## Physics beyond the Standard Model

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## Exercise Sheet 11

Discussion: Fr, 06.02.15

## Exercise 14: Sfermion sector of the MSSM

In the MSSM, under the assumption of "minimal flavour violation", the following mass terms for the sfermions appear:

$$\begin{aligned} \mathcal{L}_{\tilde{f},\text{mass}} &= -\tilde{f}^{\dagger} M_{\tilde{f}}^2 \tilde{f} \quad \text{with} \quad \tilde{f} = \begin{pmatrix} \tilde{f}_L \\ \tilde{f}_R \end{pmatrix}, \qquad M_{\tilde{f}}^2 = \begin{pmatrix} M_{\tilde{f}}^{LL} + m_{f}^2 & m_f \left( M_{\tilde{f}}^{LR} \right)^* \\ m_f M_{\tilde{f}}^{LR} & M_{\tilde{f}}^{RR} + m_f^2 \end{pmatrix} \quad \text{as well as} \\ M_{\tilde{f}}^{LL} &= M_Z^2 \left( I_3^f - Q_f s_W^2 \right) \cos 2\beta + \begin{cases} M_{\tilde{L}}^2 & \text{for left-handed sleptons} \\ M_{\tilde{Q}}^2 & \text{for left-handed squarks} \end{cases} \\ M_{\tilde{f}}^{RR} &= M_Z^2 \left( Q_f s_W^2 \right) \cos 2\beta + \begin{cases} M_{\tilde{E}}^2 & \text{for right-handed electron-type sleptons} \\ M_{\tilde{U}}^2 & \text{for right-handed up-type squarks} \\ M_{\tilde{f}}^2 & \text{for right-handed down-type squarks} \end{cases} \\ M_{\tilde{f}}^{LR} &= A_f - \mu^* \begin{cases} \frac{1}{\tan\beta} & \text{for up-type squarks} \\ \tan\beta & \text{for electron-type sleptons and down-type squarks.} \end{cases} \end{aligned}$$

Thereby,  $Q_f$  denotes the electromagnetic charge of the sfermion and  $I_3^f$  the quantum number of the third component of the weak isospin operator.

- (a) How do you obtain the mass eigenstates from there? What are the corresponding masses?
- (b) What happens in the limit of
  - (i) small fermion mass  $m_f$ ,
  - (ii) large SUSY parameters  $M_{\tilde{L},\tilde{E},\tilde{Q},\tilde{D},\tilde{U}} = A_f = \mu =: M_{\text{SUSY}}$ ?

## Exercise 15: Radiative electroweak symmetry breaking

Given are the following simplified  $(A_t = 0, \mu = 0)$  renormalization group equations for the up-like Higgs mass and the two stop masses

$$\frac{\partial}{\partial \log(Q^2/M_{\rm GUT}^2)} \begin{pmatrix} M_{H_u}^2 \\ M_{\tilde{t}_R}^2 \\ M_{\tilde{t}_L}^2 \end{pmatrix} = \frac{\alpha_t}{\pi} \begin{pmatrix} 3 & 3 & 3 \\ 2 & 2 & 2 \\ 1 & 1 & 1 \end{pmatrix} \begin{pmatrix} M_{H_u}^2 \\ M_{\tilde{t}_R}^2 \\ M_{\tilde{t}_L}^2 \end{pmatrix} ,$$

where  $\alpha_t = G_F m_t^2 / (\sqrt{2}\pi)$ .

Solve these equations under the condition  $M_{H_u} = M_{\tilde{t}_R} = M_{\tilde{t}_L} = M_0$  at  $Q = M_{\text{GUT}}$ . Show that the Higgs mass squared  $M_{H_u}^2$  gets negative at  $Q = M_Z$ , if the top mass  $m_t$  is sufficiently large, whereas the stop masses squared remain positive.