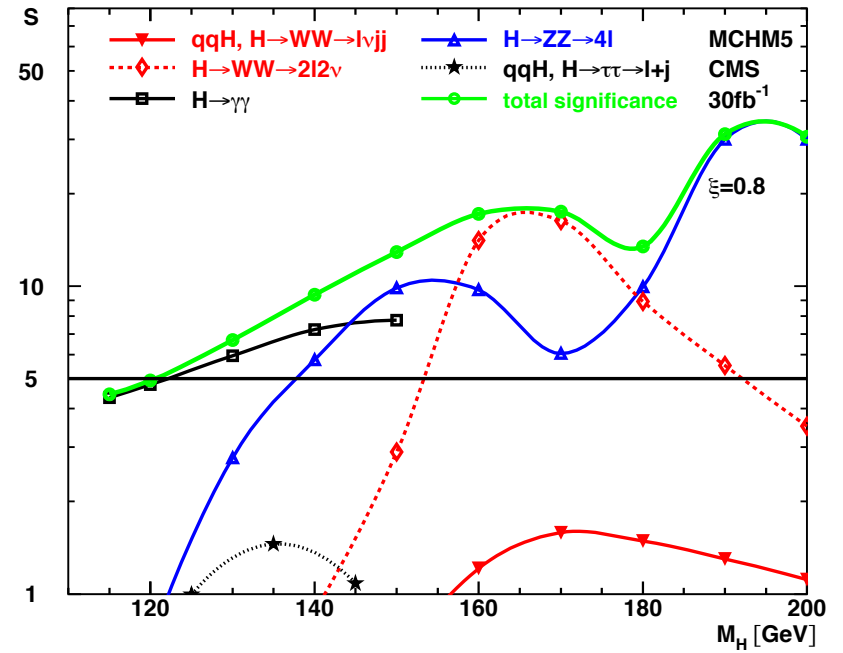
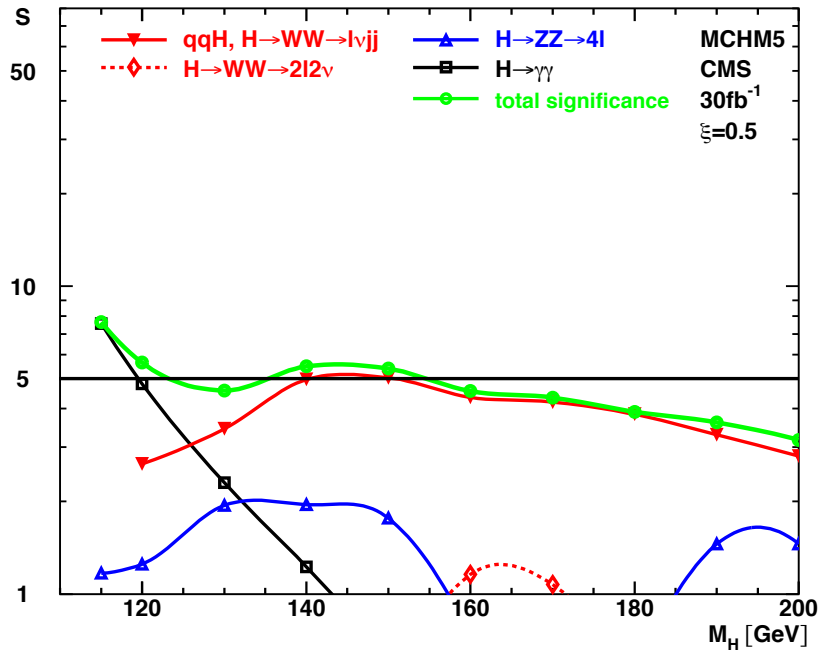
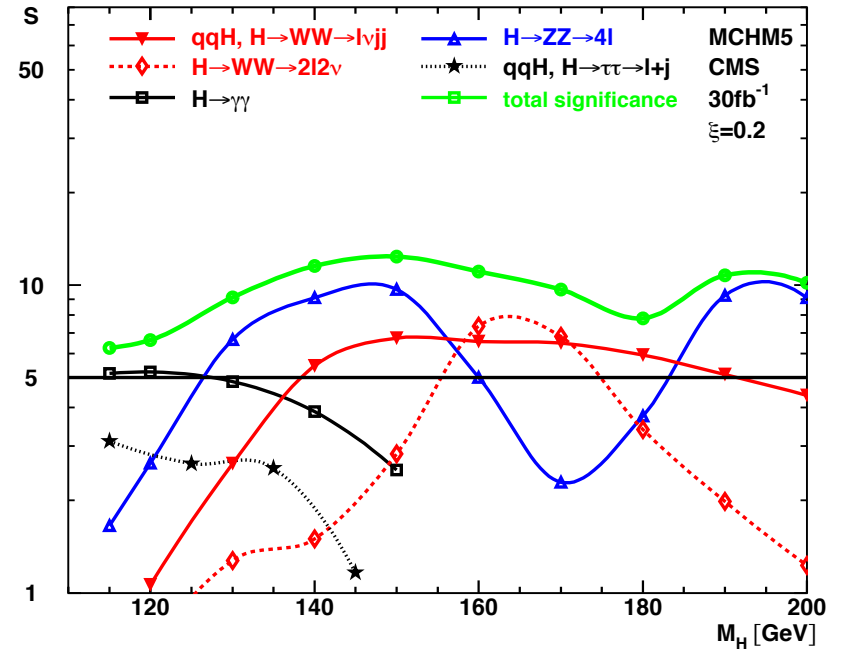
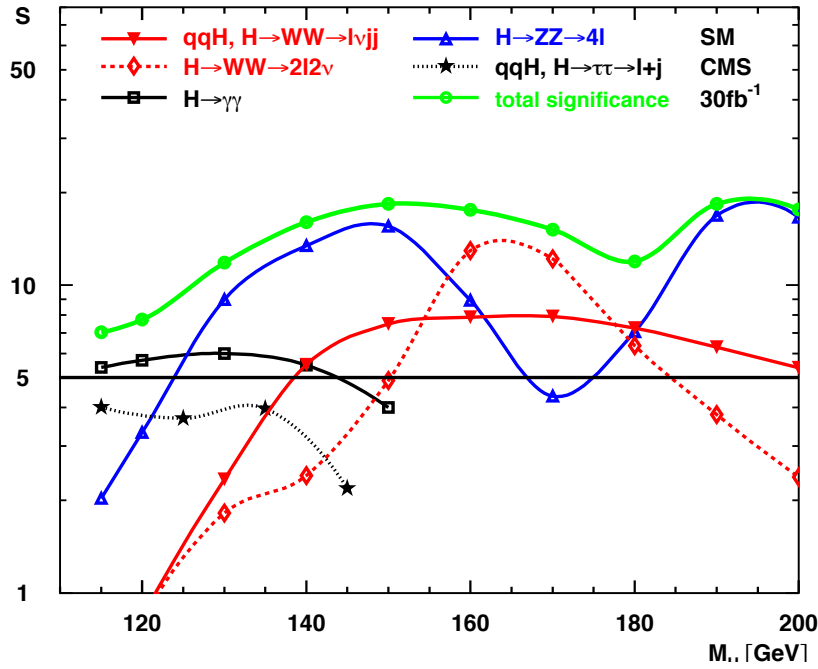
MCHM4

MCHM5

MCHM4	MCHM5
$g_{HVV} = g_{HVV}^{SM} \sqrt{1-\xi}$	$g_{HVV} = g_{HVV}^{SM} \sqrt{1-\xi}$
$g_{Hff} = g_{Hff}^{SM} \sqrt{1-\xi}$	$g_{Hff} = g_{Hff}^{SM} \frac{(1-2\xi)}{\sqrt{1-\xi}}$
universal factor \rightsquigarrow BRs unchanged	g_{Hff} coupling vanishes for $\xi = 0.5$

- **Impact on** BR's, Γ_{tot} , production cross sections, [Higgs searches at the LHC](#) Espinosa,Grojean,MMM
(gg fusion at NNLO Furlan '11)

Significances MCHM5



Sensitivity to the triple Higgs self-coupling

- Can we extract the Higgs self-coupling?

Gröber, Mühlleitner

- First step: plot sensitivity areas

- Sensitivity criteria: $\int \mathcal{L} = 300 \text{ fb}^{-1}$

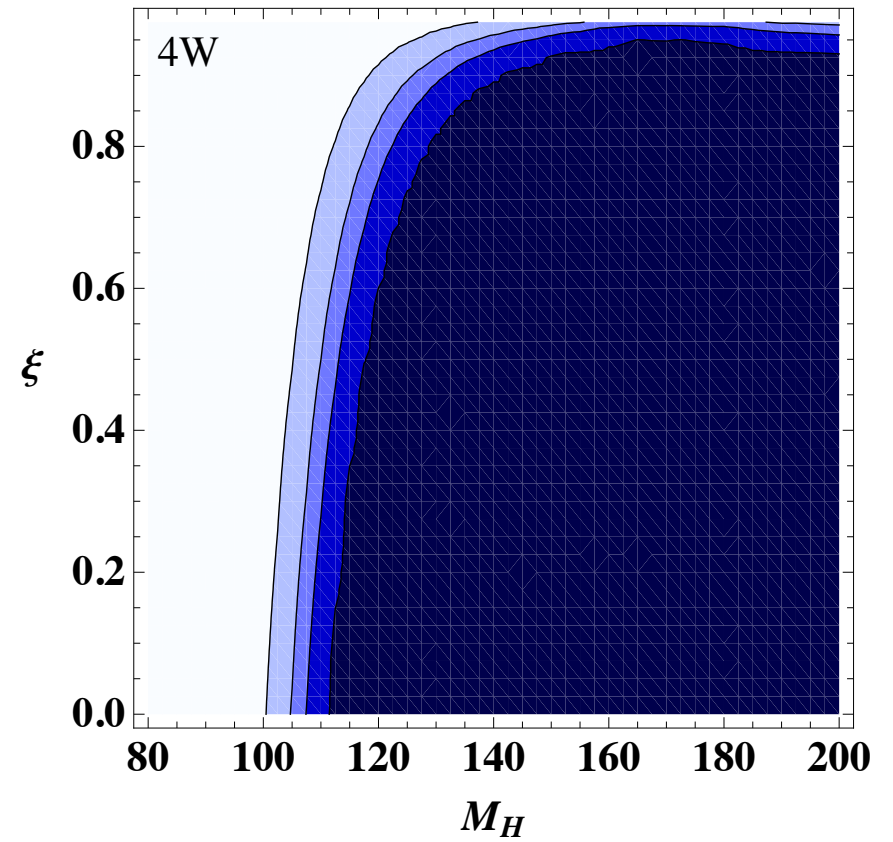
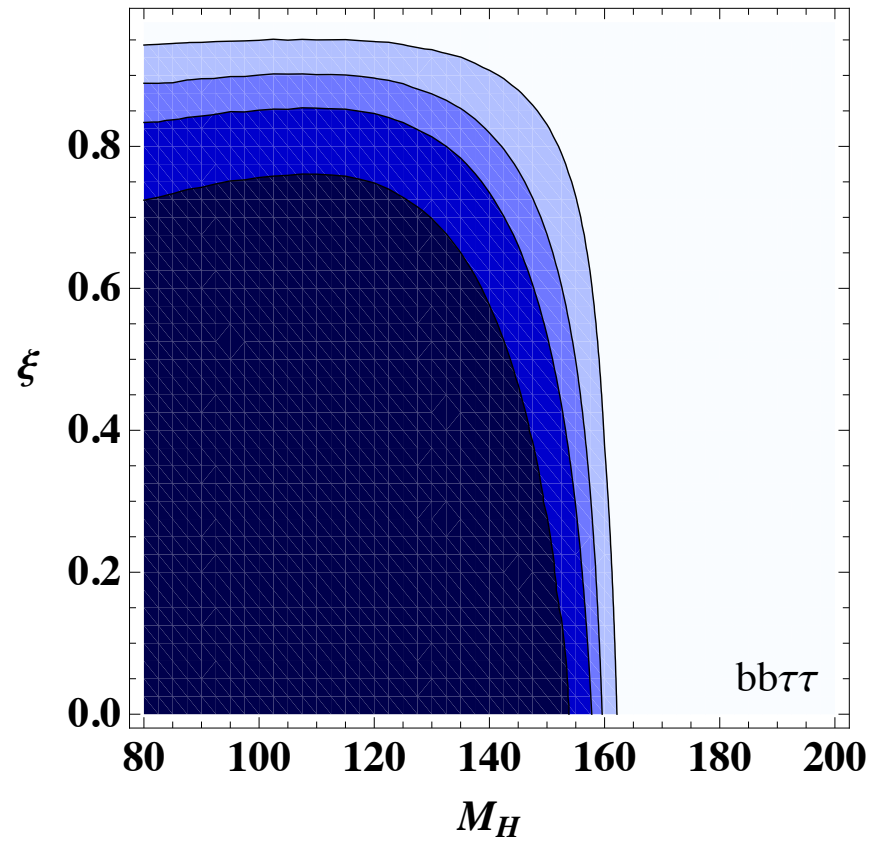
$$\begin{aligned} \lambda_{HHH} \neq 0 &\rightsquigarrow \sigma^{HH}(\lambda_{HHH}) &\rightsquigarrow N_{\text{events}}^{\lambda \neq 0} \\ \lambda_{HHH} = 0 &\rightsquigarrow \sigma^{HH}(\lambda_{HHH} = 0) &\rightsquigarrow N_{\text{events}}^{\lambda = 0} \end{aligned}$$

Demand: $N^{\lambda=0} + \beta\sqrt{N^{\lambda=0}} < N^{\lambda \neq 0}$ or $N^{\lambda=0} - \beta\sqrt{N^{\lambda=0}} > N^{\lambda \neq 0}$ ($\beta = 1, 2, 3, 5$)

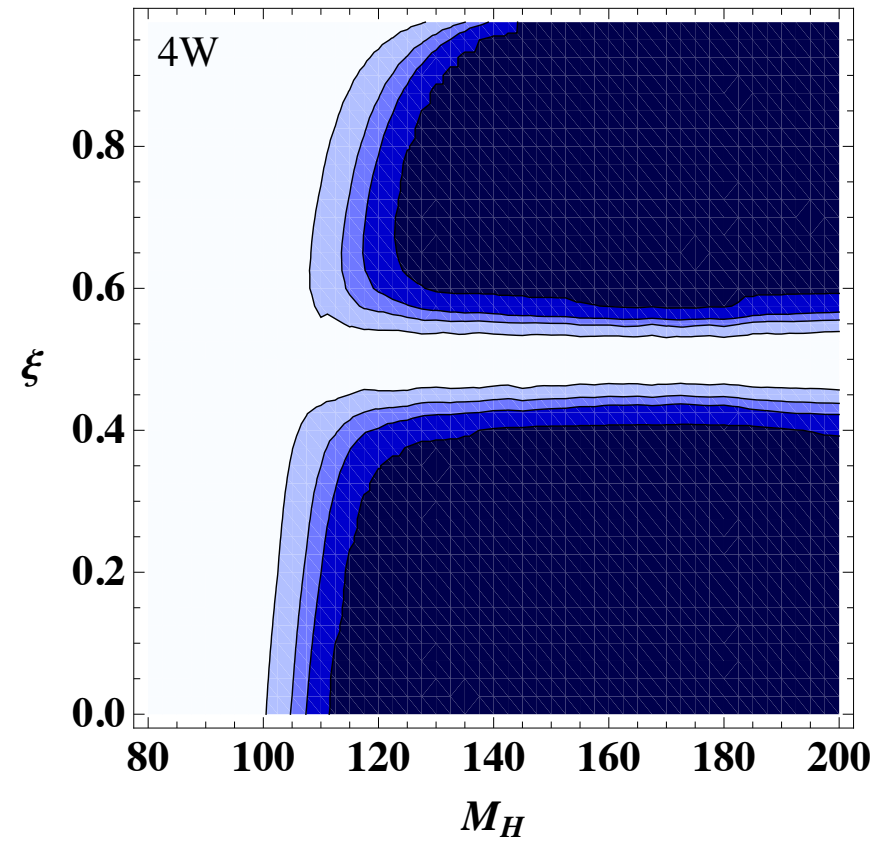
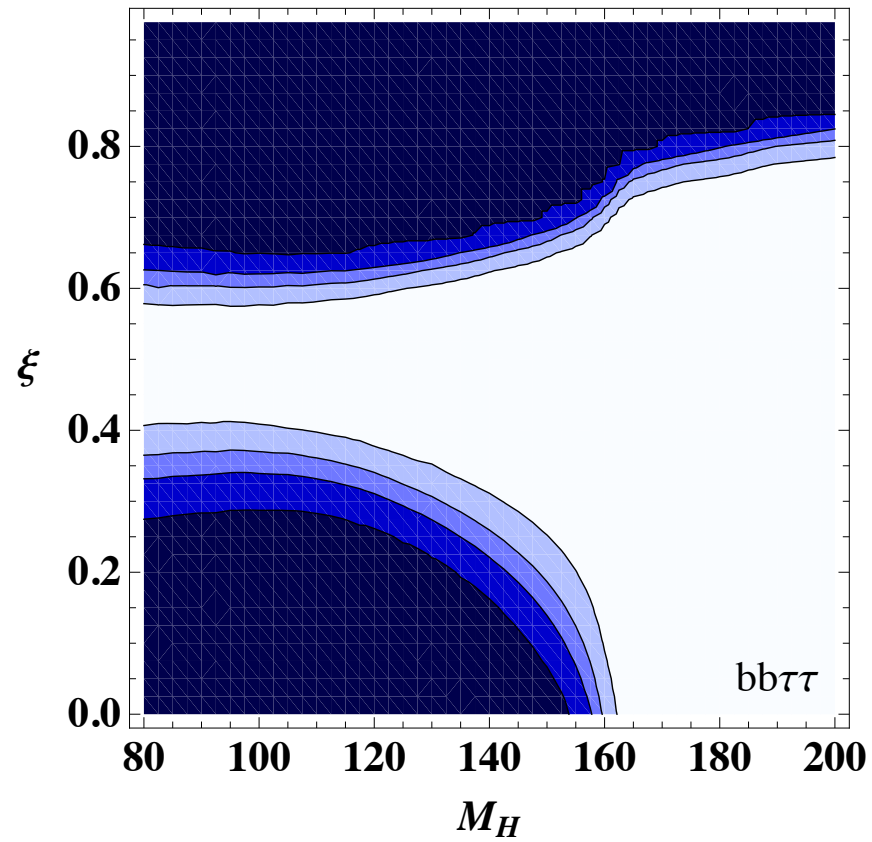
- Plots: $gg \rightarrow HH$ for MCHM4, MCHM5

- Final states: $HH \rightarrow bb\gamma\gamma$, $HH \rightarrow W^+W^-W^+W^-$, $HH \rightarrow bb\tau\tau$, $HH \rightarrow bb\mu\mu$

• Significances to non-vanishing λ_{HHH} in MCHM4



• Significances to non-vanishing λ_{HHH} in MCHM5



Conclusions

The LHC is a discovery machine

- Higgs particle(s) can be discovered
 - First tests of the Higgs mechanism are possible
 - ◇ Determination of the Higgs couplings to fermions and bosons
 - ◇ Determination of the Higgs quantum numbers (spin and CP)
 - ◇ Determination of the trilinear Higgs self-coupling(s)
- ⇒ Important steps towards the understanding of the creation of masses

Conclusions

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 - ◇ Determination of the Higgs quantum numbers (spin and CP)
 - ◇ Determination of the trilinear Higgs self-coupling(s)
- ⇒ Important steps towards the understanding of the creation of masses
- **Composite Higgs Model** Higgs as pseudo-Goldstone boson of strong sector
 - **After discovery:** Which Higgs boson have we discovered?

Conclusions

UnHiggs
Gaugephobic Higgs
Composite Higgs
Gauge Higgs
Simplest Higgs

Private Higgs
Intermediate Higgs
Fat Higgs
Twin Higgs
Phantom Higgs

Little Higgs
Littlest Higgs
Higgsless
Lone Higgs

Slim Higgs
Portal Higgs

Backup Slides

Differential distributions

◇ Double polar angular distribution (\mathcal{CP} invariant theory)

$$\begin{aligned} \frac{d\Gamma_H}{d\cos\theta_1 d\cos\theta_2} &\sim \sin^2\theta_1 \sin^2\theta_2 |\mathcal{T}_{00}|^2 + \frac{1}{2}(1 + \cos^2\theta_1)(1 + \cos^2\theta_2) [|\mathcal{T}_{11}|^2 + |\mathcal{T}_{1,-1}|^2] \\ &\quad + (1 + \cos^2\theta_1) \sin^2\theta_2 |\mathcal{T}_{10}|^2 + \sin^2\theta_1 (1 + \cos^2\theta_2) |\mathcal{T}_{01}|^2 \\ &\quad + 2\eta_1\eta_2 \cos\theta_1 \cos\theta_2 [|\mathcal{T}_{11}|^2 - |\mathcal{T}_{1,-1}|^2] \end{aligned}$$

SM: $\mathcal{T}_{00} = M_H^2/(2M_Z^2) - 1$, $\mathcal{T}_{11} = -1$, $\mathcal{T}_{10} = \mathcal{T}_{01} = \mathcal{T}_{1,-1} = 0$

◇ Azimuthal angular distribution (\mathcal{CP} invariant theory)

$$\begin{aligned} \frac{d\Gamma_H}{d\varphi} &\sim |\mathcal{T}_{11}|^2 + |\mathcal{T}_{10}|^2 + |\mathcal{T}_{1,-1}|^2 + |\mathcal{T}_{01}|^2 + |\mathcal{T}_{00}|^2/2 \\ &\quad + \eta_1\eta_2 \left(\frac{3\pi}{8}\right)^2 \Re(\mathcal{T}_{11}\mathcal{T}_{00}^* + \mathcal{T}_{10}\mathcal{T}_{0,-1}^*)\cos\varphi + \frac{1}{4}\Re(\mathcal{T}_{11}\mathcal{T}_{-1,-1}^*)\cos 2\varphi \end{aligned}$$

Determination of spin and parity

- $M_H < 2M_Z$: $d\Gamma/dM_*^2 \sim \beta$ for $\mathcal{J}^P = 0^+$

◇ $d\Gamma/dM_*^2$ rules out $\mathcal{J}^P = 0^-, 1^-, 2^-, 3^\pm, 4^\pm$

◇ $d\Gamma/dM_*^2$ and no $[1 + \cos^2 \theta_1] \sin^2 \theta_2$

$[1 + \cos^2 \theta_2] \sin^2 \theta_1$ rules out $\mathcal{J}^P = 1^+, 2^+$

- $M_H > 2M_Z$:

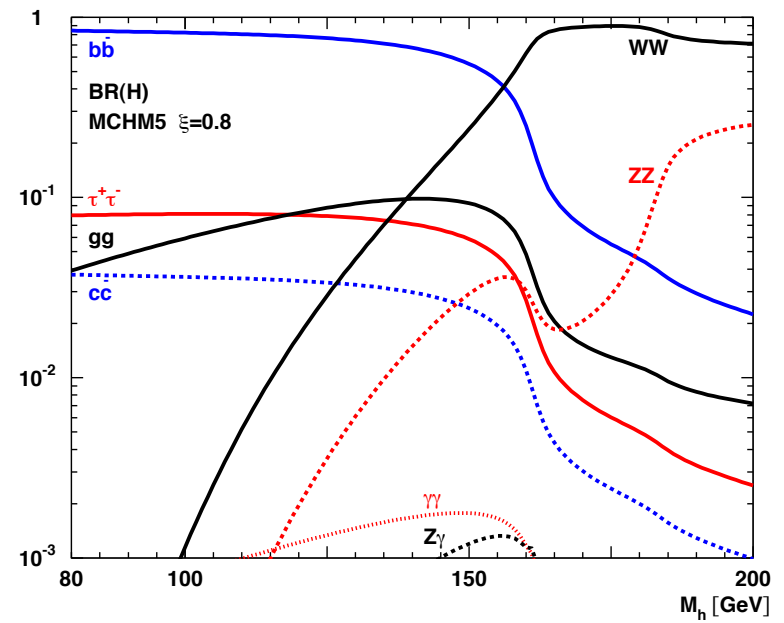
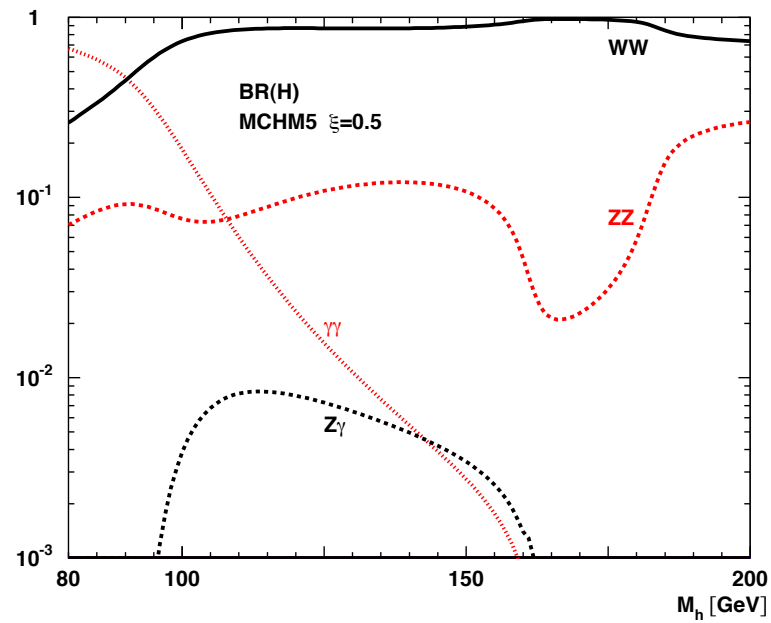
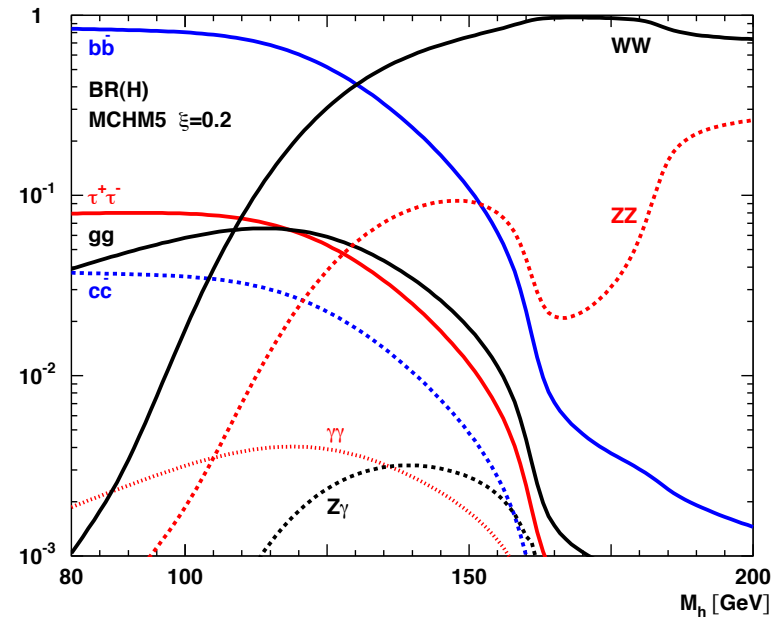
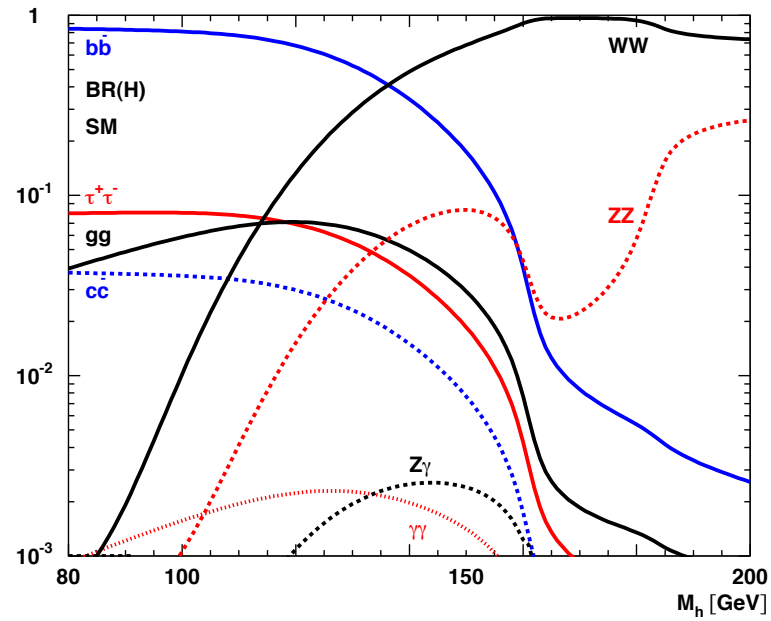
◇ odd normality: $\mathcal{J}^P = 0^-, 1^+, 2^-, 3^+, \dots$ excluded by non-zero $\sin^2 \theta_1 \sin^2 \theta_2$

◇ even normality: $\mathcal{J}^P = 1^-, 3^-, \dots$ excluded by non-zero $\sin^2 \theta_1 \sin^2 \theta_2$

◇ rule out $\mathcal{J}^P = 2^+, 4^+$ with:

$\frac{d\sigma}{d\cos\theta} [gg/\gamma\gamma \rightarrow H \rightarrow ZZ]$ only isotropic for spin 0

• Branching ratios MCHM5



Constraints from EWPT, LEP, Tevatron

- **EWPT constraints**

- ◇ $\hat{T} = c_T \frac{v^2}{f^2} \Rightarrow |c_T \frac{v^2}{f^2}| < 2 \times 10^{-3}$

removed by custodial symmetry

- ◇ $\hat{S} = (c_W + c_B) \frac{m_W^2}{m_\rho^2} \Rightarrow$

$m_\rho \geq (c_B + c_W)^{1/2} 2.5 \text{ TeV}$

- ◇ 1-loop IR effects Barbieri et al

$$\hat{S}, \hat{T} = a \ln m_H + b$$

modified Higgs coupling to matter \Rightarrow

$$\hat{S}, \hat{T} = a((1 - c_H \xi) \ln m_H + c_H \xi \ln \Lambda) + b$$

effective Higgs mass:

$$m_H^{eff} = m_H \left(\frac{\Lambda}{m_H} \right)^{c_H v^2 / f^2} > m_H$$

LEP II, $m_H \approx 115 \text{ GeV}$:

$$c_H \frac{v^2}{f^2} < \frac{1}{3} \sim \frac{1}{2}$$

IR effects can be cancelled by heavy fermions (model-dependent)

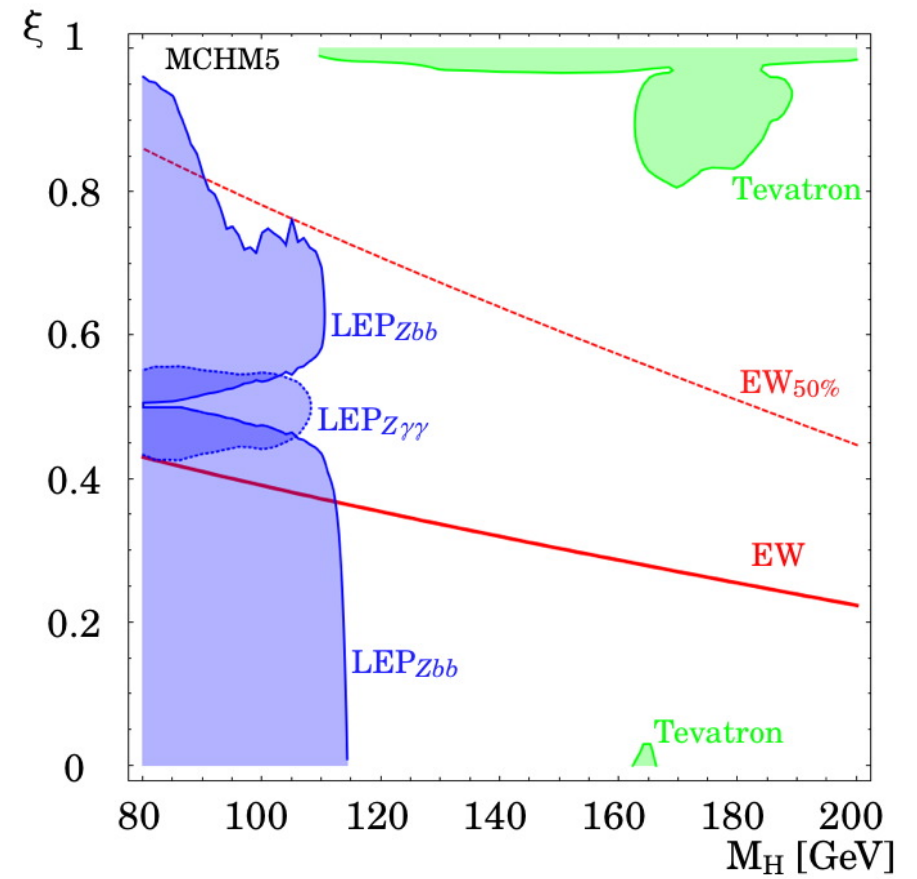
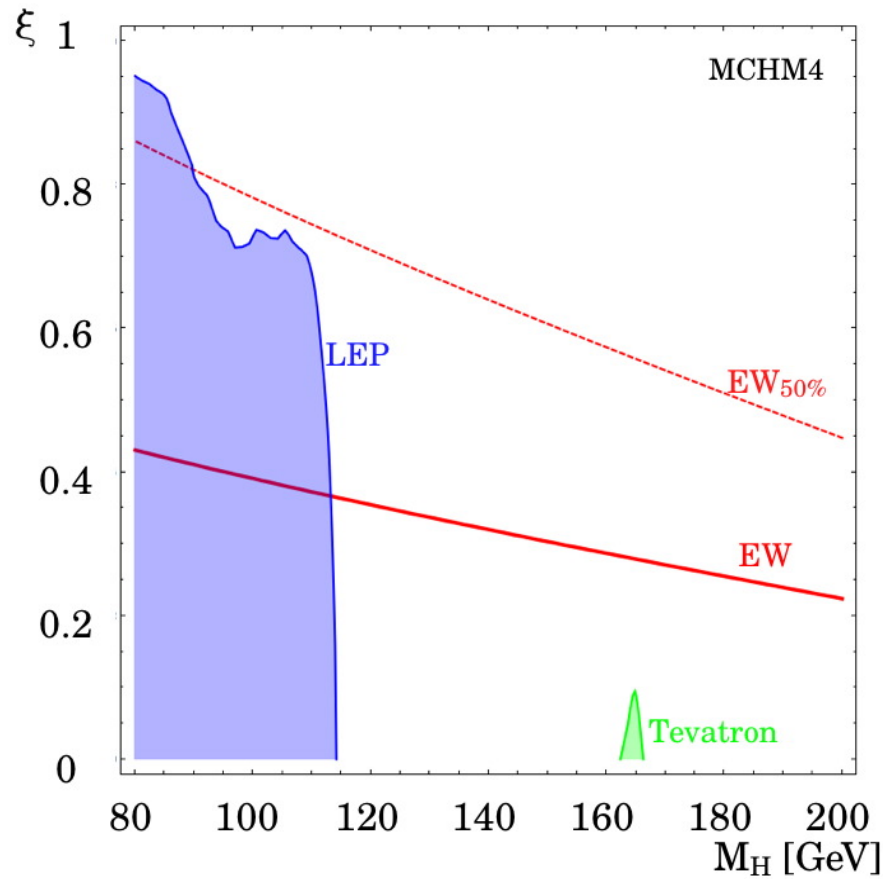
- **Searches at LEP** $e^+ e^- \rightarrow ZH \rightarrow Z b \bar{b}$

- **Tevatron search** most relevant $H \rightarrow WW$

LEP/Tevatron exclusion limits generated with Higgsbounds Bechtel et al

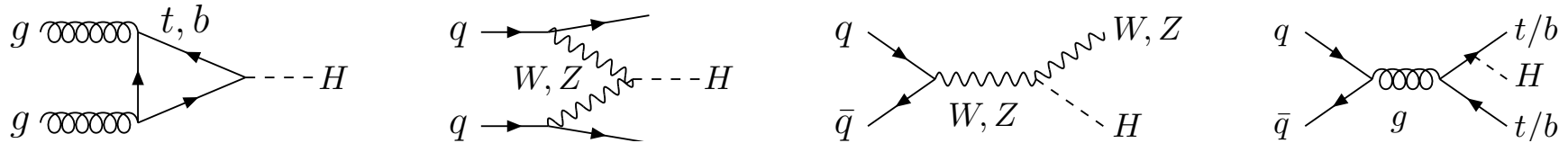
Constraints from EWPT, LEP, Tevatron

Espinosa, Grojean, Mühlleitner



Production Processes

- Production cross sections



- Higgs gauge boson couplings

$$\text{MCHM4/5: } g_{HVV} = g_{HVV}^{SM} \sqrt{1 - \xi}$$

- Higgs fermion couplings

$$\text{MCHM4: } g_{Hff} = g_{Hff}^{SM} \sqrt{1 - \xi}$$

$$\text{MCHM5: } g_{Hff} = g_{Hff}^{SM} \frac{1-2\xi}{\sqrt{1-\xi}} \quad \text{vanishes for } \xi = 0.5!!!$$

NLO QCD corrections: not affected by modified Higgs couplings

$$\sigma_{NLO}(gg \rightarrow H) = \left\{ \begin{array}{l} (1 - \xi) \\ \frac{(1-2\xi)^2}{(1-\xi)} \end{array} \right\} \sigma_{NLO}^{SM}(gg \rightarrow H) \quad \sigma_{NLO}(Ht\bar{t}) \text{ analogous}$$

$$\sigma_{NLO}(qqH) = (1 - \xi) \sigma_{NLO}^{SM}(qqH) \quad \sigma_{NLO}(VH) \text{ analogous}$$

Significances: Composite Higgs

- **Composite Higgs search:**

- ◇ Composite couplings affect signal events, not background events
- ◇ Rescaling factor

$$\kappa = \frac{\sigma_{prod}^{\xi} BR^{\xi}(H \rightarrow X)}{\sigma_{prod}^{SM} BR^{SM}(H \rightarrow X)}$$

- ◇ Exp analyses provide signal & bkg events after cuts, $s^{SM}, b^{SM} \rightsquigarrow$
significances composite model from $s^{\xi} = \kappa s^{SM}$ and $b^{\xi} = b^{SM}$

- **Investigated Channels:** CMS analyses (similar results expected for ATLAS)

Inclusive production with subsequent decay : $H \rightarrow \gamma\gamma$

$H \rightarrow ZZ \rightarrow 2e2\mu, 4e, 4\mu$

$H \rightarrow WW \rightarrow 2l2\nu$

Vector boson fusion with subsequent decay : $H \rightarrow WW \rightarrow l\nu jj$

$H \rightarrow \tau\tau \rightarrow l + j + E_T^{miss}$.